

A NO-REPLICA TRICK FOR THE FREE ENERGY

Sergio Hernández-Cuenca

based on w.i.p. with Ven Chandrasekaran, Netta Engelhardt and Sebastian Fischetti

> Gravitational Emergence in AdS/CFT BIRS

> > October 28, 2021

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INTRODUCTION

THE GENERAL PRESCRIPTION

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DISCUSSION

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► The Euclidean GPI is a mysterious object in quantum gravity

$$\mathcal{P}(B) = \int_{\partial M = B} \mathcal{D}g \ e^{-S[g]}$$

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- ► Inclusion of connected topologies has crucial consequences for:
 - von Neumann entropies (unitary Page curve)
 - gravitational correlators (factorization problem $\overline{Z(B^m)} \neq \overline{Z(B)}^m$)
 - free energies [Engelhardt-Fischetti-Maloney]

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 - free energies [Engelhardt-Fischetti-Maloney]
- Annealed vs quenched free energies:

$$F_a = -\frac{1}{\beta}\log\overline{Z}$$
 vs $F_q = -\frac{1}{\beta}\overline{\log Z}$

• How can we use \mathcal{P} to compute $\overline{\log Z}$?

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► Not so fast: continuation to no replicas is ill-defined

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THE GENERAL PRESCRIPTION

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• Constrain $\mathcal{P}(B^m)$ to replica symmetric manifolds M_m for $m \in \mathbb{Z}_+$

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- ► Work in quotient $\hat{M}_m = M_m / \mathbb{Z}_m$ for a unique extension to $m \in \mathbb{R}_{\geq 0}$ [Lewkowycz-Maldacena]

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- Constrain $\mathcal{P}(B^m)$ to replica symmetric manifolds M_m for $m \in \mathbb{Z}_+$
- ► Work in quotient $\hat{M}_m = M_m / \mathbb{Z}_m$ for a unique extension to $m \in \mathbb{R}_{\geq 0}$ [Lewkowycz-Maldacena]
- Localize the path integral to gravitational saddle points:

$$\overline{\log Z} = \lim_{m \to 0} \frac{1}{m} (e^{-I[M_m]} - 1) = \lim_{m \to 0} \frac{1}{m} (e^{-mI[\hat{M}_m]} - 1) = -I[\hat{M}_0]$$

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TWO INTERESTING OBSERVATIONS

1. Not just the LM recipe...

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Chandrasekaran, HC, Engelhardt, Fischetti

...bring the CHEF recipe!

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2. Saddle points in the $m \rightarrow 0$ limit give quenched generating functionals in quantum gravity... who are these creatures?

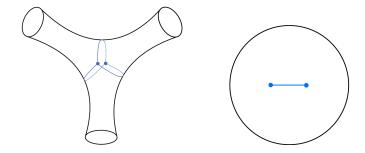
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A Story about JT

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The Quotient Geometry

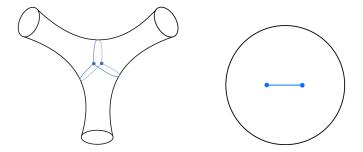
Example: replica wormhole M_3 and \mathbb{Z}_3 orbifold \hat{M}_3 [East Coast]



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THE QUOTIENT GEOMETRY

Example: replica wormhole M_3 and \mathbb{Z}_3 orbifold \hat{M}_3 [East Coast]

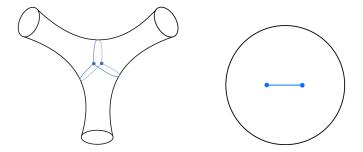


In general, M̂_m is conformal to a Poincaré disk with two conical defects of opening angle 2π/m

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THE QUOTIENT GEOMETRY

Example: replica wormhole M_3 and \mathbb{Z}_3 orbifold \hat{M}_3 [East Coast]



- ► In general, \hat{M}_m is conformal to a Poincaré disk with two conical defects of opening angle $2\pi/m$
- Wormhole throat sizes relate to proper distance between defects

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JT AND BOUNDARY CONDITIONS

► JT action:

$$I = -\frac{S_0}{4\pi} \left[\int_M R + 2 \int_{\partial M} K \right] - \frac{1}{2} \int_M \Phi(R+2) - \int_{\partial M} \Phi(K-1)$$

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• Gauss-Bonnet and Φ path integral:

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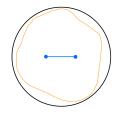
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Boundary conditions:

Cutoff boundaries identified with level sets of the dilaton $\Phi|_{\partial M} = 1/\delta$. Limit $\delta \to 0$ taken with fixed ratio $L_{\partial M}/\Phi|_{\partial M} = \beta$



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AN INTERACTING SCHWARZIAN THEORY

Near-boundary metric:

$$g = \left(\frac{1}{(1-\xi)^2} + h_a^{(m)}(\phi) + \mathcal{O}(1-\xi)\right) \left(d\xi^2 + d\phi^2\right)$$

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• Wiggle $\phi : \mathbb{S}_{\beta} \to \partial M \cong \mathbb{S}$ defined by $g|_{\partial M} = du^2/\delta^2$

$$-\int_{\partial M} \Phi(K-1) = \int_{\mathbb{S}_{\beta}} du \left(\left\{\tan\frac{\phi}{2}, u\right\} + \frac{1}{2}(1+3h_a^{(m)}(\phi)) \phi'(u)^2\right)$$

