

A large, oval-shaped map of the Cosmic Microwave Background (CMB) showing temperature fluctuations. The map is filled with a dense pattern of blue and orange dots, representing the distribution of matter and energy in the early universe. The title "Cosmological Complexity" is overlaid in the center in a large, white, bold font with a black outline.

Cosmological Complexity

Bret Underwood

Pacific Lutheran University, Tacoma, WA, USA



Based on

Bhattacharyya, Das, Haque, **BU**, 2001.08664

Bhattacharyya, Das, Haque, **BU**, 2005.10854

Haque, Jana, **BU**, 2107.08969

Haque, Jana, **BU**, 2110.08356

Cosmic Microwave Background

Small - Scale
Quantum Fluctuations

$$\hat{\zeta}_k$$

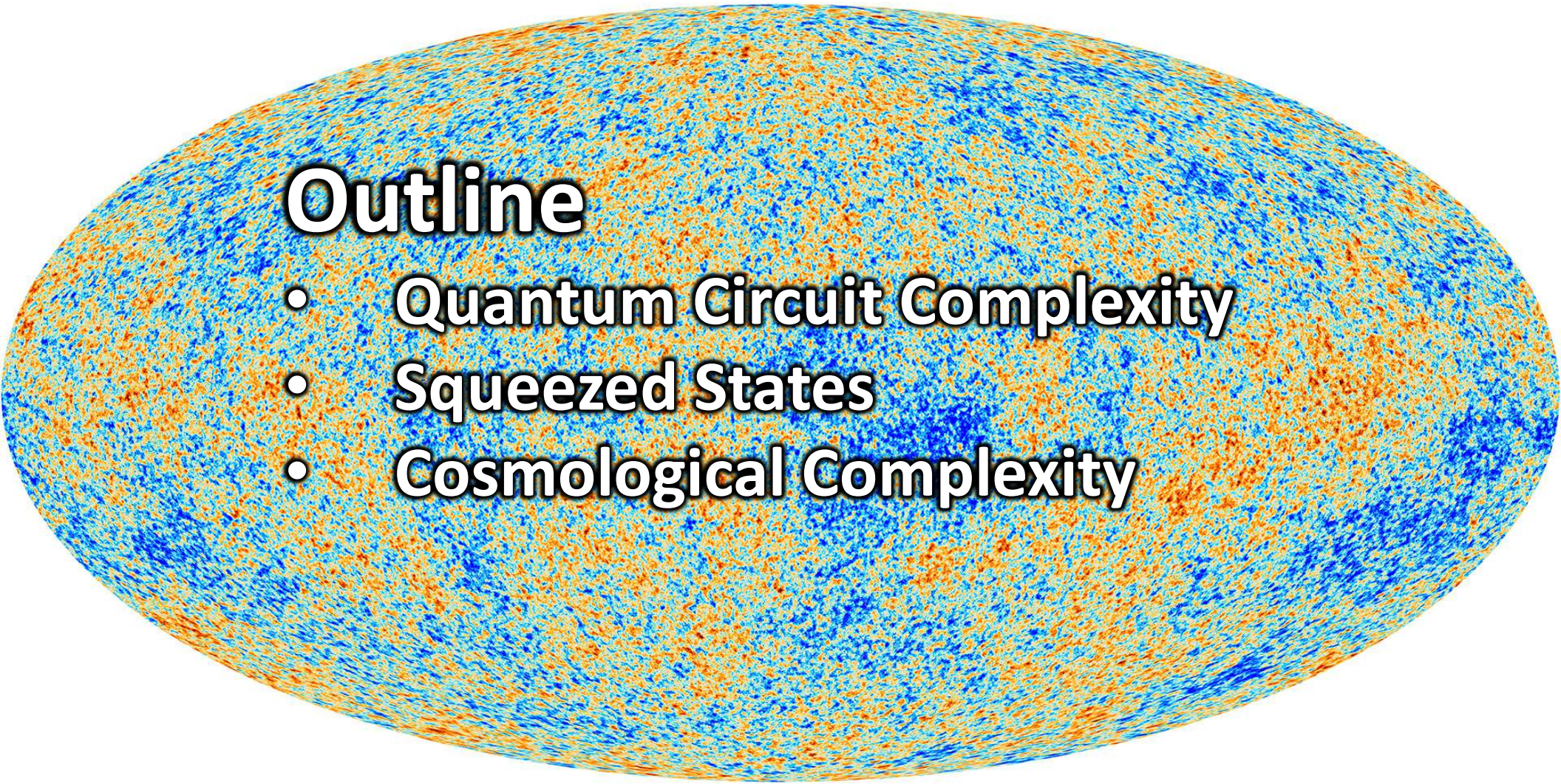
$$\hat{u}_{\text{cosmo}}$$

Large - Scale Density
Perturbations

$$\langle \delta T(k) \delta T(k') \rangle$$

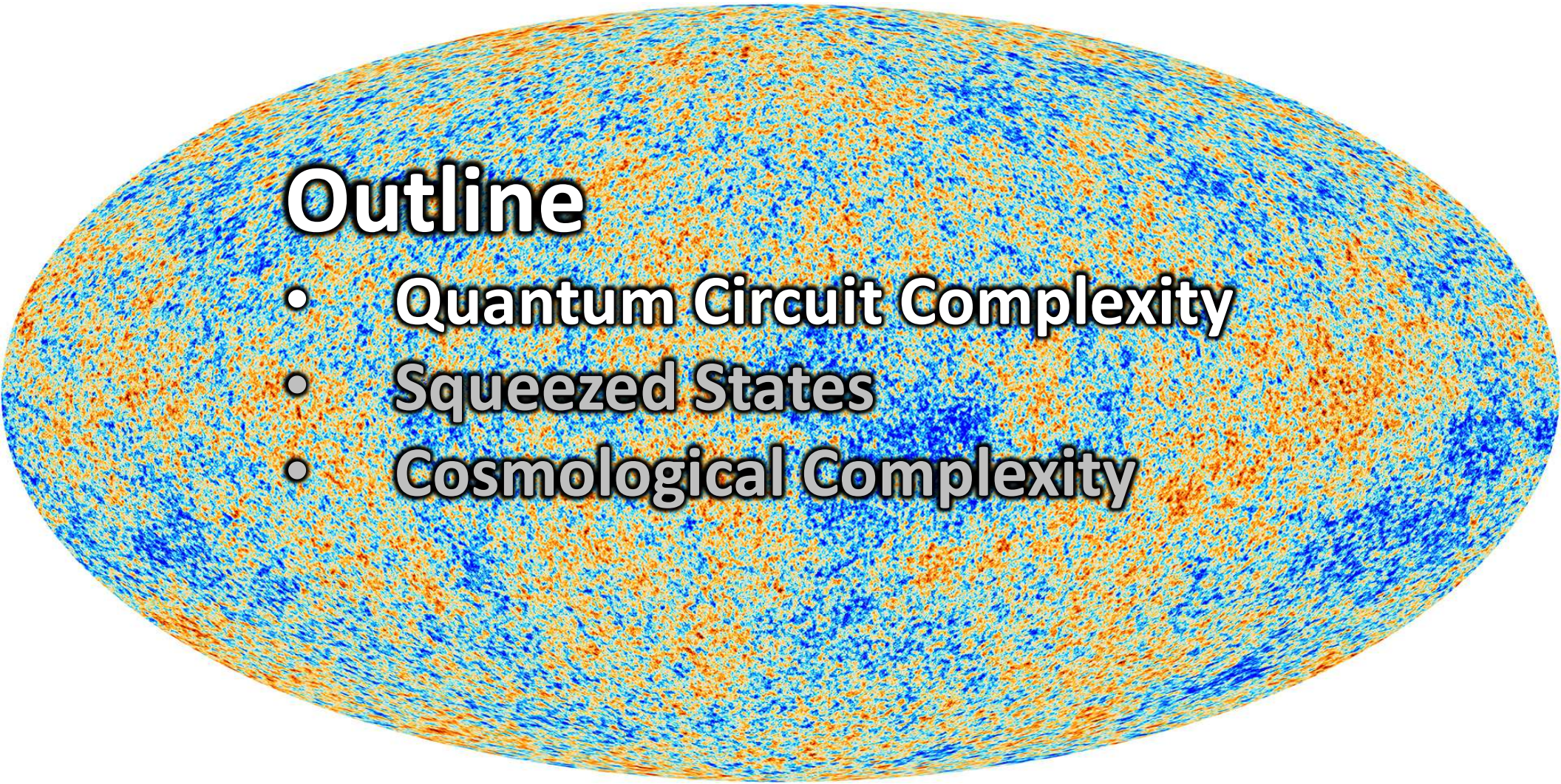
What is the (quantum circuit)
complexity of this process?

- Growth of complexity with time?
- Bounds on the growth of complexity?
- Total complexity of observed universe?



Outline

- **Quantum Circuit Complexity**
- **Squeezed States**
- **Cosmological Complexity**



Outline

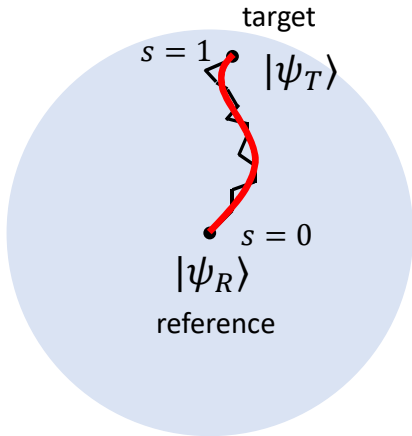
- **Quantum Circuit Complexity**
- **Squeezed States**
- **Cosmological Complexity**

Quantum Circuit Complexity

Unitary evolution from reference state $|\psi_R\rangle$ to target state $|\psi_T\rangle$

$$|\psi_R\rangle \xrightarrow{\hat{U}_{\text{target}}} |\psi_T\rangle \quad \text{e.g. } e^{-i\hat{H}t}$$

$$|\psi_T\rangle = \hat{U}_{\text{target}} |\psi_R\rangle$$



- Model as continuous application of operators Nielsen et al

$$\hat{U}_{\text{target}} = \bar{P} \exp \left[\int_0^1 V^I(s) \hat{O}_I ds \right]$$

$\{\hat{O}_I\}$: basis of gates
 $V^I(s)$: tangent vectors

- Assign a circuit depth to path

$$\mathcal{D} = \mathcal{D}[V^I]$$

$$\mathcal{D}[V^I] = \int_0^1 \sqrt{G_{IJ} V^I V^J} ds \quad \text{with } G_{IJ} = \delta_{IJ}$$

