

BPS state counting in K3 string theories

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in collaboration with:



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Introduction

- ✧ String models with **16** supersymmetries arising from type IIA/B on $K3 \times T^2$ and their orbifolds (CHL models)

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- ✧ **Goal:** Calculate space-time index counting 1/4 BPS states for each value of electric-magnetic charges
- ✧ Some of the first examples of matching BH black hole entropy vs microscopic degeneracy [Strominger, Vafa 95; Dijkgraaf, Verlinde² 96
Shih, Strominger, Yin 2005; David, Sen 2006; ...]

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- ※ All these indices organized into a generating function $1/\Phi$
- ※ Meromorphic Siegel modular form of genus 2
- ※ For type IIA/ $K3 \times T^2$, Φ is Igusa cusp form of weight 10

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2-dim superconformal field theory
- ✧ Elliptic genus of K3 $\phi(\tau, z)$
- ✧ If NLSM on K3 has a symmetry g , we can define a **twining genus** ϕ_g
- ✧ ϕ_g is weak Jacobi form wt 0 ind 1

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- ✧ some ϕ_g cannot be computed directly in the NLSM
- ✧ **Idea:** use consistency conditions on $1/\Phi$ from wall-crossing
- ✧ **Result:** for almost all g , ϕ_g is uniquely determined, otherwise only **2** possibilities



Overview

Introduction

NLSM on $K3$

Strings on $K3 \times T^2$

CHL models

NLSM on K3

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Basic facts about NLSM on $K3$

✳ 2-dim $\mathcal{N} = (4, 4)$ superconformal at $c = \bar{c} = 6$

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[Aspinwall, Morrison '95; Nahm, Wendland '99]

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- ✳ Elliptic genus of K3:

$$\phi_{K3}(\tau, z) = \text{Tr}_{RR}(q^{L_0 - \frac{c}{24}} \bar{q}^{\bar{L}_0 - \frac{\bar{c}}{24}} y^{J_0} (-1)^{J_0 + \bar{J}_0})$$

where $(\tau, z) \in \mathbb{H} \times \mathbb{C}$, $q = e^{2\pi i \tau}$, $y = e^{2\pi i z}$

J_0, \bar{J}_0 are generators in $su(2)$ in $\mathcal{N} = 4$

[Schellekens, Warner '86; Witten '87;

Eguchi, Ooguri, Taormina, Yang '88]

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$$\text{✳ } \phi(\tau, z) = \sum_{n \geq 0, l} c(n, l) q^n y^l$$

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$$\phi_g(\tau, z) = \text{Tr}_{RR}(gq^{L_0 - \frac{c}{24}}\bar{q}^{\bar{L}_0 - \frac{\bar{c}}{24}}y^{J_0}(-1)^{J_0 + \bar{J}_0})$$

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Can we find all of them?

✧ In NLSM(K3) with symmetry g , define

$$\mathcal{H}_{r,s} := \left\{ v \text{ in } g^r\text{-twisted sector} \right. \\ \left. \text{s.t. } g(v) = e^{\frac{2\pi i s}{N}} v \right\}$$

for $r, s \in \mathbb{Z}/N\mathbb{Z}$

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$$\phi_{g^n} = \sum_{s \in \mathbb{Z}/N} e^{\frac{2\pi i n s}{N}} \phi_{r,s}$$

Strings on $K3 \times T^2$

Strings on K3

$$\text{het}/T^6 \quad \leftrightarrow \quad \text{IIA}/K3 \times T^2 \quad \leftrightarrow \quad \text{IIB}/K3 \times T^2$$

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[Sen '94; Witten '94]

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- ✧ 134-dim moduli space

1/4 BPS index

✧ For each (Q, P) consider 1/4 BPS index $D(Q, P)$

$$D(Q, P) = \#\{\text{'bosonic' supermultiplets}\} \\ - \#\{\text{'fermionic' supermultiplets}\}$$

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- ✳ **Goal:** Find $d(n, m, l)$ for all $n, m, l \in \mathbb{Z}$

1/4 BPS index

✧ Organize into a generating function

$$\frac{1}{\Phi \begin{pmatrix} \sigma & z \\ z & \tau \end{pmatrix}} = \sum_{n,m,l} d(n,m,l) e^{2\pi i(m\sigma + n\tau + lz)}$$

for some complex 'chemical potentials' σ, τ, z

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✧ Given by exp-lift of elliptic genus

$$\frac{1}{\Phi\left(\begin{smallmatrix} \sigma & z \\ z & \tau \end{smallmatrix}\right)} = \prod_{n,m \geq 0, l} (1 - e^{2\pi i(m\sigma+n\tau+lz)})^{-c(mn,l)}$$

($l < 0$ if $m = n = 0$)

where $c(mn, l)$ are Fourier coeffs of ϕ_{K3}

[Dijkgraaf, Verlinde, Verlinde '96; Shih, Strominger, Yin 2005; David, Sen 2006]

Wall crossing

- ✧ Some $1/4$ BPS states can decay into pair of $1/2$ BPS in subregions of the moduli space
 $\Rightarrow D(Q, P)$ 'jumps' across wall of marginal stability [Bergman '97; Bergman, Kol '98; Denef '00]

