

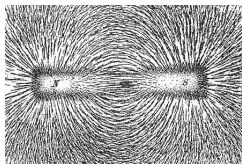
# Expanding Blaschke Products for the Lee-Yang zeros on the Diamond Hierarchical Lattice.

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Banff

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$$H_n(\sigma) = -J \cdot I(\sigma) - h \cdot M_n(\sigma),$$

where  $J > 0$ .

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An obvious danger occurs at those values of  $h, T$  for which  $Z_n(h, T) = 0$ . Luckily, this never happens for  $h, T \in \mathbb{R}$ .

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Fundamental symmetry of the Ising model!

Thermodynamic quantities in terms of zeros of  $Z_n(z, t)$ .

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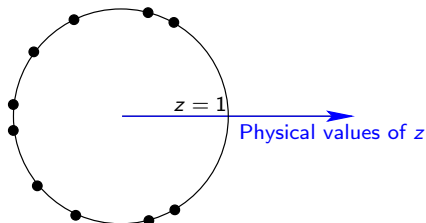
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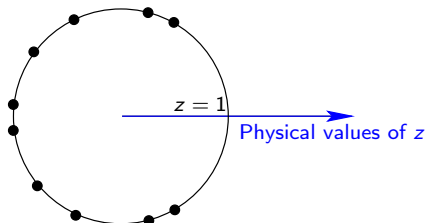


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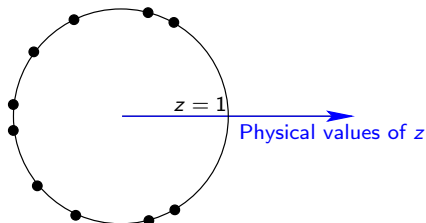
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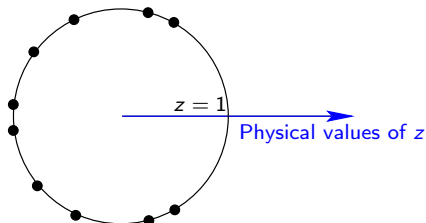
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For each  $t \in [0, 1]$  there is a measure  $\mu_t$  on  $\mathbb{T}$  describing the asymptotic distribution of Lee-Yang zeros.

## Phase transitions in terms of Lee-Yang distribution

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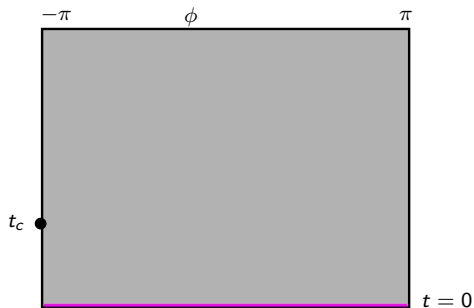
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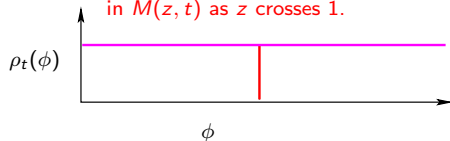
Understanding how the Lee-Yang distributions  $\mu_t(\phi)$  vary with  $t$  and  $\phi$  is essential to understanding phase transitions of the model.



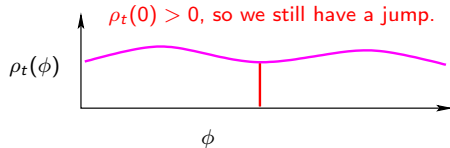
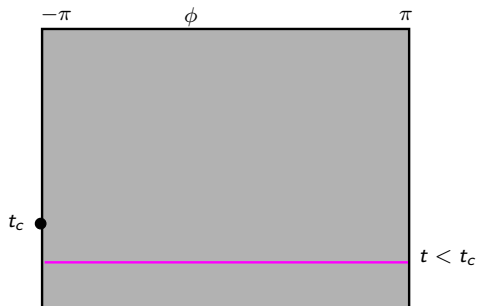
# Expected limiting distributions of Lee-Yang zeros for $\mathbb{Z}^2$



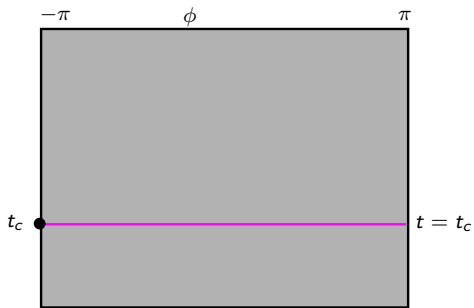
$\rho_t(0) > 0$  gives jump discontinuity  
in  $M(z, t)$  as  $z$  crosses 1.