“gate cost”

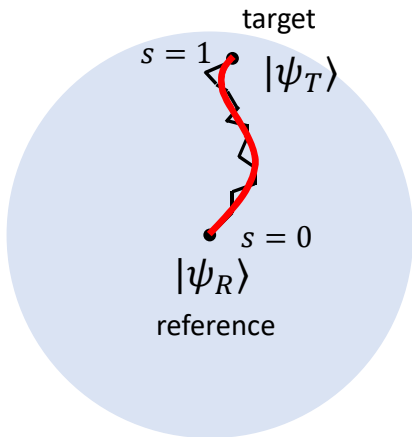
- Circuit Complexity is depth minimized over paths

$$\mathcal{C} = \min_{\{V^I\}} \mathcal{D}[V^I]$$

Quantum Circuit Complexity

Unitary evolution from reference state $|\psi_R\rangle$ to target state $|\psi_T\rangle$

$$|\psi_R\rangle \xrightarrow{\hat{U}_{\text{target}}} |\psi_T\rangle \quad |\psi_T\rangle = \hat{U}_{\text{target}} |\psi_R\rangle$$



- Model as continuous application of operators Nielsen et al

$$\hat{U}_{\text{target}} = \bar{P} \exp \left[\int_0^1 V^I(s) \hat{O}_I ds \right]$$

$\{\hat{O}_I\}$: basis of gates
 $V^I(s)$: tangent vectors

Operator Circuit Complexity

- Characterize gates by structure constants

$$[\hat{O}_I, \hat{O}_J] = i f_{IJ}^K \hat{O}_K$$

- Minimization:
 \Rightarrow Euler-Arnold eq on group manifold

$$G_{IJ} \frac{dV^J}{ds} = f_{IJ}^K G_{KL} V^J V^L$$

- Advantage:** Not restricted to subset of states
- Disadvantage:** Euler-Arnold eq can be difficult to solve

Balasubramanian, Decross, Kar, Parrikar
 Basteiro, Erdmenger, Fries, Goth, Matthaiakakis, Meyer

(Gaussian) State Circuit Complexity

- Characterize target operator by its action on Gaussian states

$$\langle x | \psi_R \rangle \sim e^{-\frac{1}{2} \omega_0 \sum_k x_k^2}$$

$$\hat{O}_k \sim e^{-i \hat{x}_k \hat{p}_k} \longrightarrow \langle x | \psi_T \rangle \sim e^{-\frac{1}{2} \sum_k \Omega_k x_k^2}$$

$\{\hat{O}_I\}$: basis for $GL(N, \mathbb{R})$ or $GL(N, \mathbb{C})$

- Advantage:** Simple to set up and find optimal path
- Disadvantage:** Restricted to Gaussian states

Jefferson, Myers
 Ali, Bhattacharyya, Haque, Kim, Moynihan, Murugan

also **K-Complexity/Spread Complexity**

Dixit, Magan, Kim, Dymarsky, Watanabe

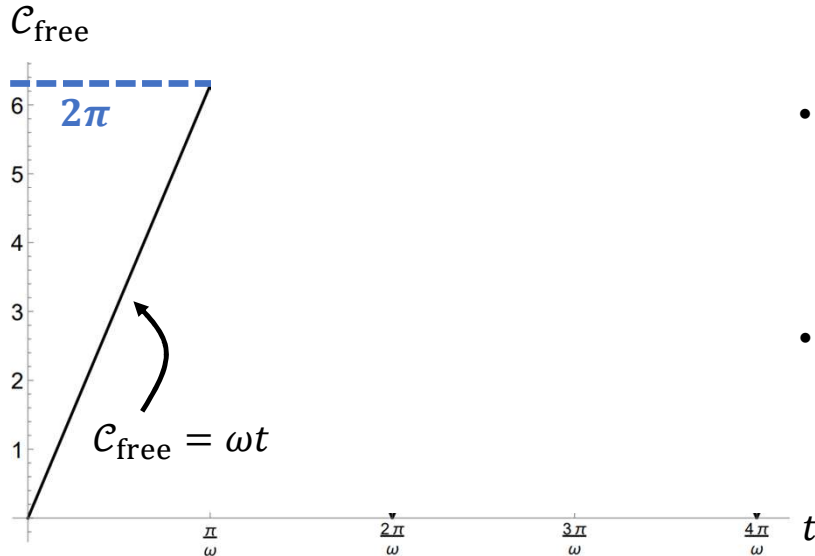
Complexity: Free Harmonic Oscillator

Example: Free Harmonic Oscillator

$$|\psi_T\rangle = \hat{U}_{\text{target}} |\psi_R\rangle$$

$$|\psi_T\rangle = e^{-i\hat{H}_0 t} |\psi_R\rangle$$

$$\hat{H}_0 = \frac{\omega}{2} (\hat{a}^\dagger \hat{a} + \hat{a} \hat{a}^\dagger) \quad \text{Haque, Jana, BU}$$



Operator Circuit Complexity

- Model as continuous application of operators

$$\hat{U}_{\text{target}} = \tilde{P} \exp \left[\int_0^1 V^I(s) \hat{O}_I ds \right] \quad \begin{array}{l} \{\hat{O}_I\}: \\ V^I(s): \end{array} \quad \begin{array}{l} \text{basis of gates} \\ \text{tangent vectors} \end{array}$$

$$\left\{ \hat{O}_1 = \frac{\hat{a}^2 + \hat{a}^{\dagger 2}}{4} \quad \hat{O}_2 = i \frac{\hat{a}^2 - \hat{a}^{\dagger 2}}{4} \quad \hat{O}_3 = \frac{\hat{a}^\dagger \hat{a} + \hat{a} \hat{a}^\dagger}{4} \right\}$$

- Characterize gates by structure constants $[\hat{O}_I, \hat{O}_J] = i f_{IJ}^K \hat{O}_K$

$$[\hat{O}_1, \hat{O}_2] = -i\hat{O}_3, \quad [\hat{O}_3, \hat{O}_1] = i\hat{O}_2, \quad [\hat{O}_2, \hat{O}_3] = i\hat{O}_1 \quad \mathbf{su(1,1)}$$

- Minimization:

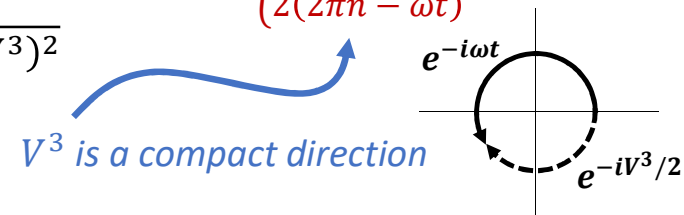
\Rightarrow Euler-Arnold eq on group manifold ($G_{IJ} = \delta_{IJ}$)

$$G_{IJ} \frac{dV^J}{ds} = f_{IJ}^K G_{KL} V^J V^L \quad \longrightarrow \quad \begin{array}{l} V^1 = 0, \\ V^2 = 0, \end{array}$$

- Complexity

$$\begin{aligned} C_{\text{free}} &= \sqrt{(V^1)^2 + (V^2)^2 + (V^3)^2} \\ &= |V^3| \end{aligned}$$

$$V^3 = \min \begin{cases} 2(\omega t - 2\pi n) \\ 2(2\pi n - \omega t) \end{cases}$$



Complexity: Free Scalar Field

Free scalar field ϕ in $(d + 1)$ -dimension, mass m , box L

Haque, Jana, BU

$$\hat{\phi} = \sum_{\vec{n}}^{N_{\max}} \frac{1}{\sqrt{2 E_{\vec{n}}}} \left(\hat{a}_{\vec{n}} e^{i\vec{p}_{\vec{n}} \cdot \vec{x}} + \hat{a}_{\vec{n}}^\dagger e^{-i\vec{p}_{\vec{n}} \cdot \vec{x}} \right)$$

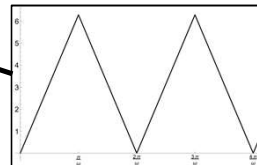
$\Lambda = N_{\max}\pi/L$ UV cutoff

Mode expansion:
$$\begin{cases} \vec{p}_{\vec{n}} = \vec{n}\pi/L \\ E_{\vec{n}} = \sqrt{p_{\vec{n}}^2 + m^2} \end{cases}$$

Target Unitary

$$\mathcal{U}_{\text{target}} = \prod_{\vec{n}}^{N_{\max}} e^{-i\frac{1}{2}E_{\vec{n}} (\hat{a}_{\vec{n}}^\dagger \hat{a}_{\vec{n}} + \hat{a}_{\vec{n}} \hat{a}_{\vec{n}}^\dagger)}$$

copies of free oscillator
for each mode



Complexity of free scalar field

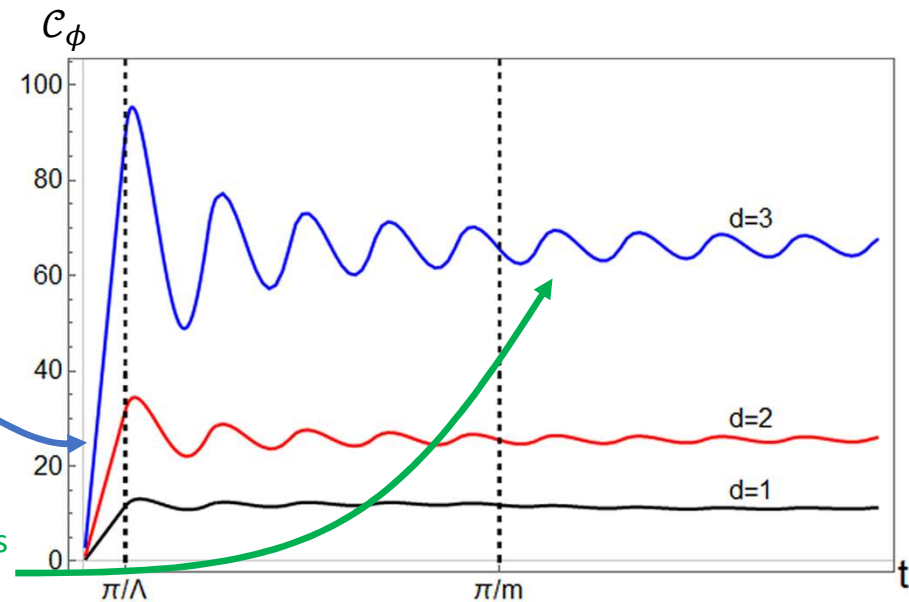
$$\mathcal{C}_\phi = \sqrt{\sum_{\vec{n}}^{N_{\max}} (V_{\vec{n}}^3)^2} \sim L^{d/2} \sqrt{\int^\Lambda (V^3(p))^2 d^d p}$$

