## UC SANTA BARBARA

## A No-Replica Trick for the Free Energy

Sergio Hernández-Cuenca<br>based on w.i.p. with<br>Ven Chandrasekaran, Netta Engelhardt and Sebastian Fischetti

Gravitational Emergence in AdS/CFT
BIRS
October 28, 2021

# INTRODUCTION 

# The General Prescription 

A Story about JT

DISCUSSION

## INTRODUCTION

## BACKGROUND

- The Euclidean GPI is a mysterious object in quantum gravity

$$
\mathcal{P}(B)=\int_{\partial M=B} \mathcal{D} g e^{-S[g]}
$$

## BACKGROUND

- The Euclidean GPI is a mysterious object in quantum gravity

$$
\mathcal{P}(B)=\int_{\partial M=B} \mathcal{D} g e^{-S[g]}
$$

- Generally assumed to compute some sort of partition function:

$$
\overline{Z(B)}=\mathcal{P}(B)
$$

## BACKGROUND

- The Euclidean GPI is a mysterious object in quantum gravity

$$
\mathcal{P}(B)=\int_{\partial M=B} \mathcal{D} g e^{-S[g]}
$$

- Generally assumed to compute some sort of partition function:

$$
\overline{Z(B)}=\mathcal{P}(B)
$$

- Inclusion of connected topologies has crucial consequences for:
- von Neumann entropies (unitary Page curve)
- gravitational correlators (factorization problem $\overline{Z\left(B^{m}\right)} \neq \overline{Z(B)^{m}}$ )
- free energies


## BACKGROUND

- The Euclidean GPI is a mysterious object in quantum gravity

$$
\mathcal{P}(B)=\int_{\partial M=B} \mathcal{D} g e^{-S[g]}
$$

- Generally assumed to compute some sort of partition function:

$$
\overline{Z(B)}=\mathcal{P}(B)
$$

- Inclusion of connected topologies has crucial consequences for:
- von Neumann entropies (unitary Page curve)
- gravitational correlators (factorization problem $\overline{Z\left(B^{m}\right)} \neq \overline{Z(B)^{m}}$ )
- free energies
[Engelhardt-Fischetti-Maloney]
- Annealed vs quenched free energies:

$$
F_{a}=-\frac{1}{\beta} \log \bar{Z} \quad \text { vs } \quad F_{q}=-\frac{1}{\beta} \overline{\log Z}
$$

## Towards a No-Replica Trick

- How can we use $\mathcal{P}$ to compute $\overline{\log Z}$ ?


## Towards a No-Replica Trick

- How can we use $\mathcal{P}$ to compute $\overline{\log Z}$ ?
- Mathematical identity:

$$
\log Z=\lim _{m \rightarrow 0} \frac{1}{m}\left(Z^{m}-1\right)
$$

## Towards a No-Replica Trick

- How can we use $\mathcal{P}$ to compute $\overline{\log Z}$ ?
- Mathematical identity:

$$
\log Z=\lim _{m \rightarrow 0} \frac{1}{m}\left(Z^{m}-1\right)
$$

- Put overlines on both sides and we are done, right?

$$
\overline{\log Z}=\lim _{m \rightarrow 0} \frac{1}{m}\left(\mathcal{P}\left(B^{m}\right)-1\right)
$$

## Towards a No-Replica Trick

- How can we use $\mathcal{P}$ to compute $\overline{\log Z}$ ?
- Mathematical identity:

$$
\log Z=\lim _{m \rightarrow 0} \frac{1}{m}\left(Z^{m}-1\right)
$$

- Put overlines on both sides and we are done, right?

$$
\overline{\log Z}=\lim _{m \rightarrow 0} \frac{1}{m}\left(\mathcal{P}\left(B^{m}\right)-1\right)
$$

- Not so fast: continuation to no replicas is ill-defined


## The General Prescription

## An LM-Style No-Replica Trick

- Assume higher topologies are parametrically suppressed


## An LM-Style No-Replica Trick

- Assume higher topologies are parametrically suppressed
- Constrain $\mathcal{P}\left(B^{m}\right)$ to replica symmetric manifolds $M_{m}$ for $m \in \mathbb{Z}_{+}$