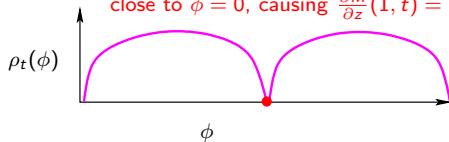
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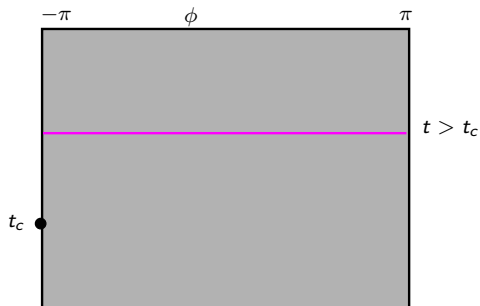
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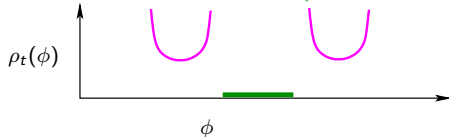
Now,  $\rho_t(0) = 0$ , so no jump.  
However, they accumulate arbitrarily close to  $\phi = 0$ , causing  $\frac{\partial M}{\partial z}(1, t) = \infty$



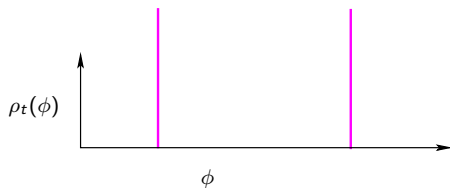
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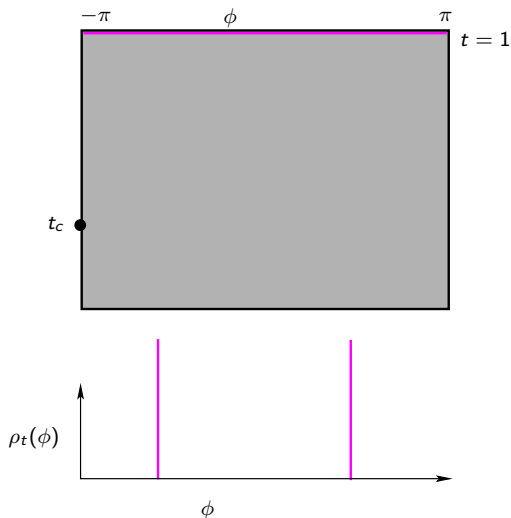
Now, we have a nice interval around  $\phi = 0$  with  $\rho_t(\phi) \equiv 0$ . Causes  $M(z, t)$  to be differentiable at  $z = 1$  (and hence everywhere).



# Expected limiting distributions of Lee-Yang zeros for $\mathbb{Z}^2$

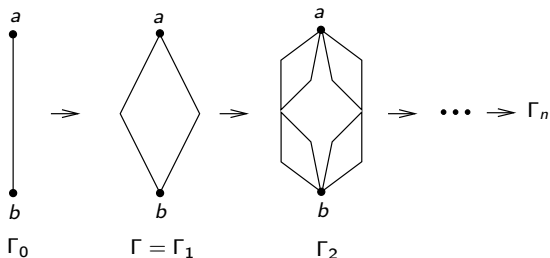


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Zakhar Kabluchko: Lee-Yang zeros for the Curie Weiss model match this conjectural description. <http://arxiv.org/pdf/2203.05533.pdf>

# Hierarchical Lattices



The Diamond Hierarchical Lattice (DHL).