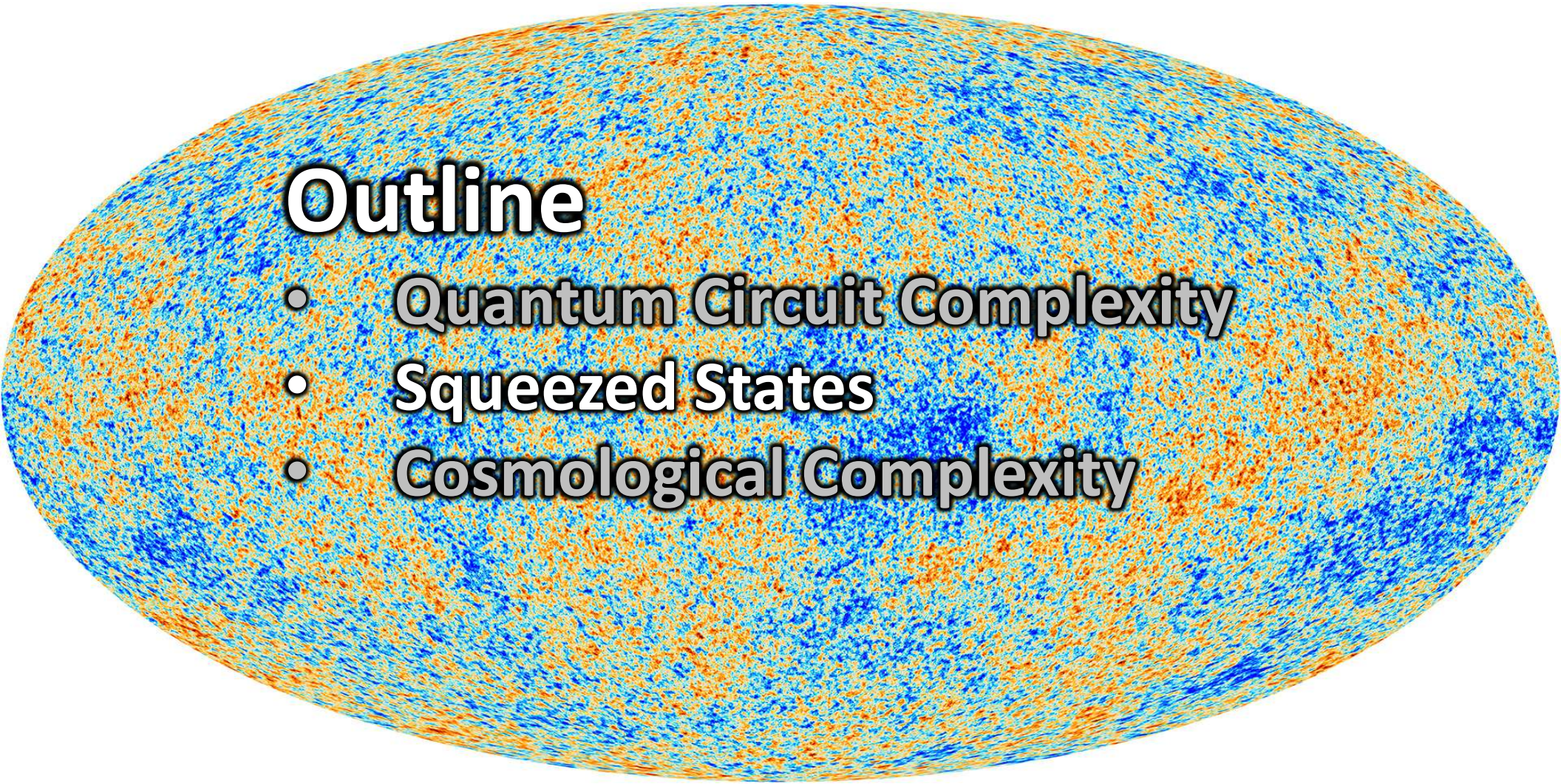
continuum limit

$$\sim \begin{cases} L^{d/2} \Lambda^{d/2} (\Lambda t) & \text{early times } t \ll \pi/\Lambda \\ L^{d/2} \Lambda^{d/2} & \text{late times } t \gg \pi/\Lambda \end{cases}$$

Linear Growth:
complexity of only one
mode growing

Saturation:
complexity of all modes
oscillating, average out





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- **Cosmological Complexity**

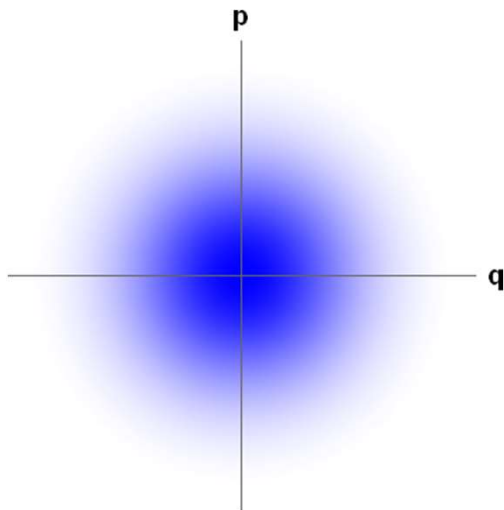
Squeezed States

Vacuum States

$$\psi(q) \sim e^{-\frac{q^2}{2}}, \quad \varphi(p) \sim e^{-\frac{p^2}{2}}$$

$$\langle \Delta \hat{q}^2 \rangle = \frac{1}{2} \quad \langle \Delta \hat{p}^2 \rangle = \frac{1}{2}$$

$$\langle \Delta \hat{q}^2 \rangle \langle \Delta \hat{p}^2 \rangle = \frac{1}{4}$$



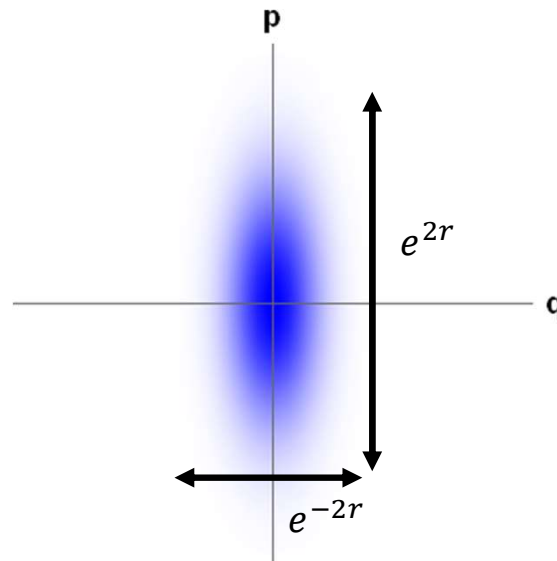
$r = 0$

Squeezed Vacuum State

$$\psi(q) \sim e^{-\frac{q^2}{2}e^{2r}}, \quad \varphi(p) \sim e^{-\frac{p^2}{2}e^{-2r}}$$

$$\langle \Delta \hat{q}^2 \rangle = \frac{1}{2}e^{-2r} \quad \langle \Delta \hat{p}^2 \rangle = \frac{1}{2}e^{2r}$$

$$\langle \Delta \hat{q}^2 \rangle \langle \Delta \hat{p}^2 \rangle = \frac{1}{4}$$



$r = 0.5$

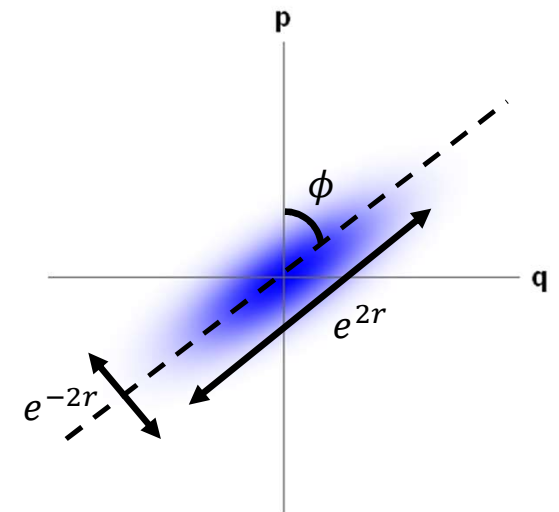
$\phi = 0$

Squeezed, **Rotated**
Vacuum State

$$\hat{q}_+ = \hat{p} \sin \phi + \hat{x} \cos \phi$$

$$\hat{q}_- = \hat{p} \cos \phi - \hat{x} \sin \phi$$

$$\langle \Delta \hat{q}_+^2 \rangle \langle \Delta \hat{q}_-^2 \rangle = \frac{1}{4}$$



$r = 0.5$

$\phi = \pi/3$

Squeezed States

Described by **squeezing parameter r** , **squeezing angle ϕ** , and **rotation angle θ**

$$|r, \phi, \theta\rangle = \hat{S}(r, \phi) \hat{R}(\theta) |0\rangle$$

$$\hat{U} = \hat{S}(r, \phi) \hat{R}(\theta)$$

where $\hat{S}(r, \phi) \equiv \exp\left[\frac{r}{2} (e^{-2i\phi} \hat{a}^2 - e^{2i\phi} \hat{a}^{\dagger 2})\right]$ **squeezing operator**

$\hat{R}(\theta) \equiv \exp[-i\theta (\hat{a}^\dagger \hat{a} + \hat{a} \hat{a}^\dagger)]$ **rotation operator**

“World record”
laboratory squeezing
 $r \approx 1.7$

Vahlbruch, et al, 2016

Squeezed States found in:

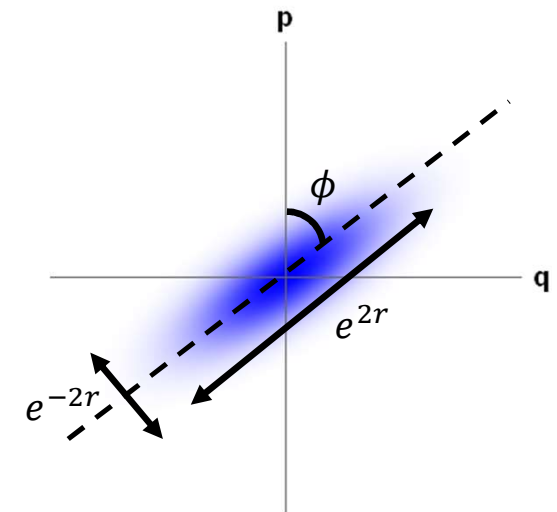
- Quantum Optics
- Gravitational Wave Detection
- **Cosmological Perturbations**

