Temperature dependence of Lanczos coefficients and integrability

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

June 8, 2023

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Outline

Introduction

Krylov space and iteration algorithm

Lanczos coefficients as dynamical variables

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

In recent years, Krylov space methods have emerged as a new way to understand complexity, operator growth and chaos.

In recent years, Krylov space methods have emerged as a new way to understand complexity, operator growth and chaos.

Krylov complexity is a measure of operator complexity; it measures the "average position" of an operator within the subspace in which it can evolve.

In recent years, Krylov space methods have emerged as a new way to understand complexity, operator growth and chaos.

Krylov complexity is a measure of operator complexity; it measures the "average position" of an operator within the subspace in which it can evolve.

The exponent of Krylov complexity bounds the exponent of OTOC

 $\lambda_{OTOC} \leq \lambda_{K}.$

Parker, Cao, Avdoshkin, Scaffidi, Altman 2019

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Lanczos coefficients b_n , obtained from the iteration algorithm, encode chaotic behavior. Linear growth of $b_n \sim n$ is an indicator of chaos.

Today's goal

The b_n can be obtained from the 2-point function $C(t) = \langle A(t), A(0) \rangle$.

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Today's goal

The b_n can be obtained from the 2-point function $C(t) = \langle A(t), A(0) \rangle$.

If C(t) is a thermal 2-point function, the b_n depend on β .

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Today's goal

The b_n can be obtained from the 2-point function $C(t) = \langle A(t), A(0) \rangle$.

If C(t) is a thermal 2-point function, the b_n depend on β .

We will show that $b_n(\beta)$ satisfy a completely integrable (non-linear) system of equations.

Nick Angelinos University of Kentucky

Introduction 0000

Outline

Introduction

Krylov space and iteration algorithm

Lanczos coefficients as dynamical variables

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Consider a Hamiltonian H and an operator A

$$A(t) = e^{-iHt}Ae^{iHt} = \sum_{n=0}^{\infty} \frac{(-it)^n \mathcal{L}^n}{n!}A,$$

where we defined the "super-operator" \mathcal{L} (Liouvillian)

 $\mathcal{L}A \equiv [H, A].$

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Consider a Hamiltonian H and an operator A

$$A(t) = e^{-iHt}Ae^{iHt} = \sum_{n=0}^{\infty} \frac{(-it)^n \mathcal{L}^n}{n!}A,$$

where we defined the "super-operator" \mathcal{L} (Liouvillian)

$$\mathcal{L}A \equiv [H, A].$$

Under time-evolution, A(t) remains inside the Krylov space

$$\mathbb{K} = \operatorname{span}\{A, \mathcal{L}A, \mathcal{L}^2A, \mathcal{L}^3A, \dots\}.$$

Nick Angelinos University of Kentucky

Inner product

We now want to construct an orthonormal basis for \mathbb{K} .

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Inner product

We now want to construct an orthonormal basis for $\mathbb K.$

We equip the space of operators with an inner product

$$\langle A|B\rangle = tr(A^{\dagger}\rho_1 B\rho_2).$$

Nick Angelinos University of Kentucky

Inner product

We now want to construct an orthonormal basis for $\mathbb K.$

We equip the space of operators with an inner product

$$\langle A|B\rangle = tr(A^{\dagger}\rho_1 B\rho_2).$$

We require that ρ_1 , ρ_2 commute with H and that the inner product is semi-positive definite and non-degenerate.

Nick Angelinos University of Kentucky

Lanczos algorithm

We obtain an orthonormal basis for ${\mathbb K}$ using the Lanczos algorithm.

Lanczos algorithm

Let $O_0 = A$. For n = 0, 1, 2, ...:

$$a_n = rac{\langle O_n | \mathcal{L} | O_n
angle}{\langle O_n | O_n
angle}, \ \ b_{n-1}^2 = rac{\langle O_n | O_n
angle}{\langle O_{n-1} | O_{n-1}
angle},$$
 $O_{n+1} = \mathcal{L}O_n - a_n O_n - b_{n-1}^2 O_{n-1},$
 $A_n = O_n / \sqrt{\langle O_n | O_n
angle}.$

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Lanczos algorithm

We obtain an orthonormal basis for ${\mathbb K}$ using the Lanczos algorithm.