$\Gamma_n$  is obtained by replacing each edge of generating graph  $\Gamma$  (a diamond) with a copy of  $\Gamma_{n-1}$ , considering the marked vertices  $a$  and  $b$  as the “endpoints” of  $\Gamma_{n-1}$ .

# Migdal-Kadanoff Renormalization<sup>123</sup>

Consider the conditional partition functions:

$$U_n := Z_n \left( \begin{array}{c} \oplus \\ \text{wavy circle} \\ \oplus \end{array} \right), \quad V_n := Z_n \left( \begin{array}{c} \oplus \\ \text{wavy circle} \\ \ominus \end{array} \right) = Z_n \left( \begin{array}{c} \ominus \\ \text{wavy circle} \\ \oplus \end{array} \right), \quad W_n := Z_n \left( \begin{array}{c} \ominus \\ \text{wavy circle} \\ \ominus \end{array} \right)$$

The total partition function is equal to  $Z_n = U_n + 2V_n + W_n$ .

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Migdal-Kadanoff RG Equations:

$$U_{n+1} = (U_n^2 + V_n^2)^2, \quad V_{n+1} = V_n^2(U_n + W_n)^2, \quad W_{n+1} = (V_n^2 + W_n^2)^2.$$

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## MK renormalization in the $(z, t)$ coordinates:

We can lift  $R$  from the  $[U : V : W]$  coordinates (downstairs) to the  $[z : t : 1]$  coordinates upstairs:

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and  $\Psi$  is some degree 2 rational map.

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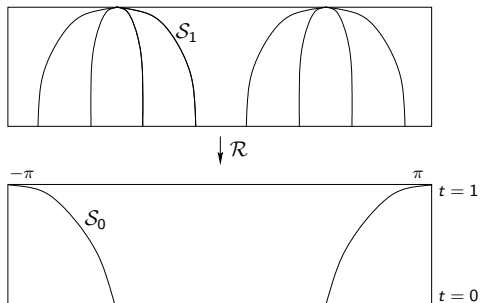
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It is this recursive relationship between  $\mathcal{S}_{n+1}$  and  $\mathcal{S}_n$  that makes a study of the Lee-Yang zeros tractable for hierarchical lattices.

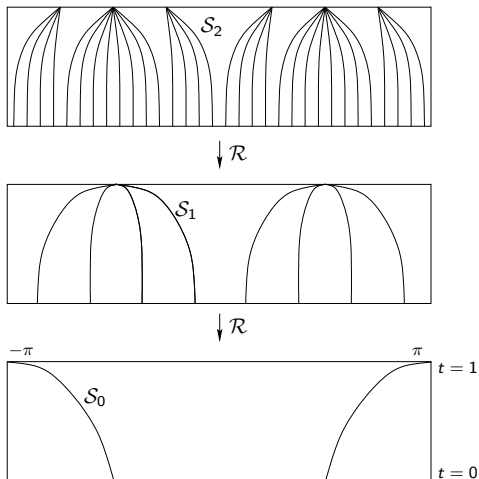
# Lee-Yang zeros as pull-backs under $\mathcal{R}$



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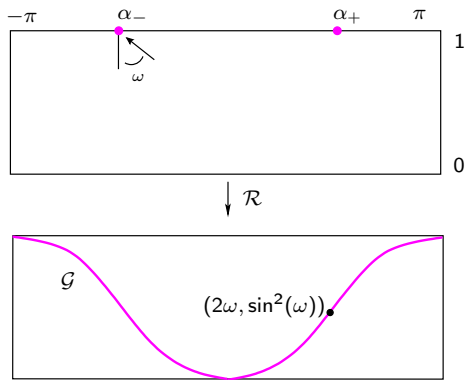
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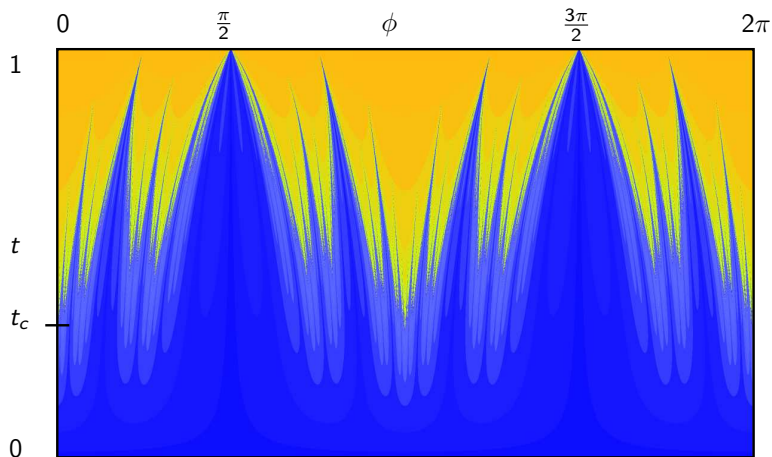
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# Numerical Experiment