Squeezed, **Rotated**
Vacuum State

$$\hat{q}_+ = \hat{p} \sin \phi + \hat{x} \cos \phi$$

$$\hat{q}_- = \hat{p} \cos \phi - \hat{x} \sin \phi$$

$$\langle \Delta \hat{q}_+^2 \rangle \langle \Delta \hat{q}_-^2 \rangle = \frac{1}{4}$$



$$r = 0.5$$

$$\phi = \pi/3$$

Squeezed States in Cosmological Perturbations

Cosmological Perturbations

$$ds^2 = a(\eta)^2 (-d\eta^2 + (1 - 2\mathcal{R}) d\vec{x}^2)$$

scale factor \rightarrow $a(\eta)$
 conformal time \rightarrow η
 curvature pert \rightarrow \mathcal{R}

Mukhanov variable

$$v = z \mathcal{R}, \quad z = a\sqrt{2\epsilon}$$

Canonical Quantization

$$\hat{v} = \int \frac{d^3k}{(2\pi)^3} \hat{v}_k e^{i\vec{k}\cdot\vec{x}} \quad \hat{v}_k = \frac{1}{\sqrt{2k}} (\hat{a}_{\vec{k}} + \hat{a}_{-\vec{k}}^\dagger)$$

$$\hat{H} = \int d^3k \hat{H}_{\vec{k}} = \int d^3k \frac{1}{2} \left[k (\hat{a}_{\vec{k}} \hat{a}_{\vec{k}}^\dagger + \hat{a}_{-\vec{k}}^\dagger \hat{a}_{-\vec{k}}) - i \frac{z'}{z} (\hat{a}_{\vec{k}} \hat{a}_{-\vec{k}} - \hat{a}_{\vec{k}}^\dagger \hat{a}_{-\vec{k}}^\dagger) \right]$$

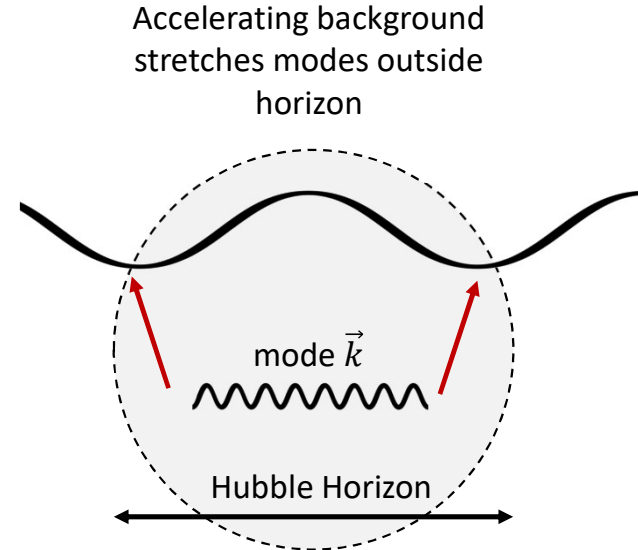
Free-particle

Inverted Oscillator

Time-dependent frequency

Two-mode squeezed state $(\vec{k}, -\vec{k})$

Grishchuk, Sidorov
Albrecht, Ferreira, Joyce, Prokopec



Squeezed States in Cosmological Perturbations

Squeezed Cosmological Perturbations

$$|r_k, \phi_k, \theta_k\rangle = \hat{S}(r_k, \phi_k) \hat{R}(\theta_k) |0_{\vec{k}}, 0_{-\vec{k}}\rangle$$

Example:

de Sitter (inflation)

Squeezing Parameters:

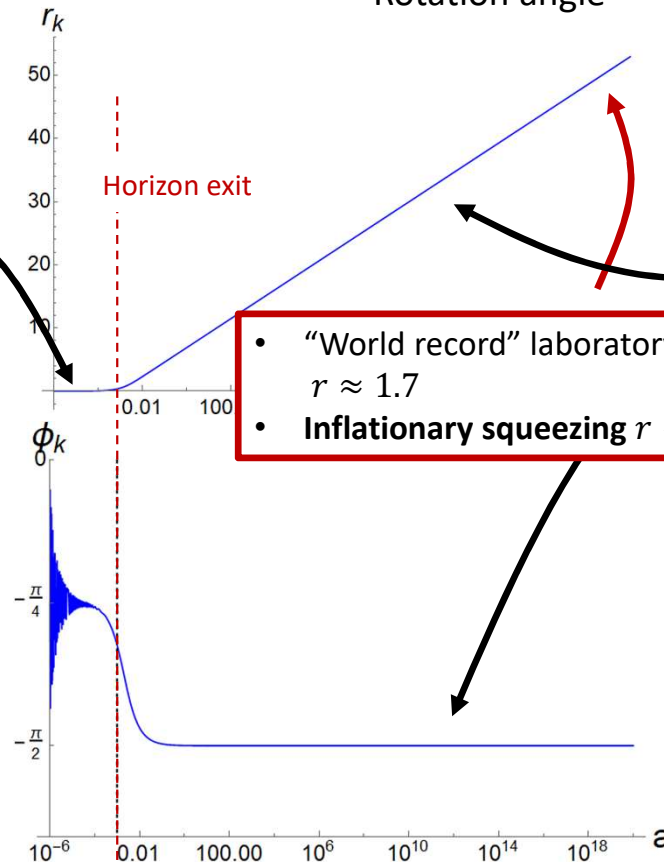
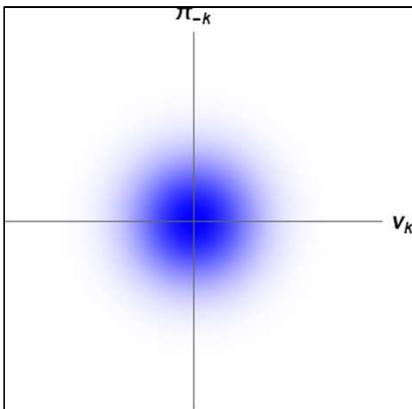
Squeezing strength
Squeezing angle
Rotation angle

$$\left. \begin{aligned} r_k &= r_k(\eta) \\ \phi_k &= \phi_k(\eta) \\ \theta_k &= \theta_k(\eta) \end{aligned} \right\} \text{Time-dependence from expanding background}$$

Inside Horizon:

$$k \gg \frac{a'}{a}$$

- $r_k \ll 1$
- ϕ_k, θ_k evolve

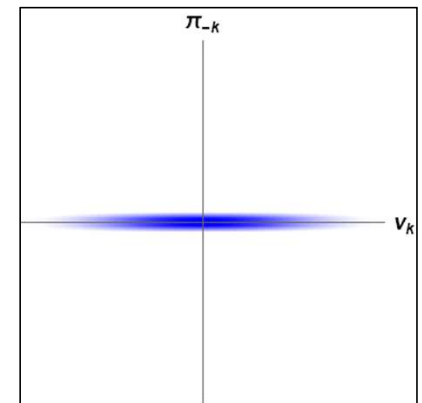


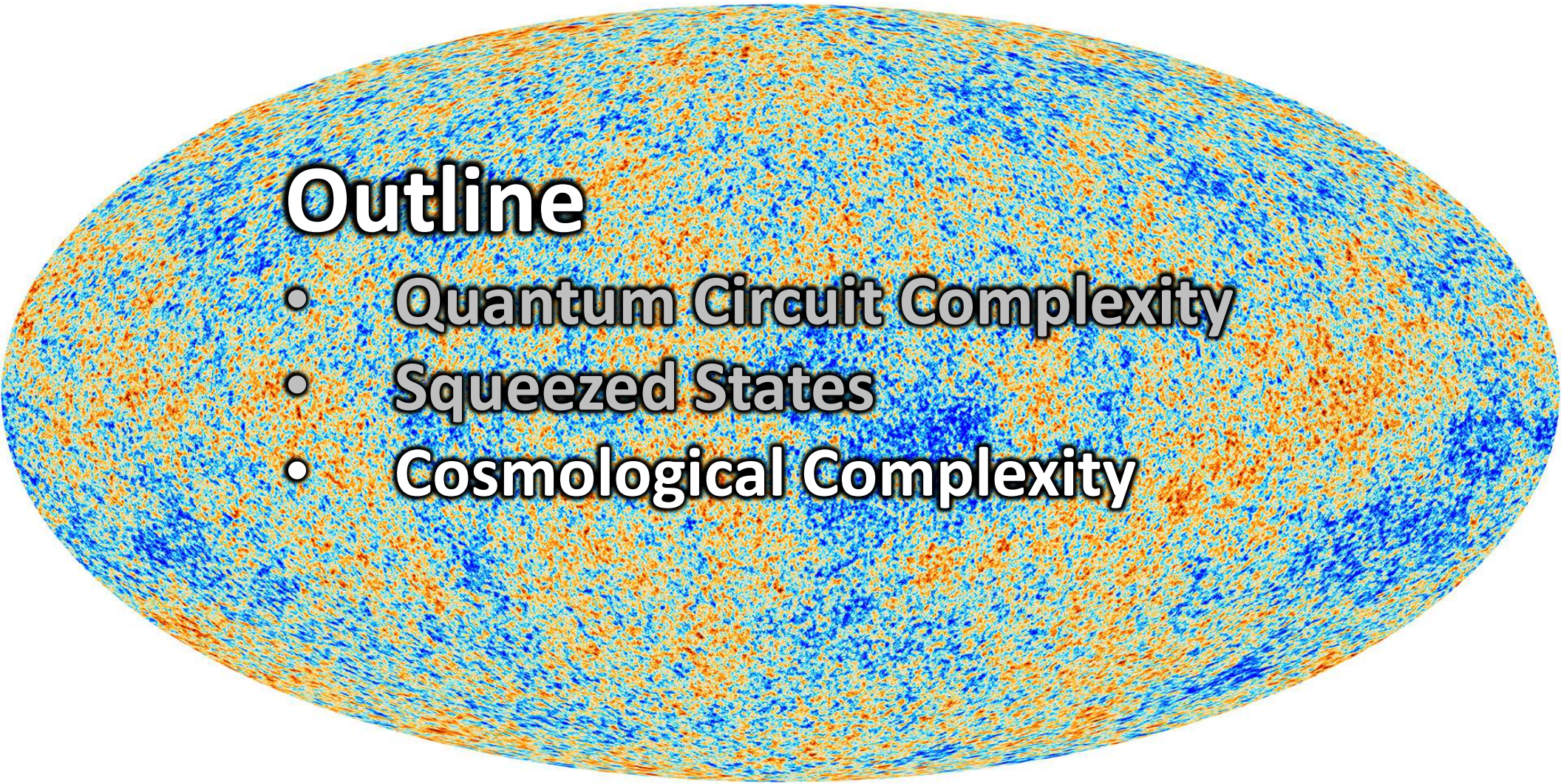
- "World record" laboratory squeezing $r \approx 1.7$
- **Inflationary squeezing** $r \sim N_e \sim 60!$