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- ✧ Fourier coefficients of $1/\Phi$

$$d(n, m, l) = \oint_{\mathcal{C}} \frac{e^{-2\pi i(m\sigma + n\tau + lz)}}{\Phi\left(\begin{matrix} \sigma & z \\ z & \tau \end{matrix}\right)}$$

where

$$\mathcal{C} := \{0 \leq \sigma_1, \tau_1, z_1 \leq 1, (\begin{matrix} \sigma_2 & z_2 \\ z_2 & \tau_2 \end{matrix}) \text{ fixed}\}$$

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- ✧ $1/\Phi$ is meromorphic $\Rightarrow d(n, m, l)$ ambiguous

Wall crossing

Contour prescription: \mathcal{C} depends on moduli and charges

$$\mathcal{C} := \{0 \leq \sigma_1, \tau_1, z_1 \leq 1, \begin{pmatrix} \sigma_2 & z_2 \\ z_2 & \tau_2 \end{pmatrix} = \epsilon^{-1} \mathcal{Z}(Q, P, \mu)\}$$

where $\epsilon \ll 1$ and \mathcal{Z} is the ‘central charge vector’

[Cheng, Verlinde 2007]

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- ✳ Poles of $1/\Phi$ exactly at walls of marginal stability
- ✳ Residue at pole matches the ‘jump’ of $D(Q, P)$
- ✳ Independent interpretations

[Banerjee, Sen, Srinistava 2008; Bossard, Cosnier-Horeau, Pioline 2016]

CHL models

CHL model definition

✧ Consider type II/ $K3 \times S^1 \times \tilde{S}^1$

[Chaudhuri, Hockney, Lykken 95; Sen, Vafa 95]

CHL model definition

- ✧ Consider type II/ $K3 \times S^1 \times \tilde{S}^1$
- ✧ Orbifold by (δ, g) , where
 - ✧ g is symmetry of NLSM on $K3$, $ord(g) = N$, preserving $\mathcal{N} = 4$ SUSY
 - ✧ δ is a shift by $1/N$ period along S^1

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Obtain 4-d $\mathcal{N} = 4$ theory with

- ✧ reduced gauge group $U(1)^d$ ($6 \leq d \leq 28$)
- ✧ reduced moduli space
- ✧ e-m charges $(Q, P) \in \Lambda_e \oplus \Lambda_m$ ($rk \Lambda_{e,m} = d$)
- ✧ Duality group $\supset O(\Lambda_e)$

[Chaudhuri, Hockney, Lykken 95; Sen, Vafa 95]

1/4 BPS states

✧ 1/4 BPS index depends on $O(\Lambda_e)$ -invariants

$$D_g(P, Q) = (-1)^{P \cdot Q + 1} d_g(P^2/2, Q^2/2, P \cdot Q)$$

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✧ Generating function $1/\Phi_g$ of $d_g(n, m, l)$ is exponential lift of twining genera [Jatkar, Sen 05]

$$\frac{1}{\Phi_g \left(\begin{smallmatrix} \sigma & z \\ z & \tau \end{smallmatrix} \right)} = \prod_{(n, m, l)} (1 - e^{2\pi i(m\sigma + \frac{n\tau}{N} + lz)})^{-c_{n, m}^{(g)}(\frac{nm}{N}, l)}$$

$c_{n, m}^{(g)}(\frac{nm}{N}, l)$: linear combinations of Fourier coeffs of ϕ_{g^d} , for $d|N$, and their $SL(2, \mathbb{Z})$ transformations

Wall crossing in CHL models

- ✧ Similar wall-crossing for 'decaying' of $1/4$ BPS to pair of $1/2$ BPS states [Sen 2007]

Wall crossing in CHL models

- ✧ Similar wall-crossing for 'decaying' of 1/4 BPS to pair of 1/2 BPS states [Sen 2007]

Main assumption: Contour prescription provides the correct $d_g(n, m, l)$ at all points in moduli space

$\Rightarrow 1/\Phi_g$ has only poles at the walls

How to compute the remaining $1/\Phi_g$?

✱ We can determine walls of marginal stability

[Paquette, Volpato, Zimet 2017]

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- ✳ Via exponential lift \rightarrow sign of $c_{m,n}^{(g)}(\frac{mn}{N}, l)$ for $4\frac{mn}{N} - l^2 < 0$

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- ✳ Via exponential lift \rightarrow sign of $c_{m,n}^{(g)}(\frac{mn}{N}, l)$ for $4\frac{mn}{N} - l^2 < 0$

This info + modularity + q^0 Fourier coeffs



ϕ_g and $1/\Phi_g$ determined almost completely
(either uniquely or up to **2** possibilities)

[Paquette, Volpato, Zimet 2017]

Open questions

- ✧ Can all $1/\Phi_g$ be determined exactly?
- ✧ Can $1/\Phi_g$ be fixed without passing through NLSM(K3)?
- ✧ Relations to Umbral + Conway moonshine