$W^s(B)$  is colored blue and  $W^s(T)$  is colored orange.

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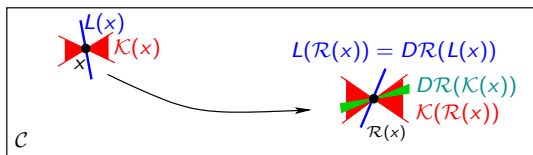
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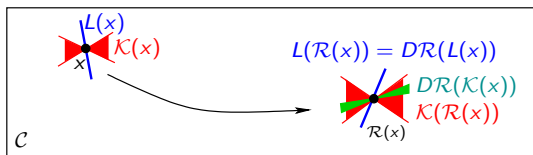
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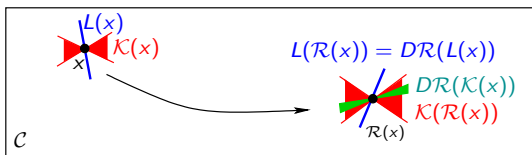
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The idea of this proof that this cone field is invariant seems to play a role in the recent work of Dang-Grigorchuk-Lyubich about the Basilica IMG.



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This is the “intertwined basins” phenomenon studied by Kan-Yorke, Bonifant-Milnor, Ilyashenko-Kleptsyn-Saltykov....

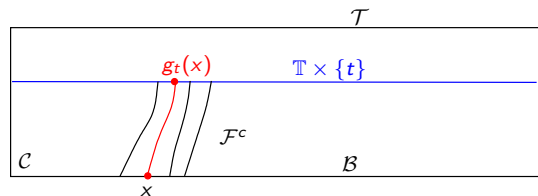
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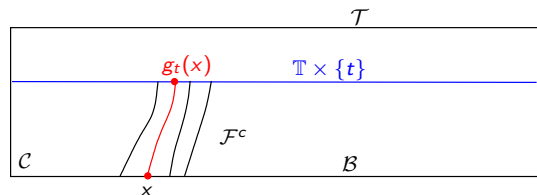
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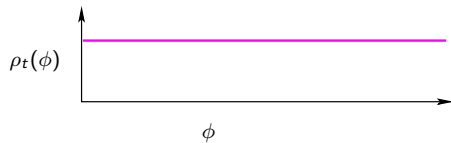
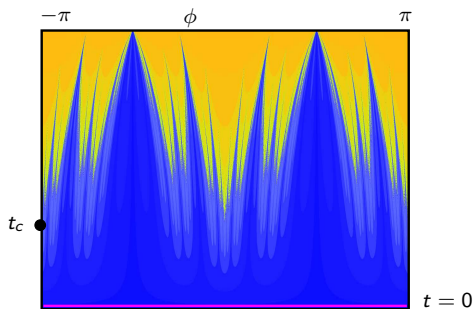
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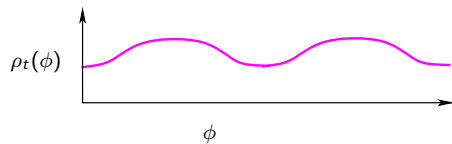
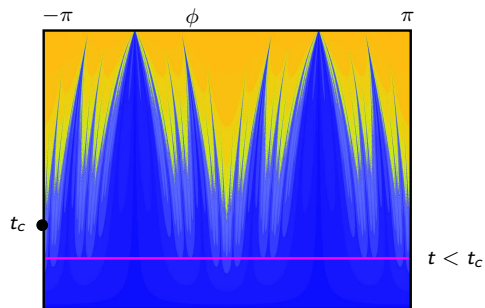
## Theorem (BLR)

The asymptotic distribution of Lee-Yang zeros at a temperature  $t_0 \in [0, 1)$  is given by under **holonomy** by  $\mu_t = (g_t)_*(\mu_0)$  where  $\mu_0$  be the Lebesgue measure on  $\mathcal{B}$ .

