## Beyond infinite time scale separation

Edgeworth approximations for subgrid-scale parameterization

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## The model reduction problem

Many systems of scientific interest are to complex to simulate numerically.

E.g. climate models can resolve only part of the relevant processes of the climate system.

Can a dynamical system of lower dimensionality be determined that approximates the full system?

## Approach: Model reduction through time scale separation

- Assume a time scale separation between slow variables $x$ and fast variables y

$$
\begin{cases}\mathrm{d} x=\frac{1}{\varepsilon} f_{0}(x, y) \mathrm{d} t+f_{1}(x, y) \mathrm{d} t & \\ \text { (resolved/slow/"climate") } \\ \mathrm{d} y=\frac{1}{\varepsilon^{2}} g(x, y) \mathrm{d} t+\frac{1}{\varepsilon} \sigma(x, y) \mathrm{d} W & \\ \text { (unresolved/fast/"weather") }\end{cases}
$$

- As $\varepsilon \rightarrow 0$ the fast $y$ variable decorrelates ever faster and acts as a Gaussian white noise on the slow variables and the slow $x$ variable converges weakly to an SDE.
- This idea can be made mathematically rigourous by the method of homogenization
stochastic: Khasminsky '66, Kurtz '73, Papanicolaou '76
deterministic: Melbourne \& Stuart '11, Gottwald \& Melbourne '13, Melbourne \& Kelly '15, De
Simoi \& Liverani '14


## Slow-fast systems in geophysics



Barotropic vorticity equation with topography

$$
\left\{\begin{array}{l}
\frac{\mathrm{d} U}{\mathrm{~d} t}=\frac{1}{4 \pi^{2}} \int h \frac{\partial \psi}{\partial x} \mathrm{~d} x \mathrm{~d} y \\
\frac{\partial q}{\partial t}+\nabla^{\perp} \psi \cdot \nabla q+U \frac{\partial \psi}{\partial x}+\beta \frac{\partial \psi}{\partial x}=0 \\
q=\Delta \psi+h
\end{array}\right.
$$

The zonal mean flow $U$ evolves slower than the fast Fourier modes $\psi_{i, j}$ of the stream function

This can be modeled by a system with a time scale separation parameter $\varepsilon$

$$
\left\{\begin{array}{l}
\frac{\mathrm{d} U}{\mathrm{~d} t}=\frac{1}{\varepsilon} f_{1}(\psi) \\
\frac{\mathrm{d} \psi \psi_{i, j}}{\mathrm{~d} t}=\frac{1}{\varepsilon^{2}} g_{2}(\psi)+\frac{1}{\varepsilon} g_{1}(U)
\end{array}\right.
$$

Reduces through homogenization, assuming infinite time scale separation to

$$
\mathrm{d} U=\alpha(U) \mathrm{d} t+\sigma(U) \mathrm{d} W
$$

This can be modeled by a system with a time scale separation parameter $\varepsilon$

$$
\begin{cases}\frac{\mathrm{d} U}{\mathrm{~d} t} & =\frac{1}{\varepsilon} f_{1}(\psi) \\ \frac{\mathrm{d} \psi_{\mathrm{i}, j}}{\mathrm{dt}} & =\frac{1}{\varepsilon^{2}} g_{2}(\psi)+\frac{1}{\varepsilon} g_{1}(U)\end{cases}
$$

Reduces through homogenization, assuming infinite time scale separation to

$$
\mathrm{d} U=\alpha(U) \mathrm{d} t+\sigma(U) \mathrm{d} W
$$

Fails when only a moderate time scale separation



## The CLT and the Edgeworth expansion

## The Central Limit Theorem

## Assume $X_{i}$ are i.i.d. random variables

$$
S_{n}:=\frac{1}{\sqrt{n}} \sum_{j=1}^{n}\left(X_{j}-\mu\right) \rightarrow_{d} \mathcal{N}\left(0, \sigma^{2}\right)
$$

where $\mu=\mathbb{E}\left[X_{i}\right]$ and $\sigma^{2}=\mathbb{E}\left[X_{i}^{2}\right]$

