#### Jones modes in Lipschitz domains

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Spectral Geometry: Theory, Numerical Analysis and Applications BIRS, Canada

July 6, 2018

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### Fluid-solid interaction

Hemholtz equation for p:

$$\Delta p + (w/c)^2 p = 0,$$

Linear elasticity:  $\mu_s > 0$ ,  $\lambda_s + \left(\frac{2}{d}\right) \mu_s > 0$ 

$$\operatorname{div} \boldsymbol{\sigma}(\mathbf{u}_s) + w^2 \rho_s \tilde{\mathbf{u}}_s = \mathbf{0}.$$

BC's on the interface:

$$\sigma(\mathbf{u}_s)\mathbf{n} = -(p + p_{inc})\mathbf{n},$$
  
$$w^2 \rho_f \mathbf{u}_s \cdot \mathbf{n} = \nabla(p + p_{inc}) \cdot \mathbf{n}.$$

Condition at infinity:

$$\frac{\partial p}{\partial r} - i(w/c)p = o(1/r), \qquad r := \|\mathbf{x}\|.$$

<sup>2</sup>Hsiao, Kleinman and Roach 2000; Hsiao, Xu and Yin 2017 ( B > ( E >



#### Non-uniqueness

#### Lemma

If  $(\mathbf{u}_s, p)$  solves the time harmonic fluid-solid interaction problem then  $(\mathbf{u}_s + \mathbf{u}_0, p)$  also solves this problem, with  $\mathbf{u}_0$  a non-zero solution of

 $\operatorname{div} \boldsymbol{\sigma}(\mathbf{u}_0) + \rho_s w^2 \mathbf{u}_0 = \mathbf{0}, \text{ in } \Omega_s, \quad \boldsymbol{\sigma}(\mathbf{u}_0)\mathbf{n} = \mathbf{0}, \quad \mathbf{u}_0 \cdot \mathbf{n} = 0, \text{ on } \Gamma.$ 



#### Note

- Condition on shear along the interface;
- Robin condition in an "artificial" boundary away from the solid;
- no eigenpairs for  $C^{\infty}$  domains in  $\mathbb{R}^3$ .

HKR, 2000; Gatica et al., 2009; Barucq et al., 2014; T. Hargé, 1990

#### The EV problem: Jones modes

This problem is not uniquely solvable when  $w^2 \rho_s$  is an eigenvalue of

$$\begin{aligned} &\operatorname{div} \boldsymbol{\sigma}(\mathbf{u}_s) + w^2 \rho_s \mathbf{u}_s = \mathbf{0}, & \operatorname{in} \, \Omega_s, \\ & \boldsymbol{\sigma}(\mathbf{u}_s)\mathbf{n} = \mathbf{0}, & \mathbf{u}_s \cdot \mathbf{n} = \mathbf{0}, & \operatorname{on} \, \partial \Omega_s. \end{aligned}$$

A non-trivial solution  $\mathbf{u}_s$  for some  $w^2$  is called a *Jones mode*.

Note

Over-determined EV:

Elasticity equation + Traction condition + extra constraint on  $\mathbf{u}_s$ .

<sup>2</sup> Jones et al.,	1983	and	1984
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#### shear and compression modes

*s*-waves: div 
$$\mathbf{u}_s = \mathbf{0}$$
 *p*-waves: rot  $\mathbf{u}_s = \mathbf{0}$ 

On  $[0, a] \times [0, b]$ ,

$$w_{mn}^{2} := \begin{cases} \left(\frac{\pi^{2}\mu}{\rho}\right) \left(\frac{m^{2}}{a^{2}} + \frac{n^{2}}{b^{2}}\right), \\ \left(\frac{\pi^{2}(\lambda+2\mu)}{\rho}\right) \left(\frac{m^{2}}{a^{2}} + \frac{n^{2}}{b^{2}}\right), \end{cases}$$

with eigenfunctions:

$$\mathbf{u}_{mn} := \begin{cases} an\sin\left(\frac{m\pi}{a}x\right)\cos\left(\frac{n\pi}{b}y\right)\hat{i} - bm\cos\left(\frac{m\pi}{a}x\right)\sin\left(\frac{n\pi}{b}y\right)\hat{j},\\ bm\sin\left(\frac{m\pi}{a}x\right)\cos\left(\frac{n\pi}{b}y\right)\hat{i} + an\cos\left(\frac{m\pi}{a}x\right)\sin\left(\frac{n\pi}{b}y\right)\hat{j}, \end{cases}$$

