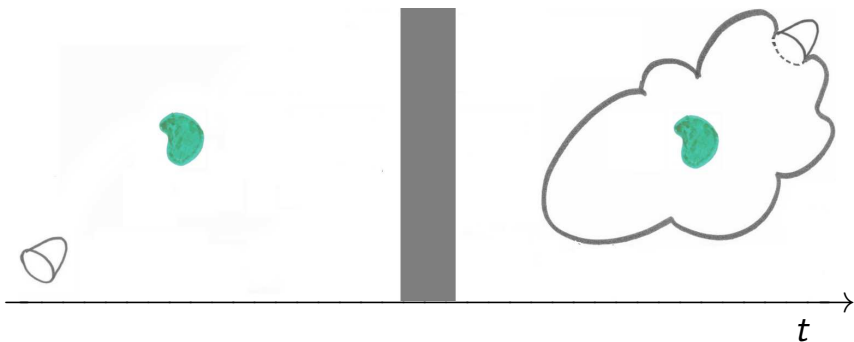


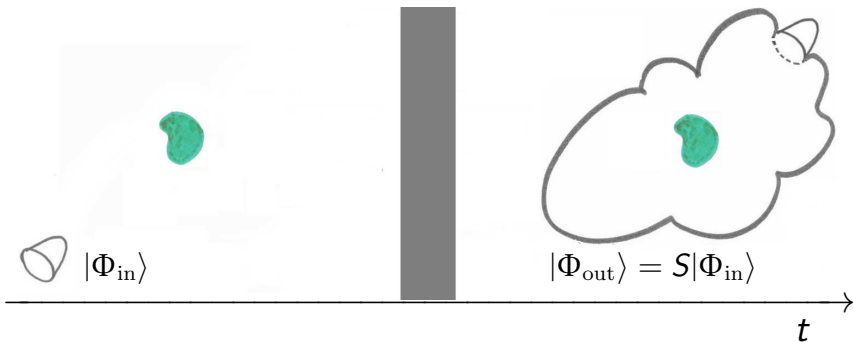
An algebraic approach to light-matter interactions

Ivan Fernandez-Corbaton

Institute of Nanotechnology, Karlsruhe Institute of Technology







Algebraic approach to light-matter interactions:

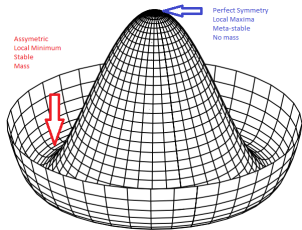
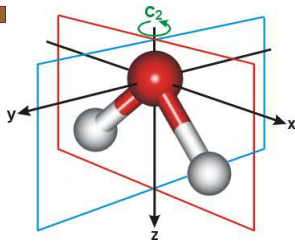
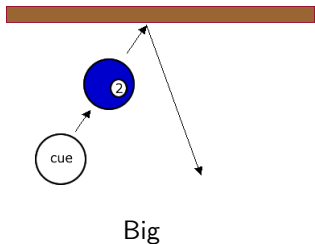
- Incoming field $|\Phi_{in}\rangle \in \mathbb{M}^{in}$
 - Vector space of **incoming** solutions of Maxwell's equations
- Outgoing field $|\Phi_{out}\rangle \in \mathbb{M}^{out}$
 - Vector space of **outgoing** solutions of Maxwell's equations
- Effect of the object represented by scattering operator S
 - Linear and bounded map between incoming and outgoing fields

- Easy to use symmetries and light-matter conservation laws
 - Symmetry results are general
 - “(Breaking of) Symmetry X is needed for effect Y”
 - Provide design guidelines
- Algebraic tools ease the solution of complicated problems
 - Given an object, find the best beam to exert force on it
 - Find eigenvalues of matrices
- Algebraic \implies very well suited for computer implementation

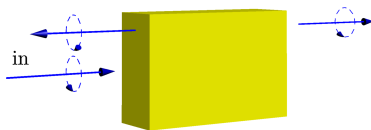
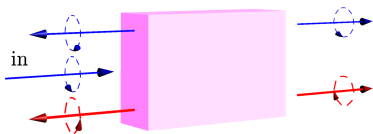
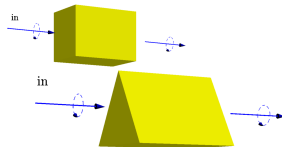
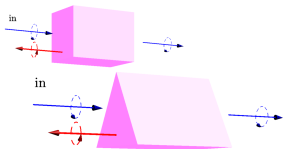
Allows straightforward use of **symmetry arguments**.