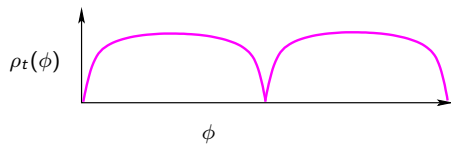
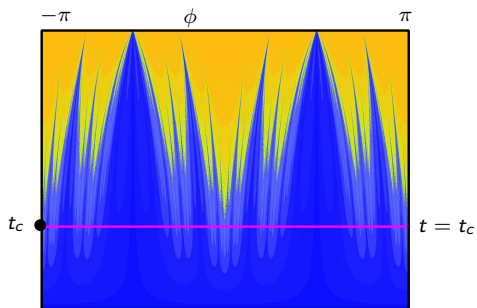
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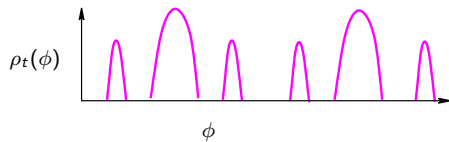
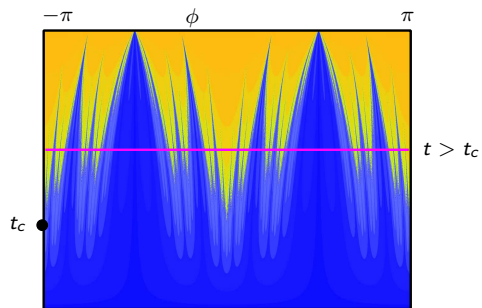
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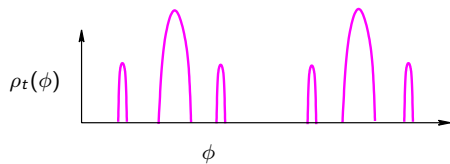
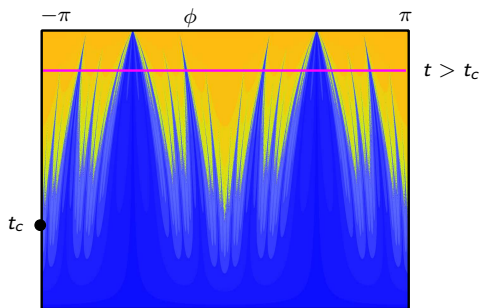
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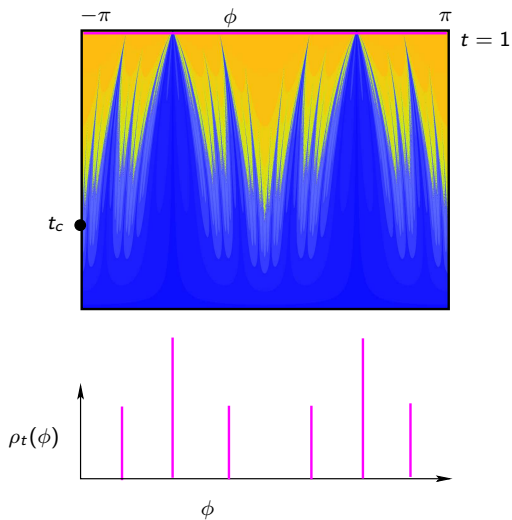
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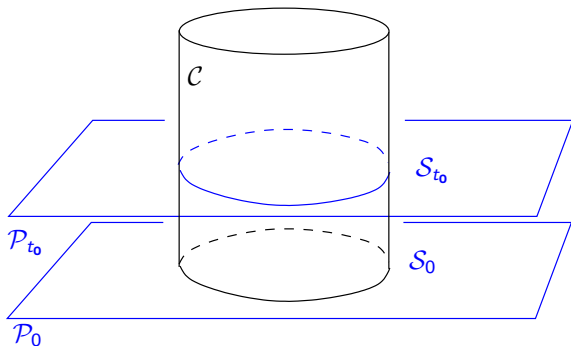