Note $h_a^{(m)}(\phi) = -\frac{1}{3}$ for m = 1,2 leave the Schwarzian alone

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Wiggle equation of motion:

$$\left(\frac{1}{\phi'}\left(\frac{\phi''}{\phi'}\right)' - 3h_a^{(m)}(\phi)\phi'\right)' + \frac{3}{2}\left(h_a^{(m)}(\phi)\right)'\phi' = 0$$

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THE JT HELLSCAPE

• Explicit JT wiggle action for all replica $m \in \mathbb{R}_{\geq 0}$ and moduli *a*:

$$I_a^{(m)}(eta) = rac{8}{eta} \operatorname{arcsinh}^2 \sqrt{\sin^2 rac{\pi}{m} \cosh^2 rac{a}{2} - 1}$$

Note $I_a^{(1)}(\beta) = -2\pi^2/\beta$, $I_a^{(2)}(\beta) = 2a^2/\beta$

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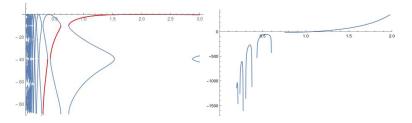
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Action and stability analysis...



... and the modulus is not stabilized

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ADDING MATTER

Conformal matter with classical sources:

$$I_{\rm mat} = \frac{1}{2} \int_M (d\psi)^2.$$

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▶ Solve in terms of boundary profile specified by $\psi_0: \partial M \to \mathbb{R}$

$$I_{\text{mat}} = \frac{1}{2} \int_{\partial M} \psi \nabla_n \psi = \frac{1}{2} \int_{\mathbb{S}} d\phi \, d\tilde{\phi} \, \psi_0(\phi) S(\phi, \tilde{\phi}) \psi_0(\tilde{\phi}),$$

where S is an integral kernel known explicitly but gross

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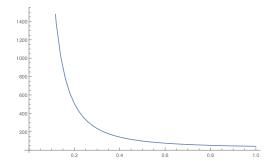
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A SILVER LINING

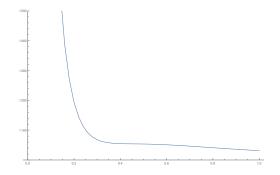
• Modulus saddles appear for m = 2 as sources are turned on:



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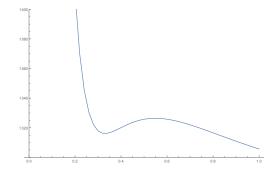
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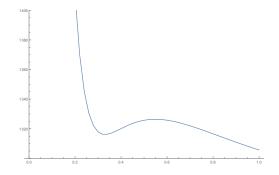
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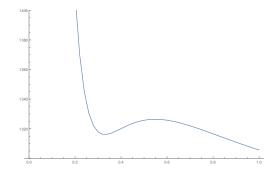


▶ Will modulus saddles make it all the way to *m* < 1?

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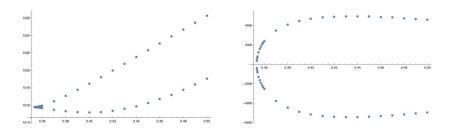
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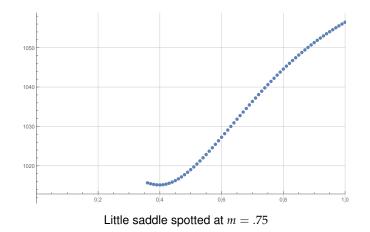
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THE LITTLE SADDLE THAT COULD

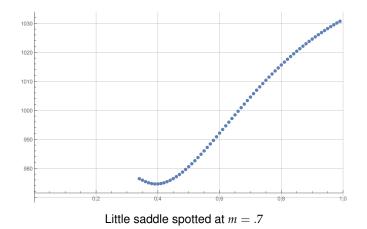
- ▶ Pair of stable/unstable branches of solutions exist for m < 1!
- ► The little wormhole can be made to dominate over the disk one
- Action and stability analysis for m = .75:



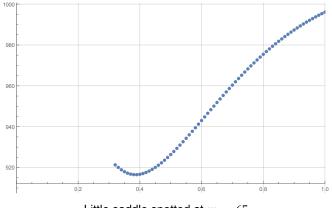
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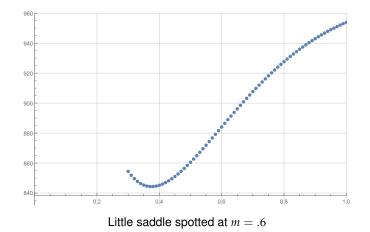


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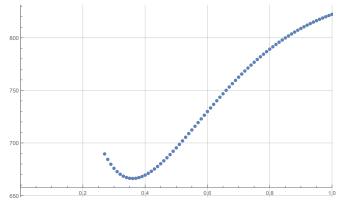


Little saddle spotted at m = .65

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DISCUSSION

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Outlook

- Will the little saddle make it to $m \rightarrow 0$?
- What properties does the resulting generating functional have?
- How does it differ from the annealed result?
- What is the effect on scalar correlation functions?

OPEN QUESTIONS AND FUTURE DIRECTIONS

- ► Is there a simple diagnostic for when quenched $m \rightarrow 0$ saddle points will differ from annealed ones?
- Is there any correlation between dominance of replica wormholes for *m* ∈ Z₊ and for 0 < *m* < 1?</p>
- Are there any universal features about $m \rightarrow 0$ saddle points and quenched generating functionals in quantum gravity?
- Other toy models for the study $m \rightarrow 0$ saddles?

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Thank you for listening!