## An LM-Style No-Replica Trick

- Assume higher topologies are parametrically suppressed
- Constrain $\mathcal{P}\left(B^{m}\right)$ to replica symmetric manifolds $M_{m}$ for $m \in \mathbb{Z}_{+}$
- Work in quotient $\hat{M}_{m}=M_{m} / \mathbb{Z}_{m}$ for a unique extension to $m \in \mathbb{R}_{\geq 0}$


## An LM-Style No-Replica Trick

- Assume higher topologies are parametrically suppressed
- Constrain $\mathcal{P}\left(B^{m}\right)$ to replica symmetric manifolds $M_{m}$ for $m \in \mathbb{Z}_{+}$
- Work in quotient $\hat{M}_{m}=M_{m} / \mathbb{Z}_{m}$ for a unique extension to $m \in \mathbb{R}_{\geq 0}$
[Lewkowycz-Maldacena]
- Localize the path integral to gravitational saddle points:

$$
\overline{\log Z}=\lim _{m \rightarrow 0} \frac{1}{m}\left(e^{-I\left[M_{m}\right]}-1\right)=\lim _{m \rightarrow 0} \frac{1}{m}\left(e^{-m I\left[\hat{M}_{m}\right]}-1\right)=-I\left[\hat{M}_{0}\right]
$$

## Two Interesting Observations

1. Not just the LM recipe...

## Two Interesting Observations

1. Not just the LM recipe...

Chandrasekaran, HC, Engelhardt, Fischetti
...bring the CHEF recipe!

## Two Interesting Observations

1. Not just the LM recipe...

Chandrasekaran, HC, Engelhardt, Fischetti
...bring the CHEF recipe!
2. Saddle points in the $m \rightarrow 0$ limit give quenched generating functionals in quantum gravity... who are these creatures?

## A Story about JT

## The Quotient Geometry

- Example: replica wormhole $M_{3}$ and $\mathbb{Z}_{3}$ orbifold $\hat{M}_{3}$



## The Quotient Geometry

- Example: replica wormhole $M_{3}$ and $\mathbb{Z}_{3}$ orbifold $\hat{M}_{3}$

- In general, $\hat{M}_{m}$ is conformal to a Poincaré disk with two conical defects of opening angle $2 \pi / \mathrm{m}$


## The Quotient Geometry

- Example: replica wormhole $M_{3}$ and $\mathbb{Z}_{3}$ orbifold $\hat{M}_{3}$

- In general, $\hat{M}_{m}$ is conformal to a Poincaré disk with two conical defects of opening angle $2 \pi / \mathrm{m}$
- Wormhole throat sizes relate to proper distance between defects


## JT and Boundary Conditions

- JT action:

$$
I=-\frac{S_{0}}{4 \pi}\left[\int_{M} R+2 \int_{\partial M} K\right]-\frac{1}{2} \int_{M} \Phi(R+2)-\int_{\partial M} \Phi(K-1)
$$

## JT and Boundary Conditions

- JT action:

$$
I=-\frac{S_{0}}{4 \pi}\left[\int_{M} R+2 \int_{\partial M} K\right]-\frac{1}{2} \int_{M} \Phi(R+2)-\int_{\partial M} \Phi(K-1)
$$

- Gauss-Bonnet and $\Phi$ path integral:

$$
I=-S_{0} \chi(M)-\int_{\partial M} \Phi(K-1)
$$

## JT and Boundary Conditions

- JT action:

$$
I=-\frac{S_{0}}{4 \pi}\left[\int_{M} R+2 \int_{\partial M} K\right]-\frac{1}{2} \int_{M} \Phi(R+2)-\int_{\partial M} \Phi(K-1)
$$

- Gauss-Bonnet and $\Phi$ path integral:

$$
I=-S_{0} \chi(M)-\int_{\partial M} \Phi(K-1)
$$

- Boundary conditions:

Cutoff boundaries identified with level sets of the dilaton $\left.\Phi\right|_{\partial M}=1 / \delta$. Limit $\delta \rightarrow 0$ taken with fixed ratio $L_{\partial M} /\left.\Phi\right|_{\partial M}=\beta$