Symmetry

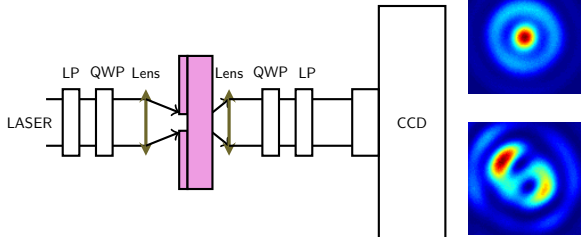
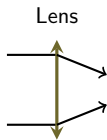
Beautifully general concept, central in theoretical physics ...



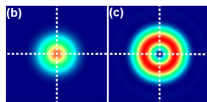
How can symmetry help in the study of light-matter interactions?



Why does zero back scattering happen?

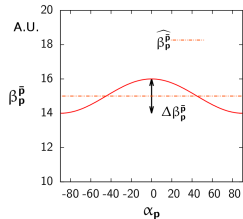
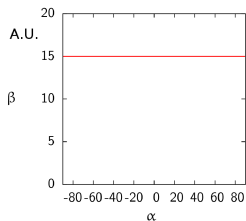
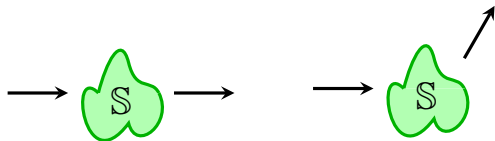
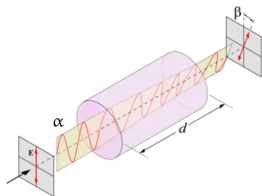


PRL 99, 073901 (2007)

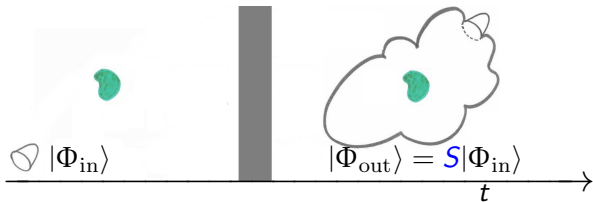


Why do optical vortices appear?

Solution of chiral molecules



What is the origin of this difference?



- Symmetry is the invariance upon transformations
- Translations, rotations, electromagnetic duality, time translation, parity, time inversion, ...
- Transformations represented by unitary operators in $\mathbb{M}^{in(out)}$
- Symmetry of S under transformation X means $XSX^{-1} = S$

$$XSX^{-1} = S \iff SX - XS = [S, X] = 0$$

$$\text{If } X|\Phi_{in}\rangle = x|\Phi_{in}\rangle \implies X|\Phi_{out}\rangle = x|\Phi_{out}\rangle$$

For continuous symmetries

$$X_\theta = \exp(-i\theta\Gamma) \text{ with } \Gamma^\dagger = \Gamma$$

$$[S, X_\theta] = 0 \iff [S, \Gamma] = 0$$

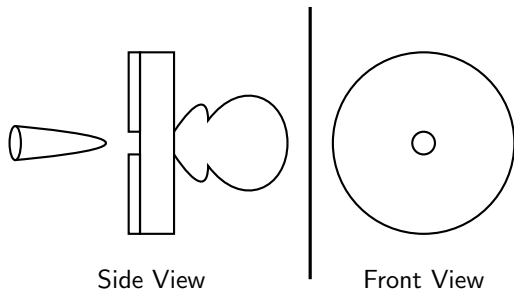
$$\text{If } \Gamma|\Phi_{\text{in}}\rangle = \gamma|\Phi_{\text{in}}\rangle \implies \Gamma|\Phi_{\text{out}}\rangle = \gamma|\Phi_{\text{out}}\rangle$$