Lanczos algorithm

Let $O_0 = A$. For n = 0, 1, 2, ...:

$$a_n = rac{\langle O_n | \mathcal{L} | O_n
angle}{\langle O_n | O_n
angle}, \quad b_{n-1}^2 = rac{\langle O_n | O_n
angle}{\langle O_{n-1} | O_{n-1}
angle},$$
 $O_{n+1} = \mathcal{L}O_n - a_n O_n - b_{n-1}^2 O_{n-1},$
 $A_n = O_n / \sqrt{\langle O_n | O_n
angle}.$

The set of operators $\{A_0, A_1, A_2, ...\}$ is called the Krylov basis. The sequences a_n, b_n are called Lanczos coefficients.

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Representation of Liouvillian

The representation of \mathcal{L} in Krylov space written in Krylov basis is, by construction, tridiagonal

$$\mathcal{L}A_n = \sum L_{nm}A_m,$$

$$L = \begin{pmatrix} a_0 & b_0 & 0 & 0 & \cdots \\ b_0 & a_1 & b_1 & 0 & \cdots \\ 0 & b_1 & a_2 & b_2 & \cdots \\ 0 & 0 & b_2 & a_3 & \cdots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

`

.

Some comments

Krylov space depends only on the choice of operator A.

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Some comments

Krylov space depends only on the choice of operator A.

Krylov basis, Lanczos coefficients and representation L of Liouvillian depend on choice of inner product.

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Outline

Introduction

Krylov space and iteration algorithm

Lanczos coefficients as dynamical variables

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Euclidean time-evolution

If the inner product depends on a parameter, the Krylov basis and Lanczos coefficients also inherit dependence on that parameter.

Euclidean time-evolution

If the inner product depends on a parameter, the Krylov basis and Lanczos coefficients also inherit dependence on that parameter.

Consider the inner product

$$\langle A|B\rangle_{\tau} \equiv \langle A(\tau)|B\rangle = tr(A^{\dagger}e^{\tau H}\rho_{1}Be^{-\tau H}\rho_{2}).$$

Nick Angelinos University of Kentucky

Euclidean time-evolution

If the inner product depends on a parameter, the Krylov basis and Lanczos coefficients also inherit dependence on that parameter.

Consider the inner product

$$\langle A|B\rangle_{\tau} \equiv \langle A(\tau)|B\rangle = tr(A^{\dagger}e^{\tau H}\rho_1 Be^{-\tau H}\rho_2).$$

The Lanczos coefficients acquire time-dependence: $a_n(\tau), b_n(\tau)$. Their evolution is governed by a system of completely integrable non-linear equations. Dymarsky, Gorsky 2019

Nick Angelinos University of Kentucky

Toda chain equations in Lax form

$$\frac{d}{d\tau}L = [B, L], \quad B = L_+ - L_-.$$

$$H_k = tr(L^k).$$

Explicitly, the equations read

$$rac{d}{d au}b_n=b_n(a_{n+1}-a_n),$$
 $rac{d}{d au}a_n=2(b_n^2-b_{n-1}^2).$

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Temperature dependence

We now consider temperature-dependent inner product (i.e. Wightmann product)

$$\langle A|B
angle \equiv tr(A^{\dagger}e^{-eta H/2}
ho Be^{-eta H/2}
ho).$$

Nick Angelinos University of Kentucky

Temperature dependence

We now consider temperature-dependent inner product (i.e. Wightmann product)

$$\langle A|B
angle \equiv tr(A^{\dagger}e^{-eta H/2}
ho Be^{-eta H/2}
ho).$$

Now the Lanczos coefficients acquire temperature dependence $b_n(\beta)$.

Nick Angelinos University of Kentucky

Temperature dependence

We now consider temperature-dependent inner product (i.e. Wightmann product)

$$\langle A|B
angle \equiv tr(A^{\dagger}e^{-eta H/2}
ho Be^{-eta H/2}
ho).$$

Now the Lanczos coefficients acquire temperature dependence $b_n(\beta)$.

Simplification: Assume that $A \in im(\mathcal{L})$, and let dim $(\mathbb{K}) = 2N$.

Nick Angelinos University of Kentucky

Representation of $\{H, \cdot\}$

The representation of the operator $\mathcal{J} = \{H, \cdot\}$ in the Krylov space is

$$\mathcal{J}A_n=\sum_{m=0}^{2N-1}J_{nm}A_m.$$

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Representation of $\{H, \cdot\}$

The representation of the operator $\mathcal{J} = \{H, \cdot\}$ in the Krylov space is

$$\mathcal{J}A_n=\sum_{m=0}^{2N-1}J_{nm}A_m.$$

The matrix J satisfies the Lax equation

$$\frac{d}{d\beta}J = [B, J], \quad B = J_+ - J_-.$$

This looks similar to Toda, however J is not tridiagonal.