For finite $n$, there are deviations to the CLT
These are described by the Edgeworth expansion

$$
\rho_{n}(x)=\Phi_{0, \sigma^{2}}(x) \times\left(1+\frac{1}{6 \sqrt{n}} \frac{\gamma}{\sigma^{3}} H_{3}(x / \sigma)\right)+o(1 / \sqrt{n})
$$

where $H_{3}$ is the third Hermite polynomial and $\gamma=\mathbb{E}\left[X_{i}^{3}\right]$
Feller (1957) "An introduction to probability theory and its applications"

## The CLT and the Edgeworth expansion (dependent version)

The Central Limit Theorem
Assume $X_{i}$ are stationary weakly dependent random variables

$$
S_{n}:=\frac{1}{\sqrt{n}} \sum_{j=1}^{n}\left(X_{j}-\mu\right) \rightarrow_{d} \mathcal{N}\left(0, \sigma^{2}\right)
$$

where $\mu=\mathbb{E}\left[X_{i}\right]$ and $\sigma^{2}=\mathbb{E}\left[X_{1}^{2}\right]+2 \sum_{j=1}^{\infty} \mathbb{E}\left[X_{1} X_{j+1}\right]$

For finite $n$, there are deviations to the CLT
These deviations are described by the Edgeworth expansion

$$
\rho_{n}(x)=\Phi_{0, \sigma^{2}+\delta \sigma^{2} / n}(x) \times\left(1+\frac{1}{\sqrt{n}} \delta \kappa H_{3}(x / \sigma)\right)+o(1 / \sqrt{n})
$$

where $H_{3}$ is the third Hermite polynomial and $\delta \sigma^{2}$ and $\delta \kappa$ are sums over correlation functions of $X_{i}$

Götze and Hipp (1983) Z Wahrscheinlichkeit, 64, 211

## Edgeworth expansion in action: deterministic processes

## Example: deterministic mod process

$$
x_{n}=\frac{1}{\sqrt{n}} \sum_{j=1}^{n} A\left(y_{j}\right) \quad \text { with } \quad \begin{aligned}
y_{j+1} & =p y_{j} \bmod 1 \\
A(y) & =y^{5}+y^{4}+y^{3}+y^{2}+y-c
\end{aligned}
$$

We can calculate $\sigma^{2}, \delta \sigma^{2}$ and $\delta \kappa$ explicitly


Edgeworth $F_{E, n}$ is closer to the truth $F_{n}$ than Gaussian $\mathcal{N}_{0, \sigma_{G K}^{2}}$

$$
n=32, p=3
$$

$$
\left|F_{n}-F_{E, n}\right| \text { is } o(1 / \sqrt{n})
$$

## Edgeworth expansion for slow/fast systems

For slow-fast systems

$$
\left\{\begin{array}{l}
\dot{x}=\frac{1}{\varepsilon} f_{0}(x, y)+f_{1}(x, y) \\
\dot{y}=\frac{1}{\varepsilon^{2}} g_{0}(y)+\frac{1}{\varepsilon} g_{1}(x, y)
\end{array}\right.
$$

we have that $\frac{x(\varepsilon)-x(0)}{\sqrt{\varepsilon}}$ converges to a Gaussian as $\varepsilon \rightarrow 0$.

$$
\rho_{t}\left(x(t) \mid x(0)=x_{0}\right)=\int \mathrm{d} y e^{\mathcal{L} t} \delta_{x_{0}}(x) \mu(\mathrm{d} y)
$$

where $\mathcal{L}=\frac{1}{\varepsilon^{2}} \mathcal{L}_{0}+\frac{1}{\varepsilon} \mathcal{L}_{1}+\mathcal{L}_{2}$ with $\mathcal{L}_{0} \rho=-\partial_{y}(g(y) \rho)$, $\mathcal{L}_{1} \rho=-\partial_{x}\left(f_{0}(x, y) \rho\right)$ and $\mathcal{L}_{2} \rho=-\partial_{x}\left(f_{1}(x, y) \rho\right)$ are generators