Outside Horizon:

$$k \ll \frac{a'}{a}$$

- $r_k \sim \ln a$
- $\phi_k \sim \pi/2$ (constant)
- θ_k constant





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- **Cosmological Complexity**

Cosmological Complexity

the complexity of quantum cosmological perturbations

Squeezed
Cosmological
Perturbations

$$|\psi_T\rangle = \hat{U}_{\text{cosmo}} |\psi_R\rangle$$

$$|r_k, \phi_k, \theta_k\rangle = \hat{S}(r_k, \phi_k) \hat{\mathcal{R}}(\theta_k) |0_{\vec{k}}, 0_{-\vec{k}}\rangle$$

Operator Circuit Complexity Haque, Jana, BU

$$\hat{U}_{\text{target}} = \bar{P} \exp \left[\int_0^1 V^I(s) \hat{O}_I ds \right]$$

$$\hat{S} = \exp \left[\frac{r_k}{2} \left(e^{-2i\phi_k} \hat{a}_{\vec{k}} \hat{a}_{-\vec{k}} - e^{2i\phi_k} \hat{a}_{\vec{k}}^\dagger \hat{a}_{-\vec{k}}^\dagger \right) \right] \text{ squeezing operator}$$

$$\hat{\mathcal{R}} = \exp \left[-i\theta_k \left(\hat{a}_{\vec{k}} \hat{a}_{\vec{k}}^\dagger + \hat{a}_{-\vec{k}}^\dagger \hat{a}_{-\vec{k}} \right) \right] \text{ rotation operator}$$

- Characterize gates by structure constants

$$[\hat{O}_I, \hat{O}_J] = i f_{IJ}^K \hat{O}_K$$

$$\left. \begin{aligned} \hat{O}_1 &= \frac{\hat{a}_{\vec{k}} \hat{a}_{-\vec{k}} + \hat{a}_{\vec{k}}^\dagger \hat{a}_{-\vec{k}}^\dagger}{2} \\ \hat{O}_2 &= i \frac{\hat{a}_{\vec{k}} \hat{a}_{-\vec{k}} - \hat{a}_{\vec{k}}^\dagger \hat{a}_{-\vec{k}}^\dagger}{2} \\ \hat{O}_3 &= \frac{\hat{a}_{\vec{k}} \hat{a}_{\vec{k}}^\dagger + \hat{a}_{-\vec{k}}^\dagger \hat{a}_{-\vec{k}}}{2} \end{aligned} \right\} \text{SU}(1,1)$$

- Minimization:

⇒ Euler-Arnold eq on group manifold

$$G_{IJ} \frac{dV^J}{ds} = f_{IJ}^K G_{KL} V^J V^L$$

↖ $V^I(s)$: tangent vectors

⇒ solve for $V^J(s)$, construct $\mathcal{U}_{\text{target}}$

- Operator Circuit

$$\text{Complexity } \mathcal{C}^{(0)} = \sqrt{G_{IJ} V^I V^J}$$

Cosmological Complexity

the complexity of quantum cosmological perturbations

Squeezed
Cosmological
Perturbations

$$|\psi_T\rangle = \hat{U}_{\text{cosmo}} |\psi_R\rangle$$

$$|r_k, \phi_k, \theta_k\rangle = \hat{S}(r_k, \phi_k) \hat{\mathcal{R}}(\theta_k) |0_{\vec{k}}, 0_{-\vec{k}}\rangle$$

Operator Circuit Complexity Haque, Jana, BU

$$\hat{U}_{\text{target}} = \tilde{P} \exp \left[\int_0^1 V^I(s) \hat{O}_I ds \right]$$

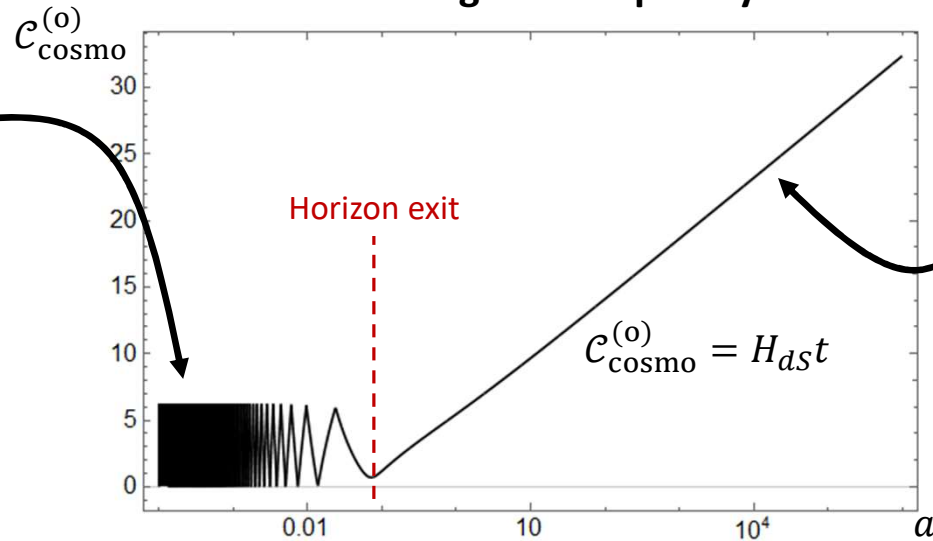
$$\hat{S} = \exp \left[\frac{r_k}{2} \left(e^{-2i\phi_k} \hat{a}_{\vec{k}} \hat{a}_{-\vec{k}} - e^{2i\phi_k} \hat{a}_{\vec{k}}^\dagger \hat{a}_{-\vec{k}}^\dagger \right) \right] \quad \text{squeezing operator}$$

$$\hat{\mathcal{R}} = \exp \left[-i\theta_k \left(\hat{a}_{\vec{k}} \hat{a}_{\vec{k}}^\dagger + \hat{a}_{-\vec{k}}^\dagger \hat{a}_{-\vec{k}} \right) \right] \quad \text{rotation operator}$$

Inside Horizon:

- $r_k \ll 1$
- Complexity oscillates
- $\mathcal{C}_{\text{cosmo}}^{(o)} \leq 2\pi$ due to rotation phase θ

Cosmological Complexity



Outside Horizon:

- Grows with squeezing, e-folds
- $\mathcal{C}_{\text{cosmo}}^{(o)} \approx r_k \approx \ln \left(\frac{a}{a_{\text{exit}}} \right)$
- Growth is linear in cosmic time t , $a(t) = e^{H_d s t}$

Cosmological Complexity

the complexity of quantum cosmological perturbations

Squeezed
Cosmological
Perturbations

$$|\psi_T\rangle = \hat{U}_{\text{cosmo}} |\psi_R\rangle$$

$$|r_k, \phi_k, \theta_k\rangle = \hat{S}(r_k, \phi_k) \hat{\mathcal{R}}(\theta_k) |0_{\vec{k}}, 0_{-\vec{k}}\rangle$$

(Gaussian) State Circuit Complexity

Bhattacharyya, Das, Haque, BU

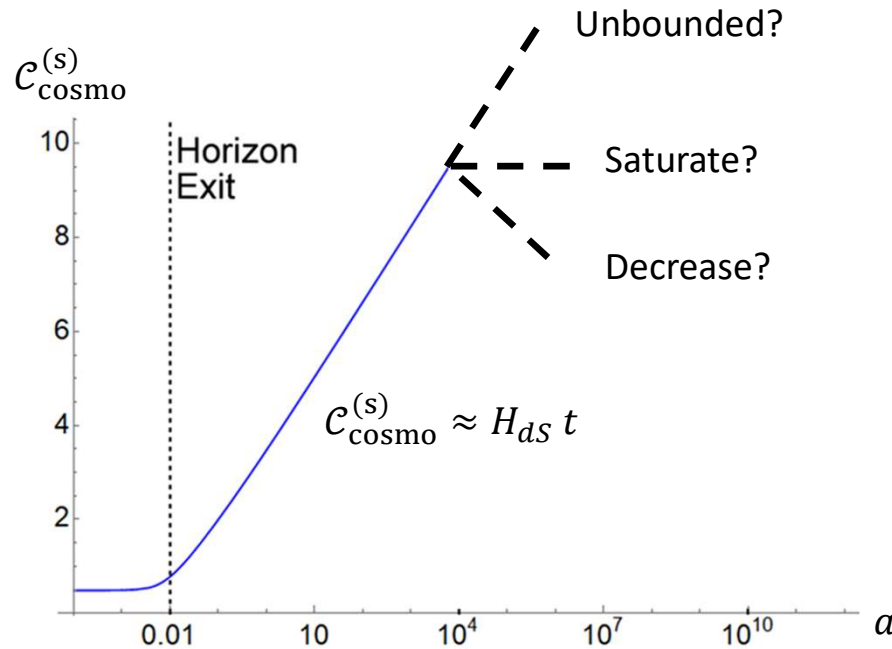
- Characterize \hat{U}_{cosmo} by its action on Gaussian states

$$\Psi_{\text{sq}} = \langle q_{\vec{k}}, q_{-\vec{k}} | r_k, \phi_k, \theta_k \rangle \sim e^{A(q_{\vec{k}}^2 + q_{-\vec{k}}^2) - B q_{\vec{k}} q_{-\vec{k}}}$$

What is the long-term behavior of cosmological complexity?

