

Convergence in time of discrete evolutions generated by alternate minimizing schemes

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Phase-field energy

Energy functional

$$\mathcal{F}(t, u, v) = \mathcal{E}(t, u, v) + G_c \mathcal{L}(v)$$

$$\mathcal{U} = \{u \in H_0^1(\Omega)\} \quad \mathcal{V} = \{v \in H^1(\Omega), 0 \leq v \leq 1\}$$

Linear elastic energy $\mathcal{E}(t, u, v) = \frac{1}{2} \int_{\Omega} (v^2 + \eta) W(\epsilon) dx - \int_{\Omega} f(t) \cdot u dx$

Dissipated energy $\mathcal{L}(v) = \frac{1}{2} \int_{\Omega} (v - 1)^2 + |\nabla v|^2 dx$

[Ambrosio-Tortorelli (90), Bourdin-Francfort-Marigo (00), etc.]

$\mathcal{F}(t, \cdot, \cdot)$ is non-convex but separately quadratic

Which quasi-static evolution ?

- i)* alternate minimization
- ii)* vanishing viscosity

Time discrete evolution by alternate minimization

Time discretization $t_k = k\Delta t$. Set $u(t_0)$ and $v(t_0)$

Define by recursion, $u_m = u(t_{k-1})$ and $v_m = v(t_{k-1})$ for $m = 0$

$$\text{for } m \geq 1 \quad \begin{cases} u_m \in \operatorname{argmin} \{ \mathcal{F}(t_k, \cdot, v_{m-1}) \text{ in } \mathcal{U} \} \\ v_m \in \operatorname{argmin} \{ \mathcal{F}(t_k, u_m, \cdot) \text{ in } \mathcal{V} \text{ with } v \leq v_{m-1} \} \end{cases}$$

minimization of quadratic functionals

Let $u(t_k) = \lim_{m \rightarrow +\infty} u_m$ and $v(t_k) = \lim_{m \rightarrow +\infty} v_m$

long/short increment

Then $(u(t_k), v(t_k))$ is equilibrium point for $\mathcal{F}(t_k, \cdot, \cdot)$ and separate minimizer

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Piecewise-affine (stair like) interpolation $u_{\Delta t}$ and $v_{\Delta t}$

For $\Delta t \rightarrow 0$. Do $u_{\Delta t}$ and $v_{\Delta t}$ converge and, if so, to which evolution?

Arc-length parametrization of the “stair-like” path w.r.t. $\|\cdot\|_u$ and $\|\cdot\|_v$

$$[0, S] \ni s \mapsto (t_{\Delta t}(s), u_{\Delta t}(s), v_{\Delta t}(s)) \in [0, T] \times H_0^1 \times H^1$$

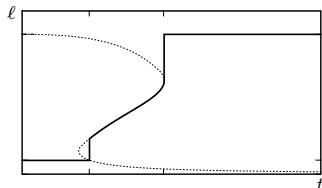
Then $(t_{\Delta t}, u_{\Delta t}, v_{\Delta t})$ is bounded in $W^{1,\infty}([0, S]; [0, T] \times H_0^1 \times H^1)$

There exists a limit evolution for $\Delta t \rightarrow 0$.

Properties and characterization?

A simple sharp crack example by Griffith's criterion

cf. "arc-length methods" [Verhoosel-Remmers-Gutierrez (09)]



dotted: critical $G(t, \ell) = G_c$
stable $G(t, \ell) < G_c$ (left)
unstable $G(t, \ell) > G_c$ (right)

bold: evolution $\ell(t)$

For $t'(s) = 0$ unstable propagation (discontinuity point in time) [long increment]

For $t'(s) > 0$ stable or critical propagation (continuity point in time) [short increment]

Continuity points: equilibrium

Equilibrium for $t'(s) > 0$.

In terms of energy variations:

[Knees-N.]

$$\begin{cases} \partial_u \mathcal{F}(t(s), u(s), v(s))[\phi] = 0 & \text{for all } \phi \\ \partial_v \mathcal{F}(t(s), u(s), v(s))[\xi] \geq 0 & \text{for all } \xi \leq 0 \end{cases}$$

Equivalently $u(s)$ and $v(s)$ are separate minimizers.

