# Temperature dependence of Lanczos coefficients and integrability 

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## Outline

## Introduction

## Krylov space and iteration algorithm

## Lanczos coefficients as dynamical variables

## Krylov space

In recent years, Krylov space methods have emerged as a new way to understand complexity, operator growth and chaos.

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Krylov complexity is a measure of operator complexity; it measures the "average position" of an operator within the subspace in which it can evolve.

The exponent of Krylov complexity bounds the exponent of OTOC

$$
\lambda_{\text {отос }} \leq \lambda_{K} .
$$

Parker, Cao, Avdoshkin, Scaffidi, Altman 2019

## Krylov space

Lanczos coefficients $b_{n}$, obtained from the iteration algorithm, encode chaotic behavior. Linear growth of $b_{n} \sim n$ is an indicator of chaos.

## Today's goal

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The $b_{n}$ can be obtained from the 2-point function $C(t)=\langle A(t), A(0)\rangle$.

If $C(t)$ is a thermal 2-point function, the $b_{n}$ depend on $\beta$.
We will show that $b_{n}(\beta)$ satisfy a completely integrable (non-linear) system of equations.

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## Krylov space

Consider a Hamiltonian $H$ and an operator $A$

$$
A(t)=e^{-i H t} A e^{i H t}=\sum_{n=0}^{\infty} \frac{(-i t)^{n} \mathcal{L}^{n}}{n!} A
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where we defined the "super-operator" $\mathcal{L}$ (Liouvillian)

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Under time-evolution, $A(t)$ remains inside the Krylov space

$$
\mathbb{K}=\operatorname{span}\left\{A, \mathcal{L} A, \mathcal{L}^{2} A, \mathcal{L}^{3} A, \ldots\right\} .
$$

## Inner product

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We require that $\rho_{1}, \rho_{2}$ commute with $H$ and that the inner product is semi-positive definite and non-degenerate.

## Lanczos algorithm

We obtain an orthonormal basis for $\mathbb{K}$ using the Lanczos algorithm.
Lanczos algorithm
Let $O_{0}=A$.
For $n=0,1,2, \ldots$ :

$$
\begin{gathered}
a_{n}=\frac{\left\langle O_{n}\right| \mathcal{L}\left|O_{n}\right\rangle}{\left\langle O_{n} \mid O_{n}\right\rangle}, \quad b_{n-1}^{2}=\frac{\left\langle O_{n} \mid O_{n}\right\rangle}{\left\langle O_{n-1} \mid O_{n-1}\right\rangle}, \\
O_{n+1}=\mathcal{L} O_{n}-a_{n} O_{n}-b_{n-1}^{2} O_{n-1} \\
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The set of operators $\left\{A_{0}, A_{1}, A_{2}, \ldots\right\}$ is called the Krylov basis. The sequences $a_{n}, b_{n}$ are called Lanczos coefficients.

## Representation of Liouvillian

The representation of $\mathcal{L}$ in Krylov space written in Krylov basis is, by construction, tridiagonal

$$
\begin{gathered}
\mathcal{L} A_{n}=\sum L_{n m} A_{m}, \\
L=\left(\begin{array}{ccccc}
a_{0} & b_{0} & 0 & 0 & \cdots \\
b_{0} & a_{1} & b_{1} & 0 & \cdots \\
0 & b_{1} & a_{2} & b_{2} & \cdots \\
0 & 0 & b_{2} & a_{3} & \cdots \\
\vdots & \vdots & \vdots & \vdots & \ddots
\end{array}\right) .
\end{gathered}
$$

## Some comments

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Krylov basis, Lanczos coefficients and representation $L$ of Liouvillian depend on choice of inner product.