Edgeworth corrections to $\rho_{t}$ can be calculated from a Dyson series for the transfer operator

$$
e^{\mathcal{L} t}=e^{\mathcal{L}_{0} t / \varepsilon^{2}}+\int_{0}^{t} d s e^{\mathcal{L}_{0}(t-s) / \varepsilon^{2}}\left(\frac{1}{\varepsilon} \mathcal{L}_{1}+\mathcal{L}_{0}\right) e^{\mathcal{L}_{0} s / \varepsilon^{2}}+\ldots
$$

## Stochastic parameterization using the Edgeworth expansion

Given a slow-fast dynamical system

$$
\left\{\begin{array}{l}
\dot{x}=\frac{1}{\varepsilon} f_{0}(y)+f_{1}(x, y) \\
\dot{y}=\frac{1}{\varepsilon^{2}} g(y)
\end{array}\right.
$$

1. determine the Edgeworth expansion coefficients $\sigma_{G K}^{2}, \delta \kappa$ associated with $f_{0}(x, y)$
2. model $x$ of the multi-scale system by $X$ of a surrogate stochastic process

$$
\left\{\begin{array}{l}
\dot{x}=\frac{1}{\varepsilon} A(\eta)+F(x) \\
\mathrm{d} \eta=-\frac{\gamma}{\varepsilon^{2}} \mathrm{~d} t+\frac{1}{\varepsilon} \mathrm{~d} W
\end{array}\right.
$$

with $A(\eta)=a \eta^{2}+b \eta+c$, where $a, b, c, \gamma$ are determined such that the Edgeworth expansion coefficients of $A(\eta)$ match those of $f_{0}$ in the true system.

## Application: parameterization of a discrete-time multiscale system

$$
\left\{\begin{array}{l}
x_{j+1}^{(\varepsilon)}=x_{j}^{(\varepsilon)}+\varepsilon f_{0}\left(y_{j}\right)+\varepsilon^{2} f_{1}\left(x_{j}^{(\varepsilon)}\right) \\
y_{j+1}=p y_{j} \bmod 1
\end{array}\right.
$$

homogenization: converges for $\varepsilon \rightarrow 0$ to a diffusion (Gottwald \&
Melbourne (2013))

$$
\mathrm{d} X=f_{1}(X) \mathrm{d} t+\sigma_{G K} \mathrm{~d} W
$$

Edgeworth: replace fast mod map $y$ by an AR1 process $\eta$

$$
\left\{\begin{array}{l}
x_{j+1}^{(\varepsilon)}=x_{j}^{(\varepsilon)}+\varepsilon f_{0, s}\left(\eta_{j}\right)+\varepsilon^{2} f_{1}\left(x_{j}^{(\varepsilon)}\right) \\
\eta_{j+1}=\phi \eta_{j}+N_{j}
\end{array}\right.
$$

with $f_{0, s}(\eta)=a_{3} \eta^{3}+a_{2} \eta^{2}+a_{1} \eta+a_{0}$ and parameters $a_{i}$ tuned to match Edgeworth corrections

mean


variance

## Parameterization of a continuous-time multiscale system

$$
\left\{\begin{array}{l}
\dot{x}=\frac{1}{\varepsilon} f_{0}(y)+f_{1}(x) \\
\dot{y}=\frac{1}{\varepsilon^{2}} g(y)
\end{array}\right.
$$

where $f_{1}(x)=-\nabla V(x)$, with $V(X)$ an assymetric double well potential and $\dot{y}=g(y)$ is the standard Lorenz '63 system.

Edgeworth: replace fast Lorenz system y by an Ornstein-Uhlenbeck process $\eta$

$$
\left\{\begin{array}{l}
\dot{X}=\frac{1}{\varepsilon} f_{0, s}(\eta)+f_{1}(X) \\
\mathrm{d} \eta=-\frac{1}{\varepsilon^{2}} \gamma \eta \mathrm{~d} t+\frac{1}{\varepsilon} \mathrm{~d} W
\end{array}\right.
$$

with $f_{0, s}(\eta)=a_{3} \eta^{3}+a_{2} \eta^{2}+a_{1} \eta+a_{0}$ and parameters $a_{i}$ tuned to match Edgeworth corrections


## mean


third moment

## Summary

- We have used the Edgeworth expansion to extend the range of time scale separation over which slow-fast sytems can be approximated
- The fast variables are replaced by a stochastic surrogate process, the parameters of which are tuned to match the Edgeworth expansion
- We have shown good agreement when reducing deterministic discrete and continuous time systems
- To do: Apply Edgeworth based reduction to the barotropic vorticity equation


## Summary

- We have used the Edgeworth expansion to extend the range of time scale separation over which slow-fast sytems can be approximated
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- We have shown good agreement when reducing deterministic discrete and continuous time systems
- To do: Apply Edgeworth based reduction to the barotropic vorticity equation

Thank you for your attention!