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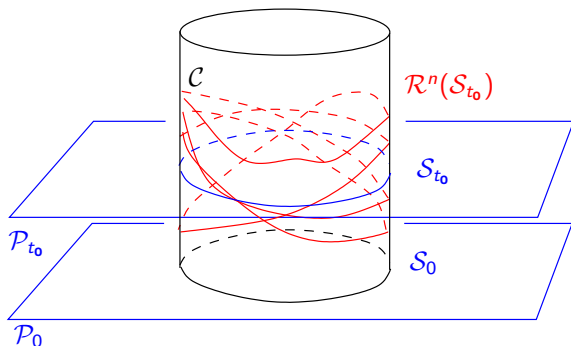
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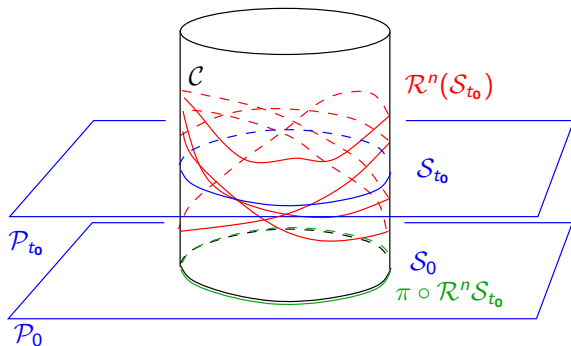
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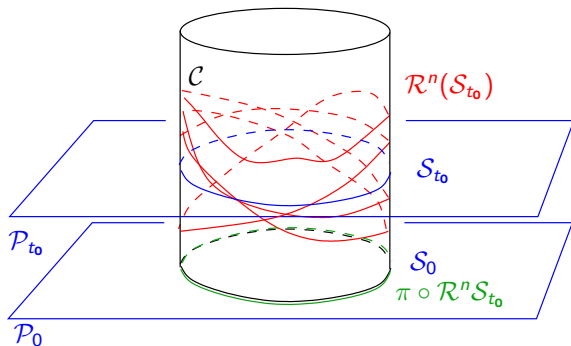
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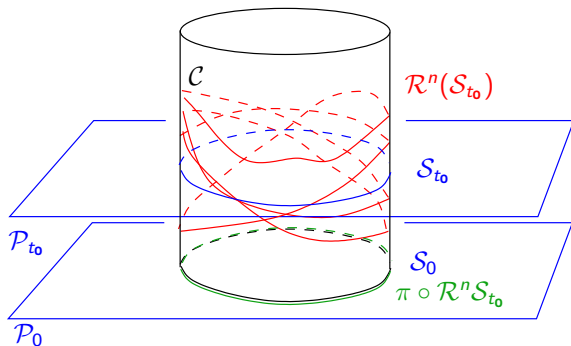
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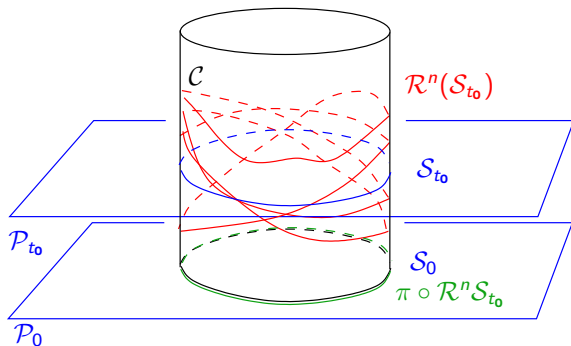


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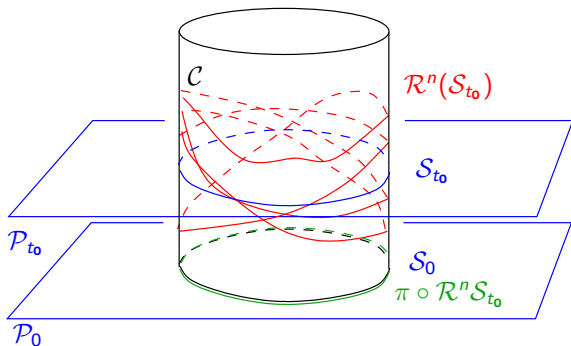
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Algebraic instability:  $4^n < \deg(\mathcal{R}^n) < (\deg(\mathcal{R}))^n = 6^n$ .