## An Interacting Schwarzian Theory

- Near-boundary metric:

$$
g=\left(\frac{1}{(1-\xi)^{2}}+h_{a}^{(m)}(\phi)+\mathcal{O}(1-\xi)\right)\left(d \xi^{2}+d \phi^{2}\right)
$$

## An Interacting Schwarzian Theory

- Near-boundary metric:

$$
g=\left(\frac{1}{(1-\xi)^{2}}+h_{a}^{(m)}(\phi)+\mathcal{O}(1-\xi)\right)\left(d \xi^{2}+d \phi^{2}\right)
$$

- Wiggle $\phi: \mathbb{S}_{\beta} \rightarrow \partial M \cong \mathbb{S}$ defined by $\left.g\right|_{\partial M}=d u^{2} / \delta^{2}$

$$
-\int_{\partial M} \Phi(K-1)=\int_{\mathbb{S}_{\beta}} d u\left(\left\{\tan \frac{\phi}{2}, u\right\}+\frac{1}{2}\left(1+3 h_{a}^{(m)}(\phi)\right) \phi^{\prime}(u)^{2}\right)
$$

Note $h_{a}^{(m)}(\phi)=-\frac{1}{3}$ for $m=1,2$ leave the Schwarzian alone

## An Interacting Schwarzian Theory

- Near-boundary metric:

$$
g=\left(\frac{1}{(1-\xi)^{2}}+h_{a}^{(m)}(\phi)+\mathcal{O}(1-\xi)\right)\left(d \xi^{2}+d \phi^{2}\right)
$$

- Wiggle $\phi: \mathbb{S}_{\beta} \rightarrow \partial M \cong \mathbb{S}$ defined by $\left.g\right|_{\partial M}=d u^{2} / \delta^{2}$

$$
-\int_{\partial M} \Phi(K-1)=\int_{\mathbb{S}_{\beta}} d u\left(\left\{\tan \frac{\phi}{2}, u\right\}+\frac{1}{2}\left(1+3 h_{a}^{(m)}(\phi)\right) \phi^{\prime}(u)^{2}\right)
$$

Note $h_{a}^{(m)}(\phi)=-\frac{1}{3}$ for $m=1,2$ leave the Schwarzian alone

- Wiggle equation of motion:

$$
\left(\frac{1}{\phi^{\prime}}\left(\frac{\phi^{\prime \prime}}{\phi^{\prime}}\right)^{\prime}-3 h_{a}^{(m)}(\phi) \phi^{\prime}\right)^{\prime}+\frac{3}{2}\left(h_{a}^{(m)}(\phi)\right)^{\prime} \phi^{\prime}=0
$$

## The JT Hellscape

- Explicit JT wiggle action for all replica $m \in \mathbb{R}_{\geq 0}$ and moduli $a$ :

$$
I_{a}^{(m)}(\beta)=\frac{8}{\beta} \operatorname{arcsinh}^{2} \sqrt{\sin ^{2} \frac{\pi}{m} \cosh ^{2} \frac{a}{2}-1}
$$

Note $I_{a}^{(1)}(\beta)=-2 \pi^{2} / \beta, I_{a}^{(2)}(\beta)=2 a^{2} / \beta$

## The JT Hellscape

- Explicit JT wiggle action for all replica $m \in \mathbb{R}_{\geq 0}$ and moduli $a$ :

$$
I_{a}^{(m)}(\beta)=\frac{8}{\beta} \operatorname{arcsinh}^{2} \sqrt{\sin ^{2} \frac{\pi}{m} \cosh ^{2} \frac{a}{2}-1}
$$

Note $I_{a}^{(1)}(\beta)=-2 \pi^{2} / \beta, I_{a}^{(2)}(\beta)=2 a^{2} / \beta$

- Action and stability analysis...