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In terms of PDEs:

$$\begin{cases} -\operatorname{div}(\sigma_{v(s)}(u(s))) = f(s) & \text{for } \sigma_v(u) = (v^2 + \eta) \sigma(u) \\ [v(s)W(Du(s)) + (v(s) - 1) - \Delta v(s)]_+ = 0 \end{cases}$$

Technically [...] is a negative finite Radon measure.

Not known if $v \in H^2$ cf. [Akagi-Kimura]

Phase field energy release?

Simultaneous evolution.

$$\tilde{\mathcal{E}}(t, v) = \mathcal{E}(t, u_{t,v}, v) \quad \text{for} \quad u_{t,v} \in \operatorname{argmin} \{ \mathcal{E}(t, \cdot, v) : u \in \mathcal{U} \}.$$

$$\partial_v \tilde{\mathcal{E}}(t, v)[\xi] = \int_{\Omega} v \xi W(\boldsymbol{\varepsilon}_{t,v}) dx = \partial_v \mathcal{E}(t, u_{t,v}, v)[\xi]$$

Normalized admissible variations $\hat{\Xi} = \{ \xi \in H^1 : \xi \leq 0 \text{ and } d\mathcal{L}(v)[\xi] = 1 \}$

$$\mathcal{G}(t, v) = \inf \{ -\partial_v \tilde{\mathcal{E}}(t, v)[\xi] : \xi \in \hat{\Xi} \} \geq 0$$

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Griffith's criterion in KT form. If $t'(s) > 0$ then $d\mathcal{L}'(v(s)) \geq 0$

- $\mathcal{G}(t(s), v(s)) \leq G_c$
- $(\mathcal{G}(t(s), v(s)) - G_c) \mathcal{L}'(v(s)) = 0$

Behind simultaneous (u, v) evolution

A couple of quantitative estimates

- If $u \in \operatorname{argmin} \{ \mathcal{E}(t, \cdot, v) : u \in \mathcal{U} \}$ then $u \in W^{1,p}(\Omega)$ for $p \gtrsim 2$ and $\Omega \subset \mathbb{R}^2$

$$\|u_1 - u_2\|_{W^{1,p}} \leq C \|v_0 - v_1\|_{L^q} \leq C' \|v_0 - v_1\|_{H^1} \quad \text{for } q \gg 2$$

[Herzog-Meyer-Wachsmuth (11) cf. also Knees-Rossi-Zanini (13)]

- If $v_k \in \operatorname{argmin} \{ \mathcal{F}(t, u_k, \cdot) : v \leq v_{k-1} \}$ then

$$\|v_1 - v_2\|_{H^1} \leq C \|u_1 - u_2\|_{W^{1,p}} \|u_1 + u_2\|_{W^{1,p}}$$

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Alternate minimization as a "simultaneous quasi-Newton" method

$$-\begin{pmatrix} \partial_v \mathcal{F}(\dots)_{k-1} \\ \partial_u \mathcal{F}(\dots)_{k-1} \end{pmatrix} = \left(\begin{array}{c|c} \partial_{vv}^2 \mathcal{F}(\dots)_{k-1} & 0 \\ \hline 0 & \partial_{uu}^2 \mathcal{F}(\dots)_{k-1} \end{array} \right) \begin{pmatrix} v_k - v_{k-1} \\ u_k - u_{k-1} \end{pmatrix}$$

Discontinuity points: a "normalized gradient flow"

A normalized gradient flow if $t'(s) = 0$

Steepest descent (normalized gradient flow): for a.e. $s \in [0, S]$ s.t. $t'(s) = 0$

$$\begin{cases} \hat{u}'(s) \in \operatorname{argmin} \{ \partial_u \mathcal{F}(t(s), u(s), v(s))[\phi] : \|\phi\|_{v(s)} \leq 1 \} \\ \hat{v}'(s) \in \operatorname{argmin} \{ \partial_v \mathcal{F}(t(s), u(s), v(s))[\xi] : \|\xi\|_{u(s)} \leq 1, \xi \leq 0 \} \end{cases}$$