Inside Horizon:

- $r_k \ll 1$
- Gaussian State Complexity insensitive to phase
- $\mathcal{C}_{\text{cosmo}}^{(s)} \ll 1$



Outside Horizon:

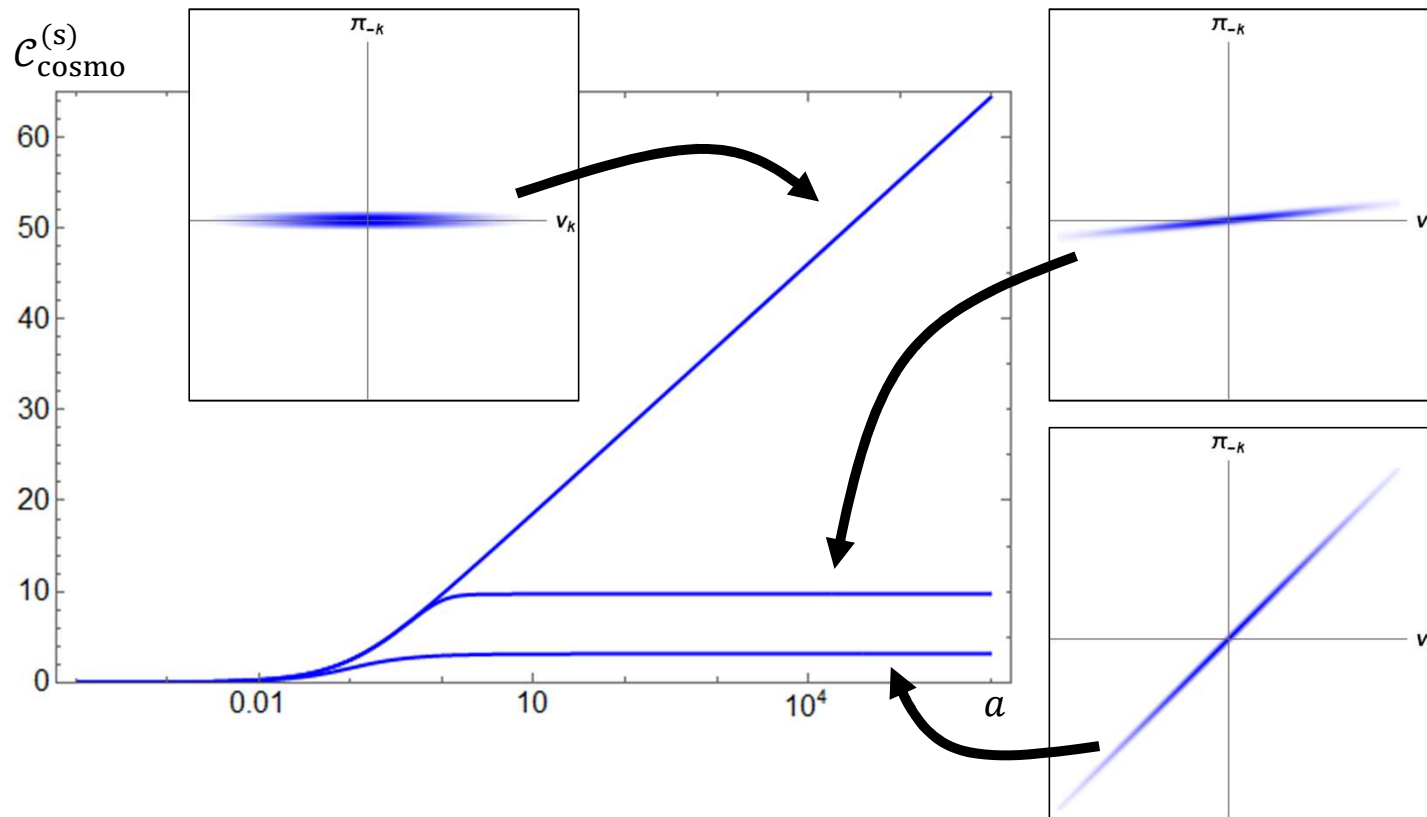
- Grows with squeezing, e-folds
- $\mathcal{C}_{\text{cosmo}}^{(s)} \approx r_k \approx \ln\left(\frac{a}{a_{\text{exit}}}\right)$
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Cosmological Complexity

the complexity of quantum cosmological perturbations

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Cosmological
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$$|r_k, \phi_k, \theta_k\rangle = \hat{S}(r_k, \phi_k) \hat{R}(\theta_k) |0_{\vec{k}}, 0_{-\vec{k}}\rangle$$

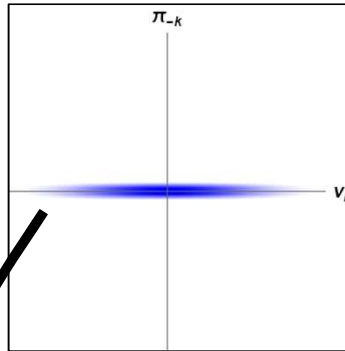


Unbounded growth of complexity depends sensitively on squeezing angle ϕ

- Complexity of dS is **maximal** w.r.t. ϕ
Why?

Cosmological Complexity

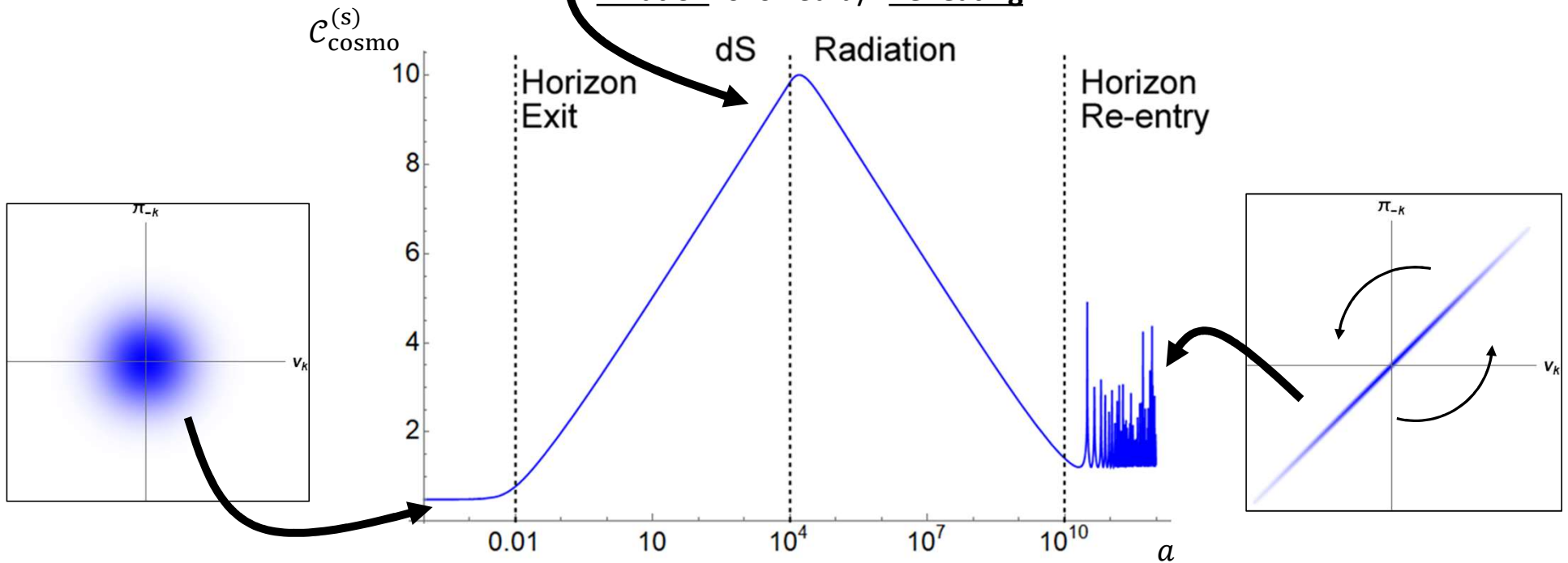
the complexity of quantum cosmological perturbations



“De-Complexification”:

- Complexity **decreases** for radiation-dominated, then “freezes-in” upon horizon re-entry!
- Modes are still highly squeezed
- dS \rightarrow radiation transition cuts off complexity growth

Inflation followed by “**Reheating**”



Bhattacharyya, Das, Haque, BU

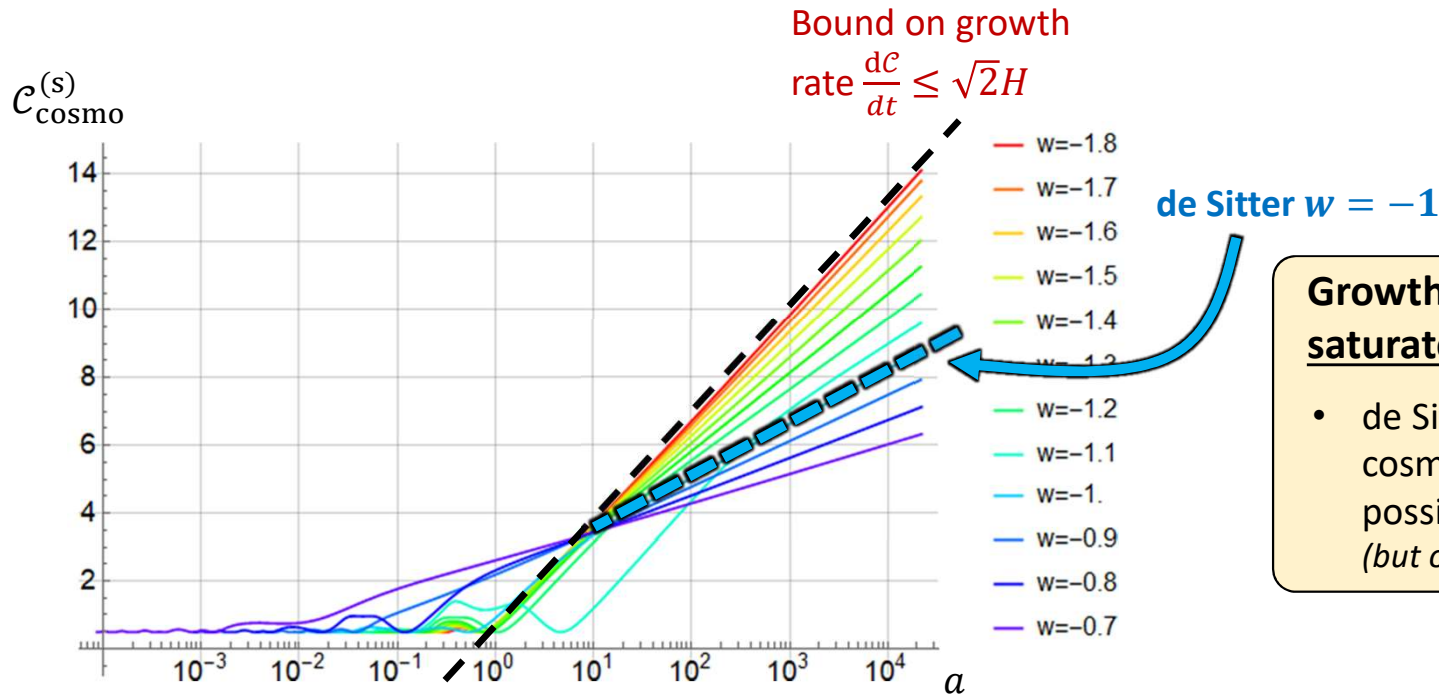
Cosmological Complexity

the complexity of quantum cosmological perturbations

Accelerating, Expanding Backgrounds

$$ds^2 = a(\eta)^2 (-d\eta^2 + d\vec{x}^2) \quad a(\eta) = \left(\frac{\eta_0}{\eta}\right)^{-2/(1+3w)}$$