... and the modulus is not stabilized


## Adding Matter

- Conformal matter with classical sources:

$$
I_{\mathrm{mat}}=\frac{1}{2} \int_{M}(d \psi)^{2} .
$$

## Adding Matter

- Conformal matter with classical sources:

$$
I_{\mathrm{mat}}=\frac{1}{2} \int_{M}(d \psi)^{2}
$$

- Solve in terms of boundary profile specified by $\psi_{0}: \partial M \rightarrow \mathbb{R}$

$$
I_{\mathrm{mat}}=\frac{1}{2} \int_{\partial M} \psi \nabla_{n} \psi=\frac{1}{2} \int_{\mathbb{S}} d \phi d \tilde{\phi} \psi_{0}(\phi) S(\phi, \tilde{\phi}) \psi_{0}(\tilde{\phi})
$$

where $S$ is an integral kernel known explicitly but gross

## Adding Matter

- Conformal matter with classical sources:

$$
I_{\mathrm{mat}}=\frac{1}{2} \int_{M}(d \psi)^{2}
$$

- Solve in terms of boundary profile specified by $\psi_{0}: \partial M \rightarrow \mathbb{R}$

$$
I_{\mathrm{mat}}=\frac{1}{2} \int_{\partial M} \psi \nabla_{n} \psi=\frac{1}{2} \int_{\mathbb{S}} d \phi d \tilde{\phi} \psi_{0}(\phi) S(\phi, \tilde{\phi}) \psi_{0}(\tilde{\phi})
$$

where $S$ is an integral kernel known explicitly but gross

- Wiggle equation of motion

$$
\left(\frac{1}{\phi^{\prime}}\left(\frac{\phi^{\prime \prime}}{\phi^{\prime}}\right)^{\prime}-3 h_{a}(\phi) \phi^{\prime}\right)^{\prime}+\frac{3}{2}\left(h_{a}(\phi)\right)^{\prime} \phi^{\prime}+\psi_{0}(\phi) \int_{\mathbb{S}} d \tilde{\phi} S(\phi, \tilde{\phi}) \psi_{0}(\tilde{\phi})=0
$$

## A Silver Lining

- Modulus saddles appear for $m=2$ as sources are turned on:



## A Silver Lining

- Modulus saddles appear for $m=2$ as sources are turned on:



## A Silver Lining

- Modulus saddles appear for $m=2$ as sources are turned on:



## A Silver Lining

- Modulus saddles appear for $m=2$ as sources are turned on:

- Will modulus saddles make it all the way to $m<1$ ?


## A Silver Lining

- Modulus saddles appear for $m=2$ as sources are turned on:

- Will modulus saddles make it all the way to $m<1$ ?

Yes!

## The Little Saddle that Could

- Pair of stable/unstable branches of solutions exist for $m<1$ !
- The little wormhole can be made to dominate over the disk one
- Action and stability analysis for $m=.75$ :




## The Journey Just Began



Little saddle spotted at $m=.75$

## The Journey Just Began



Little saddle spotted at $m=.7$

## The Journey just Began



Little saddle spotted at $m=.65$

## The Journey Just Began



Little saddle spotted at $m=.6$

## The Journey Just Began



Little saddle spotted at $m=.55$

DISCUSSION

## OUTLOOK

- Will the little saddle make it to $m \rightarrow 0$ ?
- What properties does the resulting generating functional have?
- How does it differ from the annealed result?
- What is the effect on scalar correlation functions?


## Open Questions and Future Directions

- Is there a simple diagnostic for when quenched $m \rightarrow 0$ saddle points will differ from annealed ones?
- Is there any correlation between dominance of replica wormholes for $m \in \mathbb{Z}_{+}$and for $0<m<1$ ?
- Are there any universal features about $m \rightarrow 0$ saddle points and quenched generating functionals in quantum gravity?
- Other toy models for the study $m \rightarrow 0$ saddles?


## Open Questions and Future Directions

- Is there a simple diagnostic for when quenched $m \rightarrow 0$ saddle points will differ from annealed ones?
- Is there any correlation between dominance of replica wormholes for $m \in \mathbb{Z}_{+}$and for $0<m<1$ ?
- Are there any universal features about $m \rightarrow 0$ saddle points and quenched generating functionals in quantum gravity?
- Other toy models for the study $m \rightarrow 0$ saddles?

Thank you for listening!