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Which PDEs? Up to a normalization factor

$$\begin{cases} -\operatorname{div} (\sigma_{v(s)}(u(s) + u'(s))) = f(s) & \text{[viscoelastic phase-field flow]} \\ \dots & \text{[variational inequality or a differential inclusion]} \end{cases}$$

Write the energy

$$\begin{aligned}\mathcal{F}(t, u, v) &= \frac{1}{2} \int_{\Omega} (v^2 + \eta) W(\varepsilon) dx - \int_{\Omega} f(t) \cdot u dx + G_c \mathcal{L}(v) \\ &= \frac{1}{2} \|u\|_v^2 + \beta_t(u) + c_v\end{aligned}$$

Note that $\|u\|_v$ and $\|u\|_{H_0^1}$ are equivalent

Hence $\partial_u \mathcal{F}(t, u, v)[\phi] = \langle u, \phi \rangle_v + \beta_t(\phi)$ and the slope is

$$|\partial_u \mathcal{F}(t, u, v)|_v = \max \{ -\partial_u \mathcal{F}(t, u, v)[\phi] : \|\phi\|_v \leq 1 \}$$

Similarly define the norm $\|v\|_u$ and the unilateral slope:

$$|\partial_v^- \mathcal{F}(t, u, v)|_u = \max \{ -\partial_v \mathcal{F}(t, u, v)[\xi] : \xi \leq 0, \|\xi\|_u \leq 1 \}$$

Technically: a parametrized BV evolution for $\Delta t \rightarrow 0$

A Lipschitz parametrization $s \mapsto (t(s), u(s), v(s))$ with [Knees-N. (...)]

$$t'(s) \geq 0 \quad v'(s) \leq 0 \quad 0 < t'(s) + \|u'(s)\|_{v(s)} + \|v'(s)\|_{u(s)} \leq 1$$

• for every s with $t'(s) > 0$ equilibrium

$$|\partial_u \mathcal{F}(t(s), u(s), v(s))|_{v(s)} = |\partial_v^- \mathcal{F}(t(s), u(s), v(s))|_{u(s)} = 0$$

• for every s energy balance

$$\begin{aligned} \mathcal{F}(t(s), u(s), v(s)) &= \mathcal{F}(0, u_0, v_0) + \int_0^s \mathcal{P}_{ext}(t(r), u(r), v(r)) t'(r) dr \\ &\quad - \int_0^s |\partial_u \mathcal{F}(t(r), u(r), v(r))|_{v(r)} \|u'(r)\|_{v(r)} dr \\ &\quad - \int_0^s |\partial_v^- \mathcal{F}(t(r), u(r), v(r))|_{u(r)} \|v'(r)\|_{u(r)} dr \end{aligned}$$

Proof from [N. (14)] see also [Efendiev-Mielke (06), Sandier-Serfaty (04)]

Stopping criterion

Incremental problem: $u_m = u(t_{k-1})$ and $v_m = v(t_{k-1})$ for $m = 0$

$$\text{for } m \geq 1 \quad \begin{cases} u_m \in \operatorname{argmin} \{ \mathcal{F}(t_k, \bullet, v_{m-1}) \text{ in } \mathcal{U} \} \\ v_m \in \operatorname{argmin} \{ \mathcal{F}(t_k, u_m, \bullet) \text{ in } \mathcal{V} \text{ with } v \leq v_{m-1} \} \end{cases}$$

Set $u(t_k) = u_{\bar{m}}$ and $v(t_k) = v_{\bar{m}}$ for some $\bar{m} \geq 1$ according to ...

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Set $u(t_k) = u_{\bar{m}}$ and $v(t_k) = v_{\bar{m}}$ for some $\bar{m} \geq 1$ according to ...

Piecewise-affine interpolation $u_{\Delta t}$ and $v_{\Delta t}$

$u_{\Delta t}$ and $v_{\Delta t}$ converge, as $\Delta t \rightarrow 0$, to a parametrized BV evolution

No information on uniqueness and speed of convergence.