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### Weak formulation (on $\Omega_s$ )

Consider  $\mathbf{H}^1(\Omega) := H^1(\Omega) \times H^1(\Omega)$  and define

$$\mathbf{H} := \Big\{ \mathbf{u} \in \mathbf{H}^1(\Omega) : \ \mathbf{u} \cdot \mathbf{n} = 0 \quad \text{on } \partial \Omega \Big\}.$$

FORMULATION: find  $(w^2, \mathbf{u}) \in \mathbb{C} \times \mathbf{H}$  such that

$$\mathbf{a}(\mathbf{u},\mathbf{v})=
ho w^2(\mathbf{u},\mathbf{v}) \qquad orall \mathbf{v}\in \mathbf{H},$$

where  $a(\mathbf{u}, \mathbf{v}) := \mu(\nabla \mathbf{u}, \nabla \mathbf{v}) + (\lambda + \mu)(\operatorname{div} \mathbf{u}, \operatorname{div} \mathbf{v})$  or  $(\sigma(\mathbf{u}), \epsilon(\mathbf{v}))$ .

#### Note

• 
$$a(\mathbf{u},\mathbf{v}) \leq (\lambda+\mu) \|\mathbf{u}\|_1 \|\mathbf{v}\|_1;$$

• Rayleigh quotient + properties of  $a(\cdot, \cdot)$  and  $(\cdot, \cdot) \Rightarrow w^2 \ge 0$ .

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### Ellipticity of a

We can get

### $m{a}(\mathbf{u},\mathbf{u})\geq\minig\{2\mu,d(\lambda+(2/d)\mu)ig\}\|m{\epsilon}(\mathbf{u})\|_0^2,\quadorall\,\mathbf{u}\in\mathbf{H}^1(\Omega).$

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#### Theorem (Bauer 2016, Domínguez 2018)

Let  $\Omega$  be a non-axisymmetric bounded and Lipschitz domain in  $\mathbb{R}^d$ . Then, there is a positive constant C > 0 such that

 $\|\boldsymbol{\epsilon}(\mathbf{u})\|_0 \geq C \|\mathbf{u}\|_1, \quad \forall \, \mathbf{u} \in \mathbf{H}.$ 

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#### Lemma

Under the same assumptions for  $\Omega$ , there is a constant C > 0 such that

 $a(\mathbf{u},\mathbf{u}) \geq c \|\mathbf{u}\|_1^2, \quad \forall \, \mathbf{u} \in \mathbf{H}.$ 

The corresponding solution operator T is then

- linear and bounded with  $||T||_{\mathbf{H}'} = \frac{\rho}{C}$ ;
- compact from **H** to itself;
- self-adjoint w.r.t.  $a(\cdot, \cdot)$ .
- Spectral Theorem  $\Rightarrow$  eigenpairs  $\{w_n\}$  and  $\{u_n\}$  with  $w_n \rightarrow +\infty$ .

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### **Rigid motions**

#### Lemma

- $w^2 = 0$  is an eigenvalue with:
  - (i) a pure translation as eigenfunction if  $\partial \Omega$  consists of two parallel planes;
- (ii) a pure rotation  $\mathbf{u}_0$  as eigenfunction if  $\Omega$  is axisymmetric about the axis of rotation of  $\mathbf{u}_0$ .