## homogenization extends the Central Limit Theorem

$$
\left\{\begin{aligned}
\mathrm{d} x^{(\varepsilon)} & =\frac{1}{\varepsilon} f_{0}\left(y^{(\varepsilon)}\right) \mathrm{d} t & & \text { resolved/slow } \\
\mathrm{d} y(\varepsilon) & =\frac{1}{\varepsilon^{2}} g\left(y^{(\varepsilon)}\right) \mathrm{d} t+\frac{1}{\varepsilon} \sigma \mathrm{~d} W & & \text { unresolved/fast }
\end{aligned}\right.
$$

The slow variable $x$ integrates the fast variable $y$

$$
\begin{aligned}
x^{(\varepsilon)}(t)-x^{(\varepsilon)}(0) & =\frac{1}{\varepsilon} \int_{0}^{t} f_{0}\left(y^{(\varepsilon)}(s)\right) \mathrm{d} s \\
& =\varepsilon \int_{0}^{t / \varepsilon^{2}} f_{0}\left(y^{(\varepsilon=1)}(s)\right) \mathrm{d} s \\
& =\frac{1}{\sqrt{n}} \int_{0}^{t n} f_{0}\left(y^{(\varepsilon=1)}(s)\right) \mathrm{d} s
\end{aligned}
$$

Invoking the CLT, $x(t)$ converges weakly to $\mathrm{d} X=\sigma \mathrm{d} W$ where
$\sigma^{2}=2 \int_{0}^{\infty} \mathbb{E}\left[f_{0}\left(y^{(1)}(0)\right) f_{0}\left(y^{(1)}(s)\right)\right] d s$

## homogenization

$$
\begin{cases}\mathrm{d} x=\frac{1}{\varepsilon} f_{0}(x, y) \mathrm{d} t+f_{1}(x, y) \mathrm{d} t & \text { resolved/slow } \\ \mathrm{d} y=\frac{1}{\varepsilon^{2}} g(x, y) \mathrm{d} t+\frac{1}{\varepsilon} \sigma(x, y) \mathrm{d} W & \text { unresolved/fast }\end{cases}
$$

Assumptions:

- fast $y$-process is ergodic with measure $\mu_{x}$
- $\int f_{0}(x, y) \mathrm{d} \mu_{x}=0$

In the limit $\varepsilon \rightarrow 0$, the slow $x$-dynamics is approximated by

$$
\mathrm{d} X=F(X) \mathrm{d} t+\Sigma(X) \mathrm{d} W
$$

where

$$
\begin{aligned}
& \Sigma \Sigma^{T}=2 \int_{0}^{\infty} \mathbb{E}^{\mu_{x}}\left[f_{0}(x, y) f_{0}(x, y(s))\right] \mathrm{d} s \\
& F(X)=\int f_{1}(x, y) \mathrm{d} \mu_{x}+\int_{0}^{\infty} \int \nabla_{x} f_{0}(x, y(s)) f_{0}(x, y) \mathrm{d} \mu_{x} \mathrm{~d} s
\end{aligned}
$$

stochastic: Khasminsky '66, Kurtz '73, Papanicolaou '76
deterministic: Melbourne \& Stuart '11, Gottwald \& Melbourne '13, Melbourne \& Kelly '15