Irreversibility constraint and time interpolation

Incremental problem: $u_m = u(t_{k-1})$ and $v_m = v(t_{k-1})$ for $m = 0$

$$\text{for } m \geq 1 \quad \begin{cases} u_m \in \operatorname{argmin} \{ \mathcal{F}(t_k, \bullet, v_{m-1}) \text{ in } \mathcal{U} \} \\ v_m \in \operatorname{argmin} \{ \mathcal{F}(t_k, u_m, \bullet) \text{ in } \mathcal{V} \text{ with } v \leq v(t_{k-1}) \} \end{cases}$$

Set $u(t_k) = \lim_m u_m$ and $v(t_k) = \lim_m v_m$ [Bourdin (07)] [Burke-Ortner-Süli (10)]

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Piecewise-affine interpolation of $u(t_k)$ and $v(t_k)$ in the points t_k .

Then $u_{\Delta t}$ and $v_{\Delta t}$ converge to u and v s.t. $v \in BV(0, T; L^1)$ with $\dot{v} \leq 0$

$$\mathcal{F}(t, u(t), v(t)) \leq \mathcal{F}(0, u_0, v_0) + \int_0^t \partial_t \mathcal{F}(r, u(r), v(r)) dr$$

[local q.s. evolutions]

The transition v^- and v^+ in discontinuity points is not fully characterized.

A “Ginzburg-Landau” system

Evolution governed by a system of PDEs

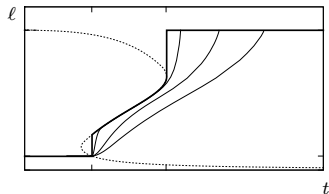
[Kuhn-Müller (10), N.]

$$\begin{cases} \varepsilon \dot{v}(t) = -[v(t)W(Du(t)) + (v(t) - 1) - \Delta v(t)]_+ \\ -\operatorname{div}(\sigma_{v(t)}(u(t))) = f(t) \end{cases}$$

We study:

- existence of solutions
(by alternate minimizing movement)
- quasi-static limit for $\varepsilon \rightarrow 0$
(by parametrized *BV*-evolutions)

[Dal Maso-Toader (03)] [Mielke et. al.]



Minimizing movements (implicit Euler)

Let $t_k = k\Delta t$. We employ an alternate (1-step) minimizing movement:

[N.] cf. [Babadjain-Millot (14)]

$$\begin{cases} u_k \in \operatorname{argmin} \{ \mathcal{F}(t_k, u, v_{k-1}) \}, \\ v_k \in \operatorname{argmin} \{ \mathcal{F}(t_k, u_k, v) + \frac{\varepsilon}{2\Delta t} \|v - v_{k-1}\|_{L^2}^2 : v \leq v_{k-1} \} \end{cases}$$

minimization of quadratic functionals and only "short" increments

Euler-Lagrange "equations"

$$\partial_u \mathcal{F}(t_k, u_k, v_{k-1}) = 0 \quad \text{"} \partial_v \mathcal{F}(t_k, u_k, v_k) + \varepsilon \dot{v}_k = 0 \text{"}$$

Piecewise affine interpolate $v_{\Delta t}$ and $u_{\Delta t}$ in the points t_k

Prove that $v_{\Delta t}$ and $u_{\Delta t}$ converge to a solution of the system of PDEs.

1) Compactness

$$v_{\Delta t} \in W^{1,2}(0, T; L^2) \cap L^\infty(0, T; H^1) \quad u_{\Delta t} \in L^\infty(0, T; W^{1,p}) \text{ for } p > 2$$

2) By l.s.c. of energy and slope and by a "upper gradient inequality"

separate convexity + argument by [Dal Maso-Francfort-Toader (05)]

$$\begin{aligned} \mathcal{F}(t, u(t), v(t)) &= \mathcal{F}(0, u_0, v_0) + \int_0^t \partial_t \mathcal{F}(r, u(r), v(r)) dr + \\ &\quad - \frac{1}{2} \int_0^t \|\varepsilon \dot{v}(r)\|_{L^2}^2 + |\partial_v^- \mathcal{F}(r, u(r), v(r))|_{L^2}^2 dr. \end{aligned}$$

(De Giorgi's representation of "gradient flows")