In light-matter interactions:

- Symmetry: Non-coupling eigenstates with different eigenvalue
- Broken symmetry: Coupling eigenstates with different eigenvalue

Generator Γ	Transformation $\exp(-i\theta\Gamma)$	Expression
Linear momentum \mathbf{P}	Spatial translations	$-i\nabla$
Hamiltonian H	Time translations	$i\partial_t$
Angular momentum \mathbf{J}	Rotations	$-\mathbf{ir} \times \nabla - i\epsilon_{knm}$
Helicity Λ	Electromagnetic duality	$\frac{\nabla \times}{k}$

Example 1: Angular momentum and Rotational symmetry



- If the input is an eigenstate $J_z|\Phi_{\text{in}}\rangle = m|\Phi_{\text{in}}\rangle$,
- so is the output: $J_z|\Phi_{\text{out}}\rangle = m|\Phi_{\text{out}}\rangle$
- **with the same eigenvalue**

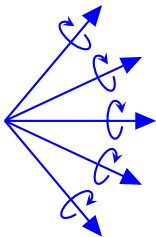
Example 2: Helicity and Duality symmetry

- Operator:

$$\Lambda = \frac{\mathbf{J} \cdot \mathbf{P}}{|\mathbf{P}|} \equiv \frac{\nabla \times}{k}$$

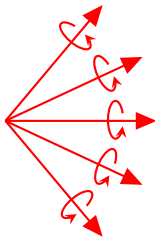
- Eigenstates: $\Lambda (\mathbf{E} \pm i\mathbf{ZH}) = \pm (\mathbf{E} \pm i\mathbf{ZH}) = \pm \mathbf{G}_{\pm}$
- Interpretation in the plane wave decomposition

$$\Lambda \mathbf{G}_+ = \mathbf{G}_+$$



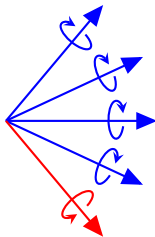
$$\mathbf{G}_- = 0$$

$$\Lambda \mathbf{G}_- = -\mathbf{G}_-$$



$$\mathbf{G}_+ = 0$$

$$\Lambda \mathbf{G} \neq c\mathbf{G}$$



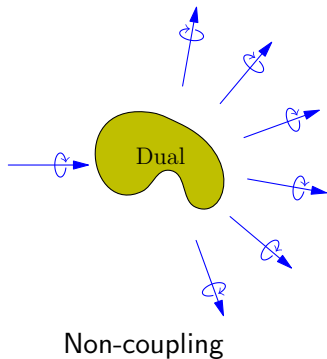
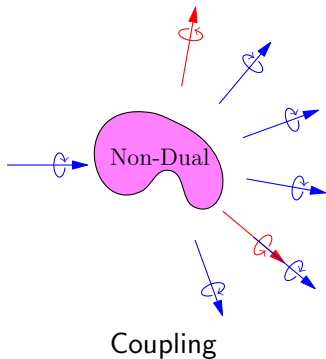
$$\mathbf{G}_+ \neq 0, \mathbf{G}_- \neq 0$$

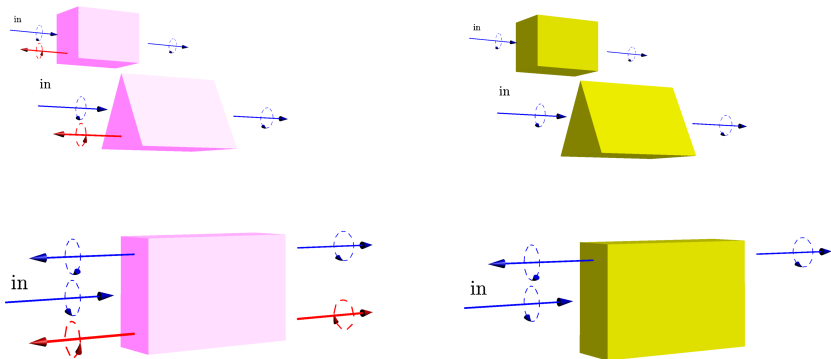
Electromagnetic duality transformation: $D(\theta) = \exp(-i\theta\Lambda)$