Equation of state $p = w\rho$



Growth rate of complexity saturates at $w = -5/3$

- de Sitter is not fastest growth in cosmological complexity among all possible accelerating backgrounds... *(but others violate NEC)*

Bhattacharyya, Das, Haque, BU

Cosmological Complexity

Decoherence

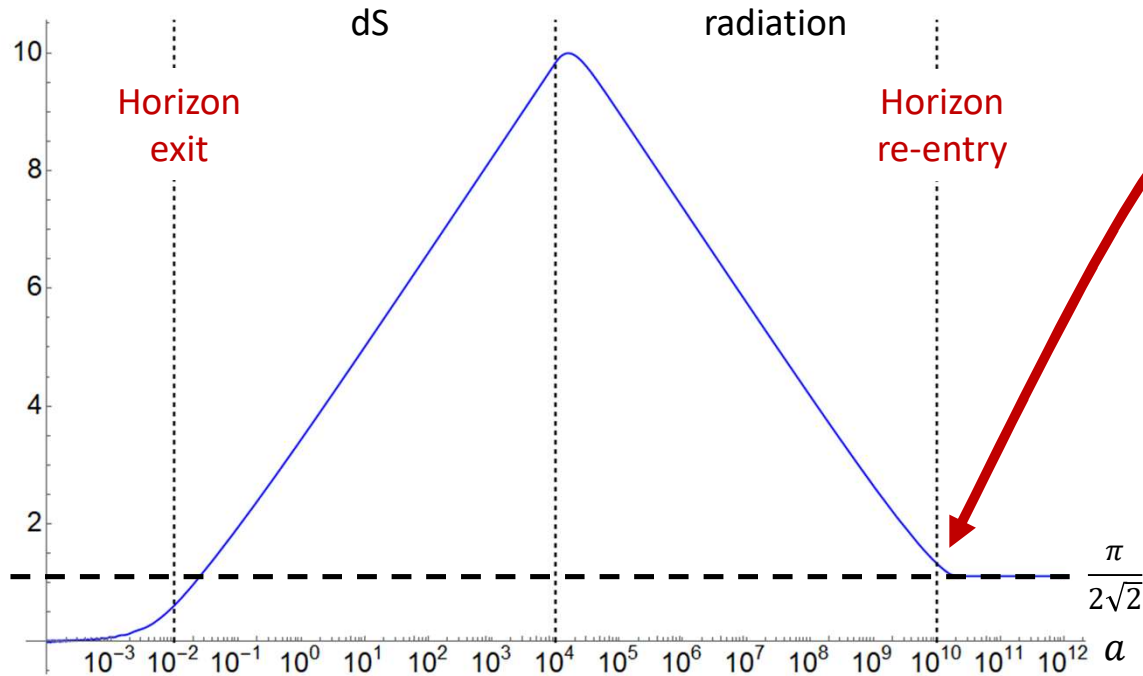
Pure State

$$\hat{\rho}_{\text{pure}} = |r_k, \phi_k, \theta_k\rangle\langle r_k, \phi_k, \theta_k|$$

Thermal Density Matrix

$$\hat{\rho}_{\text{red}} = \frac{1}{\cosh^2 r_k} \sum_{n=0}^{\infty} \tanh^{2n} r_k |n_{\vec{k}}, n_{-\vec{k}}\rangle\langle n_{\vec{k}}, n_{-\vec{k}}|$$

$\mathcal{C}_{\text{purif}}$



Complexity of Purification

- Assume decoherence occurs at re-entry
- Purification with ancillary dof
 $\mathcal{H} \rightarrow \mathcal{H} \otimes \mathcal{H}_{\text{anc}}$
- Minimize complexity over purification

$$\mathcal{C}_{\text{purif}} = \min_{\{\text{anc}\}} \mathcal{C}_{\text{tot}}$$

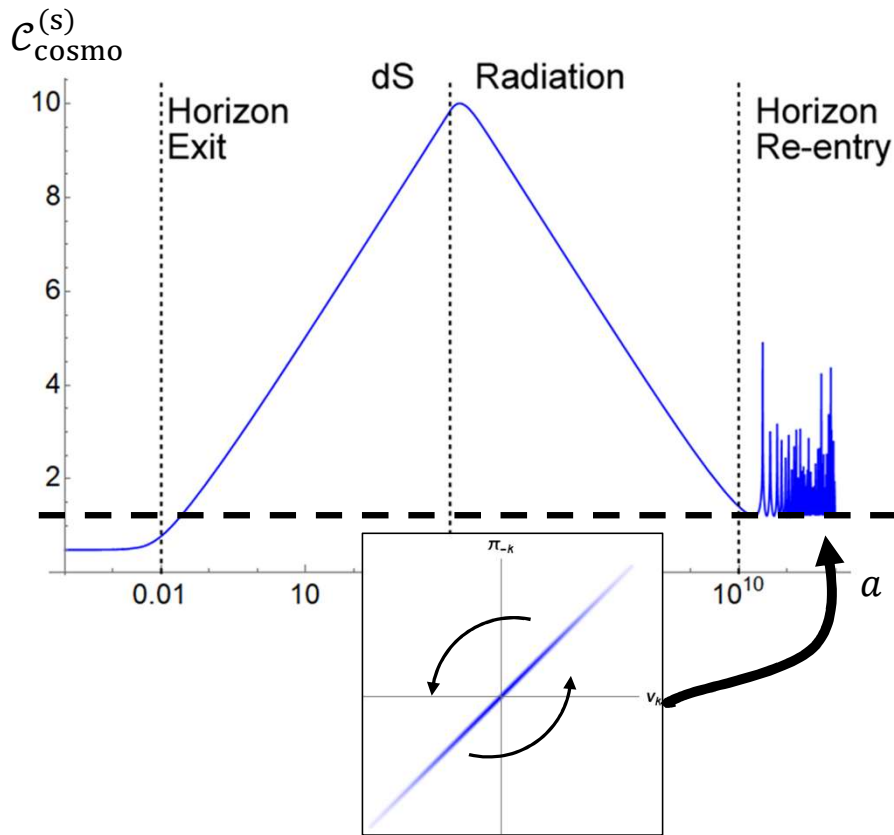
- Complexity of purification $\mathcal{O}(1)$

$$\mathcal{C}_{\text{purif}} \approx \frac{\pi}{2\sqrt{2}}$$

Cosmological Complexity

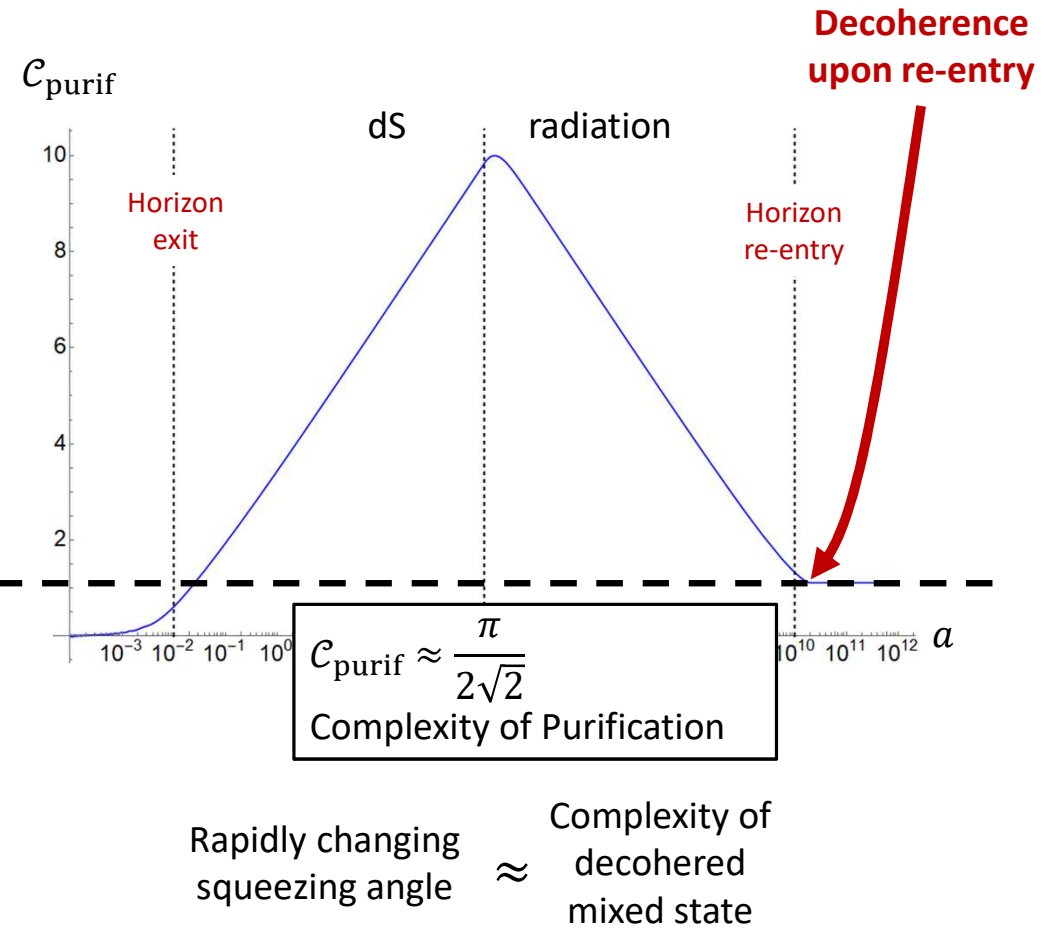
Decomplexification

Inflation followed by “Reheating”