3) From the discrete Euler-Lagrange "equations"

$$\begin{cases} \varepsilon \dot{v}(t) = -[v(t) W(Du(t)) + (v(t) - 1) - \Delta v(t)]_+ \\ -\operatorname{div}(\sigma_{v(t)}(u(t))) = f(t) \end{cases}$$

[..] is a finite Radon measure with [..]₊ in L^2

[N.] cf. [Gianazza-Savaré (94)]

No chain rule and no compactness in $W^{1,2}(0, T; H^1)$ are needed

- 1) Arc-length $t \mapsto s(t) = \int_0^t 1 + \|\dot{v}_\varepsilon(r)\|_{L^2} dr$ with inverse $t_\varepsilon(s)$
- parametrized evolutions $s \mapsto (t_\varepsilon(s), v_\varepsilon(s))$ for $s \in [0, S_\varepsilon]$

- 2) Delicate technical point:

- length S is finite unif. w.r.t. ε

Discrete Gronwall estimate [Nocetto-Savaré-Verdi (00)] cf. [Knees-Rossi-Zanini (13)]

- compactness of v_ε in $W^{1,\infty}(0, S; L^2)$

- 3) Equilibrium and parametrized energy balance for the limit evolution

$$\begin{aligned} \mathcal{F}(t(s), u(s), v(s)) &= \mathcal{F}(0, u_0, v_0) + \int_0^s \partial_t \mathcal{F}(t(r), u(r), v(r)) t'(r) dr + \\ &\quad - \int_0^s |\partial_v^- \mathcal{F}(t(r), u(r), v(r))|_{L^2} dr \end{aligned}$$

From energy balance to PDEs

From the energy balance (not from the PDE) get "two regimes"

- If $t'(s) > 0$ (continuity points)

$$\begin{cases} [v(s)W(Du(s)) + (v(s) - 1) - \Delta v(s)]_+ = 0 \\ -\operatorname{div}(\sigma_{v(s)}(u(s))) = f(t) \end{cases}$$

- If $t'(s) = 0$ (discontinuity points)

$$\begin{cases} \hat{v}'(s) = -[v(s)W(Du(s)) + (v(s) - 1) - \Delta v(s)]_+ \\ -\operatorname{div}(\sigma_{v(s)}(u(s))) = f(t) \end{cases}$$

Employ the chain rule and local compactness in $W^{1,2}(\dots, H^1)$

Not same as the q.s. alternate minimization (underlying norms are different)

From phase-field to sharp crack (I)

Γ -convergence of the energies

$$\mathcal{F}_\delta(t, u, v) = \int_{\Omega} (v^2 + \eta_\delta) W(\varepsilon) dx + G_c \int_{\Omega} v^2 / (4\delta) + \delta |\nabla v|^2 dx$$

↓

$$\mathcal{F}_0(t, u) = \int_{\Omega \setminus J_u} W(\varepsilon) dx + G_c \mathcal{H}^{n-1}(J_u)$$

(convergence holds in a very general setting)

[Ambrosio-Tortorelli (90)] [Chambolle (03)] [Iurlano (12)]

Γ -convergence implies convergence of global minimizers (upon compactness)

What about critical points, slopes and energy release?

[Francfort-Le-Serfaty (09)]

From phase-field to sharp crack (II)

Convergence of energies and derivatives imply convergence of evolutions

for gradient flows [Sandier-Serfaty (04)] for BV-evolutions [N. (14)]

At least we should have ...

$$\checkmark \mathcal{F}_0(t, u) \leq \liminf_{\delta \rightarrow 0} \mathcal{F}_\delta(t, u_\delta, v_\delta)$$

$$? \mathcal{G}_0(t, J_u) \leq \liminf_{\delta \rightarrow 0} \mathcal{G}_\delta(t, v_\delta) \text{ or for a slope ...}$$

Some partial results

- $G(t, \ell) = \lim_{\delta \rightarrow 0} \mathcal{G}_\delta(t, v_\delta)$ in a simplified geometry [Sicsic-Marigo (13)] [N. (13)]
- a different notion of evolution [Babadjian-Millot (14)]