## "Proof" of the Edgeworth expansion

Expand the characteristic function of $X / \sqrt{n}$ (assuming $\mu=0, \sigma=1$ ):

$$
\begin{aligned}
\mathbb{E}\left[e^{\imath t X / \sqrt{n}}\right] & =\mathbb{E}\left[1+\frac{\imath t X}{\sqrt{n}}+\frac{(\imath t)^{2} X^{2}}{2 n}+\frac{(\imath t)^{3} X^{3}}{6 n \sqrt{n}}+\ldots\right] \\
& =\left(1-\frac{t^{2}}{2 n}\right)+\frac{(\imath t)^{3}}{6 n \sqrt{n}} \mathbb{E}\left[X^{3}\right]+\ldots
\end{aligned}
$$

The characteristic function of $\sum_{j=1}^{n} X_{j} / \sqrt{n}$

$$
\begin{aligned}
\mathbb{E}\left[e^{\imath t X / \sqrt{n}}\right]^{n} & =\left(1-\frac{t^{2}}{2 n}\right)^{n}+\left(1-\frac{t^{2}}{2 n}\right)^{(n-1)} \frac{(\imath t)^{3} \gamma}{6 \sqrt{n}}+\ldots \\
& =e^{-t^{2} / 2}\left(1+\frac{(\imath t)^{3} \gamma}{6 \sqrt{n}}\right)+O\left(\frac{1}{n}\right)
\end{aligned}
$$

## Application: stochastic approximation of a deterministic map

Replace $\left\{\begin{array}{l}x_{j+1}=x_{j}+\varepsilon A\left(y_{j}\right) \\ y_{j+1}=p y_{j} \bmod 1\end{array}\right.$ with $A(y)=y^{5}+y^{4}+y^{3}+y^{2}+y-c$ by a surrogate AR1 process

$$
\left\{\begin{array}{l}
x_{j+1}=x_{j}+\varepsilon B\left(\eta_{j}\right) \\
\eta_{j+1}=\phi \eta_{j}+N_{j}
\end{array} \text { with } B(y)=a_{s} \eta^{2}+b_{s} \eta+c_{s}\right.
$$

such that $\sigma_{G K}^{2}$ (homogenization), as well as $\delta \kappa_{3}$ (1st Edgeworth term) match



## Edgeworth expansion in action: stochastic processes

## Example 1: AR1 process

$$
\left\{\begin{array} { l } 
{ x _ { j + 1 } = x _ { j } + \varepsilon A ( \eta _ { j } ) } \\
{ \eta _ { j + 1 } = \phi \eta _ { j } + N _ { j } }
\end{array} \quad \text { with } \quad \left\{\begin{array}{l}
A(\eta)=a \eta^{2}+b \eta+c \\
N_{j} \sim \mathcal{N}(0,1)
\end{array}\right.\right.
$$

We can calculate $\sigma^{2}, \delta \sigma^{2}$ and $\delta \kappa$ explicitly (everything is Gaussian)


Edgeworth $F_{E, n}$ is closer to the truth $F_{n}$ than Gaussian $\mathcal{N}_{0, \sigma_{G K}^{2}}$

$\left|F_{n}-F_{E, n}\right|$ is $o(1 / \sqrt{n})$

$$
n=32, \phi=1 / 3, a=b=1
$$

Triad system Majda et al (2001)

$$
\begin{cases}\dot{x} & =\frac{1}{\varepsilon} B_{0} y_{1} y_{2} \\ \dot{y}_{1} & =\frac{1}{\varepsilon} B_{1} y_{2} x-\frac{1}{\varepsilon^{2}} \gamma_{1} y_{1}-\frac{1}{\varepsilon} \sigma_{1} \dot{W}_{1} \\ \dot{y}_{2} & =\frac{1}{\varepsilon} B_{2} x y_{1}-\frac{1}{\varepsilon^{2}} \gamma_{2} y_{2}-\frac{1}{\varepsilon} \sigma_{2} \dot{W}_{2}\end{cases}
$$

Edgeworth:

$$
\begin{aligned}
& \left\{\begin{array}{l}
\dot{X}=\frac{1}{\varepsilon} A(\eta) \\
\dot{\eta}=\frac{1}{\varepsilon} \alpha X-\frac{1}{\varepsilon^{2}} \eta-\frac{1}{\varepsilon} \sigma \dot{W}
\end{array}\right. \\
& \text { with } A(\eta)=a_{s} \eta^{2}+b_{s} \eta+c_{s}
\end{aligned}
$$