Shifted formulation: find  $(\mathbf{u}, w^2) \in \mathbf{H} imes \mathbb{R}$  such that

$$\widetilde{a}(\mathbf{u},\mathbf{v}):=a(\mathbf{u},\mathbf{v})+
ho(\mathbf{u},\mathbf{v})=
ho(w^2+1)(\mathbf{u},\mathbf{v})\qquadorall\mathbf{v}\in\mathbf{H},$$

- $\tilde{a}(\mathbf{u},\mathbf{u}) \geq \min\{\mu,\rho\} \|\mathbf{u}\|_{1}^{2};$
- the corresponding solution operator  $\tilde{\mathcal{T}}$  is well-defined, compact and self-adjoint;

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#### Discrete scheme

Let  $\mathbf{H}_h \subseteq \mathbf{H}$  (Lagrange elements): find  $\mathbf{u}_h \in \mathbf{H}_h$  such that

$$a(\mathbf{u}_h,\mathbf{v}_h) = \rho \kappa_h(\mathbf{u}_h,\mathbf{v}_h), \quad \forall \, \mathbf{v}_h \in \mathbf{H}_h,$$

with  $\kappa_h := w_h^2$  or  $w_h^2 + 1$ .

- a is H<sub>h</sub>-elliptic;
- a discrete solution operator  $T_h$  (cf.  $\tilde{T}_h$ ) is well defined;
- error bound for evs:

$$\frac{|\kappa-\kappa_h|}{\kappa} \leq Ch^{2(t-1)}, \quad t>1.$$

operator approximation:

$$\|T-T_h\|\leq ch^{t-1}.$$

<sup>2</sup>Babuška and Osborn, 1991.

# Square ( $\mu = \lambda = 1$ )





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# Square (contd.)

j	w <sup>2</sup>	$w^2/\pi^2$	$\ \operatorname{div} \mathbf{u}\ _0^2$	$\ \operatorname{rot} \mathbf{u}\ _0^2$	x-component	y-component
1	19.74	2.000	9.870	0.0002633		
2	19.74	2.000	9.870	0.0001704		
3	19.74	2.000	5.883e-05	19.72		
4	39.48	4.000	19.74	0.0005061		

Table: Unit square with parameters  $\mu = \rho = 1$ ,  $\lambda = 0$ .

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# Square (contd.)

j	$\nu_j$	$\nu_j/\pi^2$	$\ \operatorname{div} \mathbf{u}\ _0^2$	$\ \operatorname{rot} \mathbf{u}\ _0^2$	x-component	y-component
1	12.34	1.250	1.231e-07	12.27		
2	19.74	2.000	9.515e-07	19.69	estina 🗐	
3	29.61	3.000	2.467	2.271e-05		
4	32.08	3.250	2.74e-06	32.05		
5	41.95	4.250	2.223e-06	41.62		

Table: 2X1 rectangle with parameters  $\mu = \rho = 1, \lambda = 10$ .

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# Square (contd.)

j	$\nu_j$	$\nu_j/\pi^2$	$\ \operatorname{div} \mathbf{u}\ _0^2$	$\ \operatorname{rot} \mathbf{u}\ _0^2$	x-component	y-component
1	51.82	5.25	2.467	2.271e-05		
2	123.4	12.5	2.271e-07	12.27		
3	197.4	20	1.757e-06	19.69		
4	207.3	21	9.869	0.0004004		
5	207.3	21	9.87	0.000243		

Table: 2X1 rectangle with parameters  $\mu = 10, \lambda = \rho = 1$ .

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## Triangle

j	w <sup>2</sup>	$w^2/\pi^2$	$\ \operatorname{div} \mathbf{u}\ _0^2$	$\ \operatorname{rot} \mathbf{u}\ _0^2$	x-component	y-component
1	4.6563	0.4718	0.7007	24.36		
2	8.3125	0.8422	0.4333	14.42		
3	11.84674	1.200	2.527	4.15		
4	21.0647	2.134	1.640	75.96		

Table: Isosceles triangle of vertices (0,0), (2,0) and (1,2) with parameters  $\lambda = \mu = \rho = 1$ .

# L-shape, $\rho = \mu = \lambda = 1$



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#### CV properties on polyhedron



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### Conclusions and future work

Conclusions:

- FEM provides a reliable scheme to approximate Jones modes on polyhedral domains;
- the extra constraint makes the problem "domain dependent".

Future work:

- Does this scheme work for more general smooth domains?
- A posteriori error analysis would help to improve the convergence for computations on non-convex domains.

# Thanks!

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