$$\mathbf{G}_{\pm} \rightarrow \exp(\mp i\theta) \mathbf{G}_{\pm}$$

$$\mathbf{E} \rightarrow \mathbf{E} \cos \theta - Z\mathbf{H} \sin \theta$$

$$Z\mathbf{H} \rightarrow \mathbf{E} \sin \theta + Z\mathbf{H} \cos \theta$$

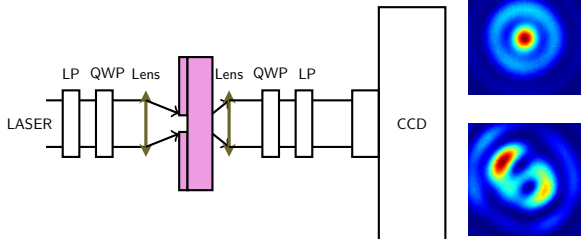
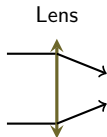




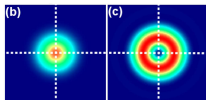
Why does zero back scattering happen?

Duality symmetry and discrete rotational symmetry $n \geq 3$

- Prisms on the left are not dual symmetric
- Prisms on the right are dual symmetric
- Common behavior due to discrete rotational symmetry



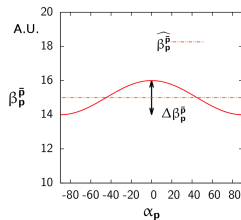
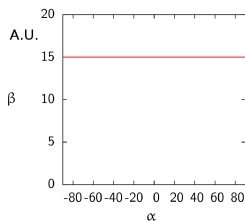
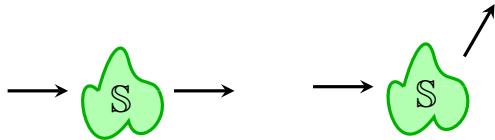
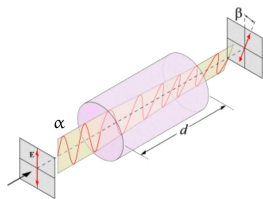
PRL 99, 073901 (2007)



Why do optical vortices appear?

- Focusing: Breaking of transverse translational symmetry
- Scattering: Breaking of duality symmetry

Solution of chiral molecules



What is the origin of this difference?

- Lack of duality symmetry of the sample
- **Chirality is necessary but not sufficient for optical activity**

Algebraic approach: Joint light-matter conservation laws

Properties like energy, momentum, angular momentum, etc .. are conserved in the **joint system of fields+matter**.

During the light-matter interaction

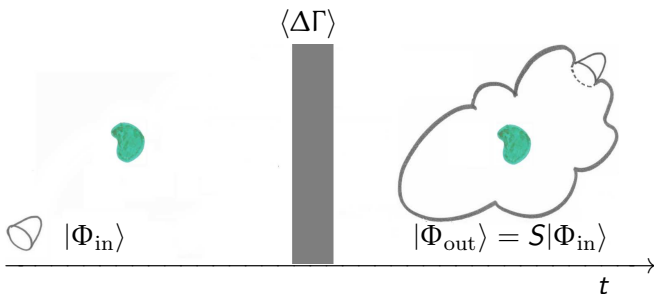
- Field loses(gains) **momentum** \implies Matter gains(looses) it
- This results in optical **force**

Algebraic approach: Joint light-matter conservation laws

Properties like energy, momentum, angular momentum, etc .. are conserved in the **joint system of fields+matter**.