vs

Decoherence

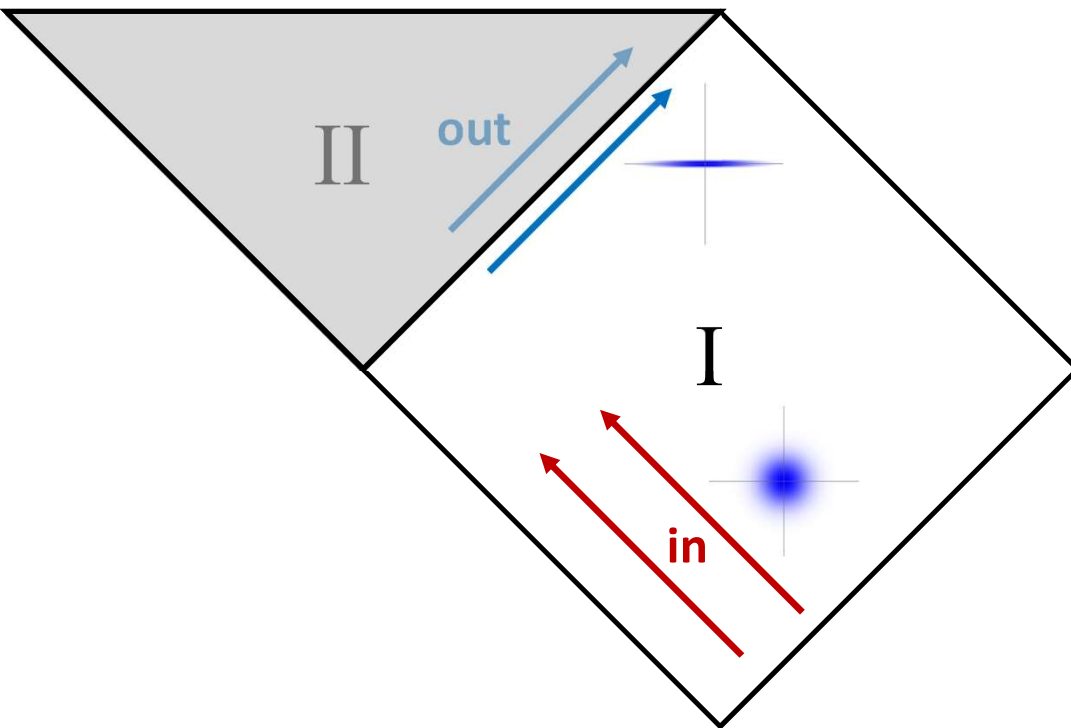


Aside: Complexity of Hawking Radiation

Hawking radiation: two-mode squeezed states

$$|0\rangle_{\text{in}} = \mathcal{N}_k \sum_{n_k} e^{-(4\pi GMk) n_k} |n_k\rangle_{\text{I}} \otimes |n_k\rangle_{\text{II}}$$

$\tanh r_k = e^{-4\pi GM}$ $\begin{cases} r_k \ll 1 & k \gg (GM)^{-1} \text{ high freq} \\ r_k \gg 1 & k \ll (GM)^{-1} \text{ low freq} \end{cases}$
 $\phi_k = \pi/2$



- Complexity of Hawking Radiation

$$\mathcal{C}_{\text{Hawk}}(k) = \frac{1}{\sqrt{2}} \ln \left(\frac{1 + e^{-4\pi GMk}}{1 - e^{-4\pi GMk}} \right)$$

$$\approx \begin{cases} 0 & k \gg (GM)^{-1} \text{ high freq} \\ \ln \left(\frac{1}{GMk} \right) & k \ll (GM)^{-1} \text{ low freq} \end{cases}$$

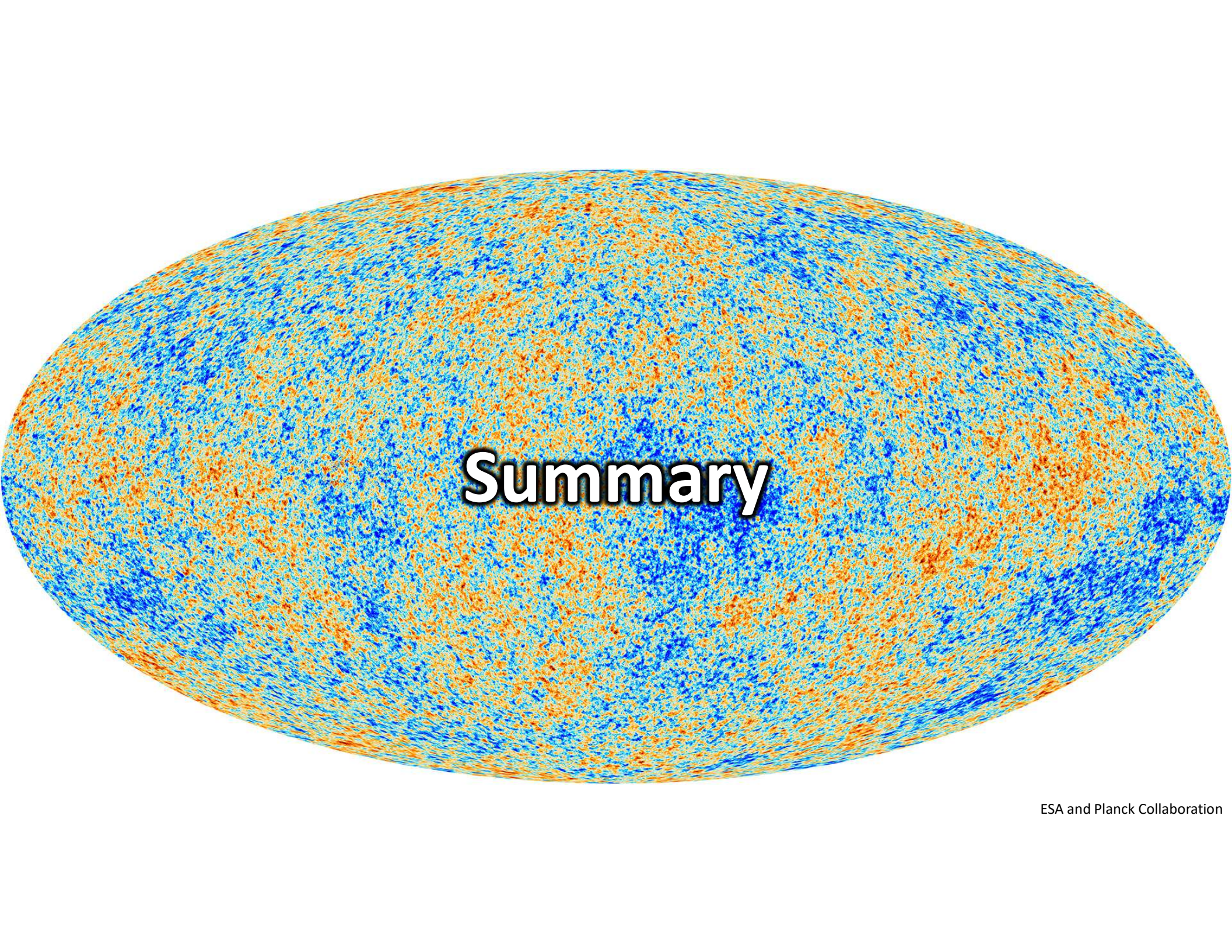
- Complexity of Hawking radiation is maximal w.r.t. ϕ
Why?

- Tracing out modes inside horizon

- Thermal density matrix
- Complexity of purification

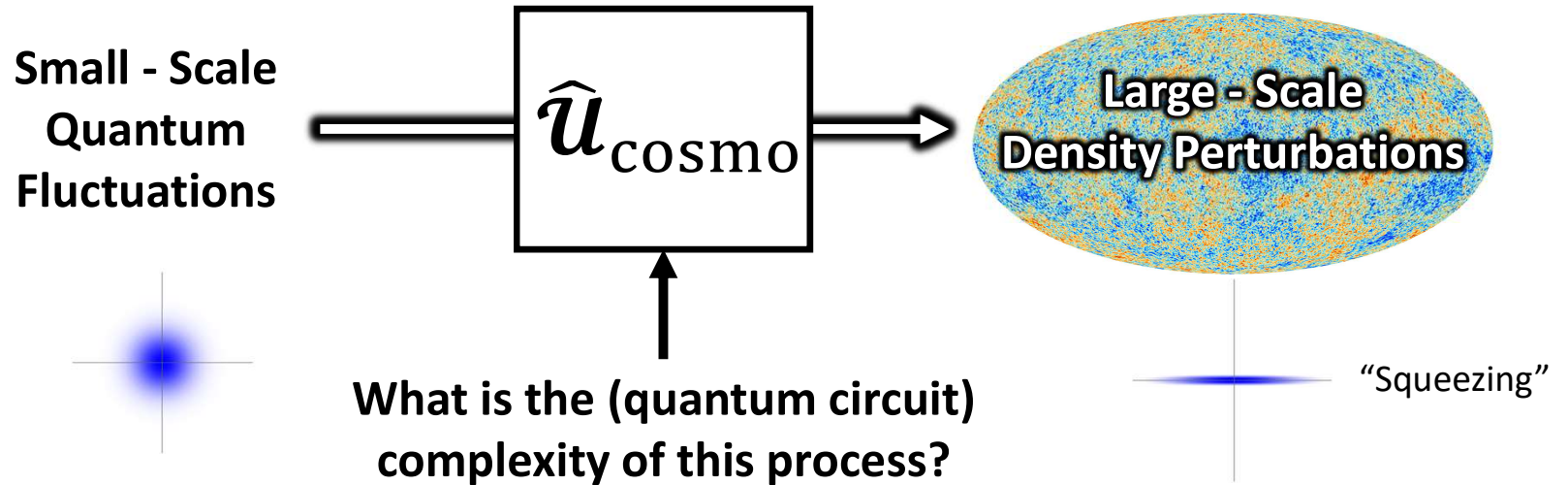
$$\mathcal{C}_{\text{purif}}(k) \approx \frac{\pi}{2\sqrt{2}}$$

- Complexity of reduced state smaller than pure state



Summary

Summary



- ❑ Cosmological Complexity in dS grows linearly with time $\mathcal{C}_{\text{cosmo}} = H_{dS} t$
- ❑ Complexity depends sensitively on squeezing angle ϕ
 - Complexity of dS is **maximal** w.r.t. ϕ . Why?
- ❑ Growth rate of complexity $\frac{d\mathcal{C}}{dt}$ is bounded from above for accelerating backgrounds
- ❑ **Decomplexification** during radiation-domination phase
 - Connection between decomplexification and decoherence?