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Consider the inner product

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The Lanczos coefficients acquire time-dependence: $a_{n}(\tau), b_{n}(\tau)$. Their evolution is governed by a system of completely integrable non-linear equations.
Dymarsky, Gorsky 2019

## Toda chain

Toda chain equations in Lax form

$$
\frac{d}{d \tau} L=[B, L], \quad B=L_{+}-L_{-}
$$

Completely integrable, with the following independent integrals of motion

$$
H_{k}=\operatorname{tr}\left(L^{k}\right)
$$

Explicitly, the equations read

$$
\begin{aligned}
\frac{d}{d \tau} b_{n} & =b_{n}\left(a_{n+1}-a_{n}\right) \\
\frac{d}{d \tau} a_{n} & =2\left(b_{n}^{2}-b_{n-1}^{2}\right)
\end{aligned}
$$

## Temperature dependence

We now consider temperature-dependent inner product (i.e. Wightmann product)

$$
\langle A \mid B\rangle \equiv \operatorname{tr}\left(A^{\dagger} e^{-\beta H / 2} \rho B e^{-\beta H / 2} \rho\right)
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Simplification: Assume that $A \in \operatorname{im}(\mathcal{L})$, and let $\operatorname{dim}(\mathbb{K})=2 N$.

## Representation of $\{H, \cdot\}$

The representation of the operator $\mathcal{J}=\{H, \cdot\}$ in the Krylov space is

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$$

The matrix $J$ satisfies the Lax equation

$$
\frac{d}{d \beta} J=[B, J], \quad B=J_{+}-J_{-} .
$$

This looks similar to Toda, however $J$ is not tridiagonal.

## Even-odd decoupling

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Since $\left\langle A_{2 n+1}(\beta) \mid A_{2 m}\left(\beta^{\prime}\right)\right\rangle=0$, we can write

$$
J=J_{\text {even }} \oplus J_{\text {odd }} .
$$

Better, but we still have $O\left(N^{2}\right)$ parameters.

## Relation between $L, J$

Consider the representation $L$ of $\mathcal{L}=[H, \cdot]$. We saw earlier that this is a tridiagonal matrix.

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The identity

$$
[H,\{H, \cdot\}]=\{H,[H, \cdot]\}
$$

can be written as

$$
[\mathcal{L}, \mathcal{J}]=0 \Longrightarrow[L, J]=0
$$

## Integrability

Now $[L, J]=0$ can be used to determine all entries of $J$ in terms of the diagonal entries of $J$ and $b_{n}$. This reduces the independent parameters to $4 N$.

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This reduces the independent parameters to $4 N$.
The independent integrals of motion are

$$
\begin{array}{ll}
\mathcal{I}_{k}=\operatorname{tr}\left(J_{\text {even }}^{k}\right), & k=1,2, \ldots, N \\
\mathcal{M}_{k}=\operatorname{tr}\left(L^{2 k}\right), & k=1,2, \ldots, N
\end{array}
$$

We have $4 N$-dimensional phase-space and $2 N$ integrals of motion, so this is a fully integrable system.

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Use Lanczos algorithm to tridiagonalize $J_{\text {even }}$ (and $J_{o d d}$ ):

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J_{\text {even }}=C \tilde{J}_{\text {even }} C^{T}, \quad C C^{T}=C^{T} C=1
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Now $\tilde{J}_{\text {even }}$ satisfies Toda equations

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\frac{d}{d \beta} \tilde{J}_{\text {even }}=\left[\tilde{B}_{\text {even }}, \tilde{J}_{\text {even }}\right], \quad \tilde{B}_{\text {even }}=\tilde{J}_{\text {even }}^{+}-\tilde{J}_{\text {even }}^{-} .
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In this basis, we have 2 decoupled Toda chains $J_{\text {even }}, J_{\text {odd }}$.

## Outlook and future directions

- We formulated the temperature dependence of Lanczos coefficients as a fully integrable Hamiltonian system (two decoupled Toda chains).


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- Potential as a powerful numerical or analytical tool.
- Temperature dependence of Lanczos coefficients can be solved as an initial value problem.
- Given a 2 pf at $\beta=0$, we can calculate the 2 pf at finite $\beta$.
- Study scaling of $b_{n}$ with $n$ as $\beta$ is varied.