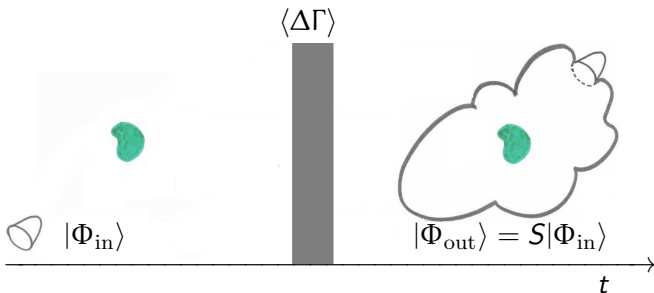
During the light-matter interaction

- Field loses(gains) **angular momentum** \implies Matter gains(looses) it
- This results in optical **torque**

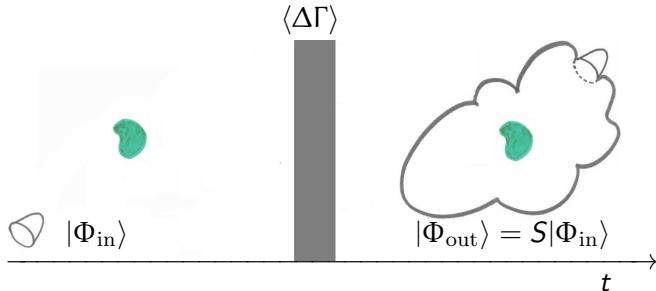


Γ is a given property (momentum, helicity, energy ...)

- ① There is some light-matter interaction
- ② Γ changes for the field
- ③ There is a corresponding change of Γ in the matter
- ④ Objective: **Quantify such change**



- ① Measure property Γ in incoming field
- ② Measure property Γ in outgoing field
- ③ Difference is the amount of Γ gained(lost) by the matter



- ① **Measure** property Γ in incoming field
- ② **Measure** property Γ in outgoing field
- ③ The difference must be the amount of Γ gained by the matter

Γ represented by a Hermitian operator

Measure of Γ in $|\Psi\rangle$:

scalar product between $|\Psi\rangle$ and $\Gamma|\Psi\rangle$: $\langle \Psi | \Gamma | \Psi \rangle$

$$\langle \Delta \Gamma \rangle = \langle \Phi_{\text{in}} | \Gamma | \Phi_{\text{in}} \rangle - \langle \Phi_{\text{out}} | \Gamma | \Phi_{\text{out}} \rangle = \langle \Phi_{\text{in}} | \Gamma - S^\dagger \Gamma S | \Phi_{\text{in}} \rangle$$

Introduce basis (e.g. multipoles, plane waves, ...).

Convenient choice: $\Gamma|\eta \gamma\rangle^{\text{in(out)}} = \gamma|\eta \gamma\rangle^{\text{in(out)}}$

$$|\Phi_{\text{in}}\rangle = \sum_{\eta, \gamma} \alpha_{\gamma}^{\eta} |\eta \gamma\rangle^{\text{in}}, \quad |\Phi_{\text{out}}\rangle = \sum_{\eta, \gamma} \beta_{\gamma}^{\eta} |\eta \gamma\rangle^{\text{out}}$$

Incoming and outgoing fields represented by column vectors

$$\underline{\alpha} \equiv \begin{bmatrix} \vdots \\ \alpha_{\gamma}^{\eta} \\ \vdots \end{bmatrix}, \quad \underline{\beta} \equiv \begin{bmatrix} \vdots \\ \beta_{\gamma}^{\eta} \\ \vdots \end{bmatrix}$$

Operators S and Γ are represented by matrices

$$\underline{\underline{S}}(\eta\gamma, \bar{\eta}\bar{\gamma}) = {}^{\text{out}}\langle \bar{\gamma} \bar{\eta} | S | \eta \gamma \rangle^{\text{in}}, \quad \underline{\underline{\Gamma}}(\eta\gamma, \bar{\eta}\bar{\gamma}) = {}^{\text{in(out)}}\langle \bar{\gamma} \bar{\eta} | \Gamma | \eta \gamma \rangle^{\text{in(out)}} = \gamma \delta_{\gamma\bar{\gamma}}$$