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Even-odd decoupling

The matrix J seemingly has too many parameters $(O(N^2))$. We expect O(N) parameters.

Even-odd decoupling

The matrix J seemingly has too many parameters $(O(N^2))$. We expect O(N) parameters.

Since
$$\langle A_{2n+1}(\beta)|A_{2m}(\beta')\rangle = 0$$
, we can write

$$J = J_{even} \oplus J_{odd}.$$

Better, but we still have $O(N^2)$ parameters.

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Relation between L, J

Consider the representation L of $\mathcal{L} = [H, \cdot]$. We saw earlier that this is a tridiagonal matrix.

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Relation between L, J

Consider the representation L of $\mathcal{L} = [H, \cdot]$. We saw earlier that this is a tridiagonal matrix.

The identity

$$[H, \{H, \cdot\}] = \{H, [H, \cdot]\}$$

can be written as

$$[\mathcal{L},\mathcal{J}]=0\implies [L,J]=0.$$

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Integrability

Now [L, J] = 0 can be used to determine all entries of J in terms of the diagonal entries of J and b_n . This reduces the independent parameters to 4N.

Integrability

Now [L, J] = 0 can be used to determine all entries of J in terms of the diagonal entries of J and b_n . This reduces the independent parameters to 4N. The independent integrals of motion are

$$\mathcal{I}_k = tr(J_{even}^k), \quad k = 1, 2, \dots, N$$
$$\mathcal{M}_k = tr(L^{2k}), \quad k = 1, 2, \dots, N.$$

We have 4N-dimensional phase-space and 2N integrals of motion, so this is a fully integrable system.

Nick Angelinos University of Kentucky

Can we relate these equations to Toda chain?

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Can we relate these equations to Toda chain?

Use Lanczos algorithm to tridiagonalize J_{even} (and J_{odd}):

$$J_{even} = C \tilde{J}_{even} C^T, \quad C C^T = C^T C = 1.$$

Nick Angelinos University of Kentucky

Can we relate these equations to Toda chain?

Use Lanczos algorithm to tridiagonalize J_{even} (and J_{odd}):

$$J_{even} = C \tilde{J}_{even} C^T, \quad C C^T = C^T C = 1.$$

Now
$$\tilde{J}_{even}$$
 satisfies Toda equations

$$\frac{d}{d\beta}\tilde{J}_{even} = [\tilde{B}_{even}, \tilde{J}_{even}], \quad \tilde{B}_{even} = \tilde{J}^+_{even} - \tilde{J}^-_{even}.$$

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

Can we relate these equations to Toda chain?

Use Lanczos algorithm to tridiagonalize J_{even} (and J_{odd}):

$$J_{even} = C \tilde{J}_{even} C^T, \quad C C^T = C^T C = 1.$$

Now \tilde{J}_{even} satisfies Toda equations

$$rac{d}{deta} ilde{J}_{even} = [ilde{B}_{even}, ilde{J}_{even}], \quad ilde{B}_{even} = ilde{J}^+_{even} - ilde{J}^-_{even}.$$

In this basis, we have 2 decoupled Toda chains J_{even}, J_{odd} .

Nick Angelinos University of Kentucky

Quantum Information in Quantum Field Theory and Cosmology, BIRS

We formulated the temperature dependence of Lanczos coefficients as a fully integrable Hamiltonian system (two decoupled Toda chains).

- We formulated the temperature dependence of Lanczos coefficients as a fully integrable Hamiltonian system (two decoupled Toda chains).
- Potential as a powerful numerical or analytical tool.

Nick Angelinos University of Kentucky

- We formulated the temperature dependence of Lanczos coefficients as a fully integrable Hamiltonian system (two decoupled Toda chains).
- Potential as a powerful numerical or analytical tool.
- Temperature dependence of Lanczos coefficients can be solved as an initial value problem.

- We formulated the temperature dependence of Lanczos coefficients as a fully integrable Hamiltonian system (two decoupled Toda chains).
- Potential as a powerful numerical or analytical tool.
- Temperature dependence of Lanczos coefficients can be solved as an initial value problem.
- Given a 2pf at $\beta = 0$, we can calculate the 2pf at finite β .

Nick Angelinos University of Kentucky

- We formulated the temperature dependence of Lanczos coefficients as a fully integrable Hamiltonian system (two decoupled Toda chains).
- Potential as a powerful numerical or analytical tool.
- Temperature dependence of Lanczos coefficients can be solved as an initial value problem.
- Given a 2pf at $\beta = 0$, we can calculate the 2pf at finite β .
- Study scaling of b_n with n as β is varied.

Nick Angelinos University of Kentucky