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Original idea actually works in these coordinates!

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**Claim:**  $\psi_n : \mathbb{C} \rightarrow \mathbb{C}$  is an Blaschke product preserving the unit disc  $\mathbb{D}$ , expanding the circle  $\mathbb{T} = \partial\mathbb{D}$  by a factor of  $2^{n+1}$ .

## Conditional partition functions and their symmetries

$$\begin{aligned} U_n(z, t) &= \sum_{\sigma(a)=\sigma(b)=+1} W(\sigma) = \sum_{\sigma(a)=\sigma(b)=+1} t^{-I(\sigma)} z^{-M(\sigma)} \\ &= a_d^+(t) z^d + \cdots + a_{-d}^+(t) z^{-d}, \end{aligned}$$

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$$a_i^+(t) = a_{-i}^-(t) \quad \text{for each } i = -d \dots d$$

2. Since  $\Gamma_n$  has valence  $2^n$  at marked vertices  $a$  and  $b$  we have

$$a_i^-(t) = 0 \quad \text{for } i < -4^n + 2^{n+1}$$

Reason for 2: With  $-1$  spins at the marked vertices  $a, b$ , we can't get more than  $4^n - 2^{n+1}$  edges with  $++$ , so  $M(\sigma) \leq 4^n - 2^{n+1}$ .

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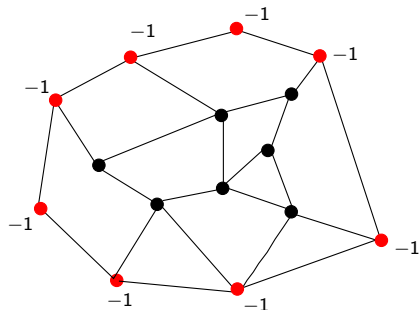
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so we'd be done!

# Lee-Yang Theorem with Boundary conditions



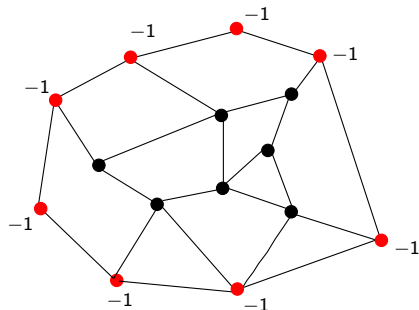
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Consider a ferromagnetic Ising model on a connected graph  $\Gamma$  and let  $\sigma_S \equiv -1$  on a nonempty subset  $S$  of the vertex set  $V$ .



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## Theorem (Bleher, Lyubich, R)

Consider a ferromagnetic Ising model on a connected graph  $\Gamma$  and let  $\sigma_S \equiv -1$  on a nonempty subset  $S$  of the vertex set  $V$ .

Then, for any temperature  $t \in (0, 1)$  the Lee-Yang zeros  $z_i^-(t)$  of the conditional partition function  $Z_{\Gamma|\sigma_S}$  lie inside the open disc  $\mathbb{D}$ .

# Thank you for listening!

Pavel Bleher, Mikhail Lyubich, and Roland Roeder. *Lee-Yang Zeros for the DHL and 2D Rational Dynamics, I. Foliation of the Physical Cylinder*. Journal de Mathématiques Pures et Appliquées, 107(5): 491-590, 2017.

For those who like joint spectra:

Pavel Bleher, Mikhail Lyubich, and Roland Roeder. *Lee-Yang-Fisher zeros for DHL and 2D rational dynamics, II. Global Pluripotential Interpretation*. Journal of Geometric Analysis, 30(1): 777-833, 2020.