In practice:

- Need to compute $\underline{\underline{S}}$
- Compute T-matrix: $\underline{\underline{S}} = \underline{\underline{I}} + 2\underline{\underline{T}}$
- There are publicly available T-matrix computer codes

Abstract operations become vector/matrix algebra in a computer

$$\langle \Delta \Gamma \rangle = \langle \Phi_{\text{in}} | \Gamma - S^\dagger \Gamma S | \Phi_{\text{in}} \rangle \implies \langle \Delta \Gamma \rangle = \underline{\alpha}^\dagger \left(\underline{\Gamma} - \underline{S}^\dagger \underline{\Gamma} \underline{S} \right) \underline{\alpha},$$

and complicated things become simple.

Measure the incoming energy

$$E_{\text{in}} = \underline{\alpha}^\dagger \underline{H} \underline{\alpha}$$

The expression

$$\frac{\underline{\alpha}^\dagger \left(\underline{\Gamma} - \underline{S}^\dagger \underline{\Gamma} \underline{S} \right) \underline{\alpha}}{\underline{\alpha}^\dagger \underline{H} \underline{\alpha}}$$

is the transfer of $\langle \Delta \Gamma \rangle$ per Joule of energy of the incoming field

$$\frac{\underline{\alpha}^\dagger (\underline{\Gamma} - \underline{S}^\dagger \underline{\Gamma} \underline{S}) \underline{\alpha}}{\underline{\alpha}^\dagger \underline{H} \underline{\alpha}}$$

Q: **Most efficient beam to transfer Γ to a given object ?**

$$\max_{\underline{\alpha}} \frac{\underline{\alpha}^\dagger (\underline{\Gamma} - \underline{S}^\dagger \underline{\Gamma} \underline{S}) \underline{\alpha}}{\underline{\alpha}^\dagger \underline{H} \underline{\alpha}}$$

Known solution using generalized eigenproblem ($\underline{A}\underline{v} = \lambda\underline{B}\underline{v}$):

- Optimal beam $\underline{\alpha}_{\max}$:
Eigenvector corresponding to maximum generalized eigenvalue of $(\underline{\Gamma} - \underline{S}^\dagger \underline{\Gamma} \underline{S})$ and \underline{H}
- Optimal transfer efficiency: Maximum generalized eigenvalue
- Multipolar decomposition / angular spectrum of $\underline{\alpha}_{\max}$
- Absolute bound for Γ transfer

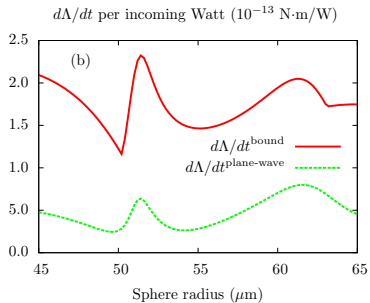
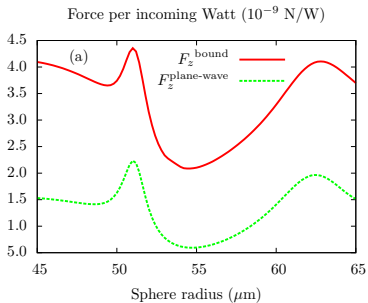
Example with a challenging object:



- CdSe sphere radius $\approx 55 \mu\text{m}$
- 20 helical objects in an icosahedral arrangement

Incoming beams:

- LCP plane wave at $222.8 \mu\text{m}$
 - Object not electromagnetically small
- Optimal monochromatic beam at $222.8 \mu\text{m}$



Also useful for theoretical considerations

$$\langle \Delta \Gamma \rangle = \langle \Phi_{\text{in}} | \Gamma - S^\dagger \Gamma S | \Phi_{\text{in}} \rangle$$

$$\begin{aligned} \langle \Delta \Gamma \rangle &= \langle \Phi_{\text{in}} | \frac{A\Gamma + \Gamma A}{2} | \Phi_{\text{in}} \rangle + \langle \Phi_{\text{in}} | \frac{S^\dagger [S, \Gamma] + [S, \Gamma]^\dagger S}{2} | \Phi_{\text{in}} \rangle \\ &= \langle \Delta \Gamma \rangle_{\text{absorption}} + \langle \Delta \Gamma \rangle_{\text{asymmetry}}. \end{aligned}$$

Symmetry $\iff [S, \Gamma] = 0$: Non-coupling of Γ eigenstates

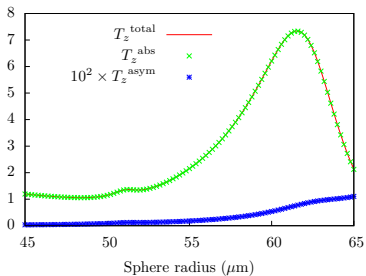
$\langle \Delta \Gamma \rangle_{\text{absorption}}$: $A = I - S^\dagger S$ is the non-unitary part of S

$$\langle \Phi_{\text{in}} | \Phi_{\text{in}} \rangle - \langle \Phi_{\text{out}} | \Phi_{\text{out}} \rangle = \langle \Phi_{\text{in}} | \Phi_{\text{in}} \rangle - \langle \Phi_{\text{in}} | S^\dagger S | \Phi_{\text{in}} \rangle = \langle \Phi_{\text{in}} | I - S^\dagger S | \Phi_{\text{in}} \rangle$$

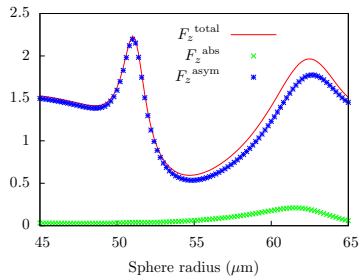
Absorption can be responsible for $\langle \Delta \Gamma \rangle \neq 0$ even if $[S, \Gamma] = 0$



Torque per incoming Watt (10^{-15} N-m/W)



Force per incoming Watt (10^{-9} N/W)



Algebraic approach to light-matter interactions

- Easy to use symmetries and light-matter conservation laws
- Algebraic tools ease the solution of complicated problems
- Algebraic \implies very well suited for computer implementation
- Single Hilbert space \mathbb{M} : Classical fields or single photon
- Should be extensible to multiphoton states of light:

$$\mathbb{M} \otimes \mathbb{M} \otimes \dots$$

Algebraic approach

- [1] I.F-C. Helicity and duality symmetry in light matter interactions: Theory and applications. *Doctoral Thesis, Macquarie University*, 2014.
- [2] I.F-C, and Carsten Rockstuhl. A unified theory to describe and engineer conservation laws in light-matter interactions. *Phys. Rev. A*, Accepted (2017).

Optical Activity

- [1] I.F-C, X. Vidal, N. Tischler, and G. Molina-Terriza. Necessary symmetry conditions for the rotation of light. *J. Chem. Phys.* 138(21), 2013.
- [2] X. Vidal, I.F-C, A. Barbara, and G. Molina-Terriza. Polarization properties of light scattered off solutions of chiral molecules in non-forward direction. *App. Phys. Lett.* 107, 211107 (2015).
- [3] I.F-C, M. Fruhnert, and C. Rockstuhl. Dual and Chiral Objects for Optical Activity in General Scattering Directions. *ACS Photonics*, 2(3), 2015.

Zero backscattering

- [1] I.F-C. Forward and backward helicity scattering coefficients for systems with discrete rotational symmetry. *Optics Express*, 21, Dec. 2013.

Symmetries in Optical Vortices

- [1] I.F-C, X. Zambrana-Puyalto, and G. Molina-Terriza. Helicity and angular momentum: A symmetry-based framework for the study of light-matter interactions. *Phys. Rev. A*, 86, Oct. 2012.
- [2] I.F-C et al. Electromagnetic duality symmetry and helicity conservation for the macroscopic maxwell's equations. *Phys. Rev. Lett.*, 111, Aug. 2013.

I am looking for a tenured(-track) position **Thank you for your time !**