Specialized Macdonald polynomials, quantum *K*-theory, and Kirillov-Reshetikhin crystals

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Recursive construction procedure (for the non-symmetric ones $E_{\mu}(x; q, t)$), based on Cherednik's intertwiners I_i .

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Theorem (Braverman-Finkelberg, Ion) *We have*

$$P_{\lambda}(x;q,t=0)=\widehat{\Psi}_{\lambda}(x;q).$$

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 $H^*(G/B)$ and K(G/B) have bases of Schubert classes; for *K*-theory, they are the classes $[\mathcal{O}_w] = [\mathcal{O}_{X_w}]$ of structure sheaves of X_w .

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A k-point GW invariant (of degree d) counts curves of degree d passing through k given Schubert varieties.

Givental and Lee defined *K*-theoretic GW invariants by applying the *K*-theory Euler characteristic when the space of curves (through given Schubert varieties) is infinite.

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The *K*-theoretic *J*-function is the generating function of 1-point K-theoretic GW invariants.

Theorem (Braverman-Finkelberg)

In simply-laced types, the q-Whittaker function $\Psi_{\lambda}(x; q)$ (viewed as a function of λ) coincides with the K-theoretic J-function.

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Let $\mathbf{p} = (p_1, p_2, \ldots)$ be a composition, and

 $W^{\otimes \mathbf{p}} = W^{\mathbf{p}_1,1} \otimes W^{\mathbf{p}_2,1} \otimes \ldots, \quad \lambda = \omega_{\mathbf{p}_1} + \omega_{\mathbf{p}_2} + \ldots.$

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Main Theorem (L.-Naito-Sagaki-Schilling-Shimozono) For all untwisted affine root systems $A_{n-1}^{(1)} - G_2^{(1)}$, we have

 $P_{\lambda}(x;q,0)=X_{\lambda}(x;q).$

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Remarks. (1) The result is believed to extend to the twisted types. (2) In simply-laced types, certain affine Demazure characters were identified with $P_{\lambda}(x; q, 0)$ (lon), and $X_{\lambda}(x; q)$ (Fourier-Littelmann).

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- ► the specialized Macdonald polynomials P_λ(x; q, 0) and the q-Whittaker functions (Ram-Yip formula),
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The model is uniform for all Lie types $A_{n-1} - G_2$.

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Reflections s_{α} , $\alpha \in \Phi$.



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The quantum Bruhat graph QBG(W) on W is the directed graph with labeled edges

$$w \xrightarrow{\alpha} ws_{\alpha}$$
,

where

$$\begin{split} \ell(\mathit{ws}_{\alpha}) &= \ell(\mathit{w}) + 1 \quad (\text{covers of the Bruhat order}), \quad \text{or} \\ \ell(\mathit{ws}_{\alpha}) &= \ell(\mathit{w}) - 2\mathrm{ht}(\alpha^{\vee}) + 1 \qquad (\mathrm{ht}(\alpha^{\vee}) = \langle \rho, \alpha^{\vee} \rangle) \,. \end{split}$$

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$$\begin{split} \ell(\mathit{ws}_{\alpha}) &= \ell(\mathit{w}) + 1 \quad (\text{covers of the Bruhat order}), \quad \text{or} \\ \ell(\mathit{ws}_{\alpha}) &= \ell(\mathit{w}) - 2\mathrm{ht}(\alpha^{\vee}) + 1 \qquad (\mathrm{ht}(\alpha^{\vee}) = \langle \rho, \alpha^{\vee} \rangle) \,. \end{split}$$

Comes from the multiplication of Schubert classes in the quantum cohomology of flag varieties $QH^*(G/B)$ (Fulton and Woodward).

Bruhat graph for S_3 :



Quantum Bruhat graph for S_3 :



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Given a dominant weight λ , we associate with it a sequence of roots, called a λ -chain:

$$\Gamma = (\beta_1, \ldots, \beta_m).$$

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Fact. The construction of a λ -chain is based on a reduced decomposition of the affine Weyl group element corresponding to $A_{\circ} - \lambda$. This gives a sequence of alcoves from A_{\circ} to $A_{\circ} - \lambda$.

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Important structures:

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Let $\mathcal{A}_q(\Gamma) := \mathcal{A}_q(\Gamma, 1_W)$ and $\mathcal{A}_{\sphericalangle}(\Gamma) := \mathcal{A}_{\sphericalangle}(\Gamma, 1_W)$.

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Theorem (Ram-Yip, L.)

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Remark. For q = 0, we retrieve the alcove model (L. and Postnikov, cf. Gaussent and Littelmann, Littelmann):

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Let Γ_{rev} = reverse of an ω_k -chain (ω_k a fundamental weight).

Theorem (L.-Postnikov, L.-Shimozono) In K(G/B) (finite-type or Kac-Moody), we have

$$[\mathcal{O}_w] \cdot [\mathcal{O}_{s_k}] = \sum_{J \in \mathcal{A}_{\triangleleft}(\Gamma_{\mathrm{rev}}, w) \setminus \{\emptyset\}} (-1)^{|J|-1} [\mathcal{O}_{\mathrm{end}(w, J)}].$$

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Conjecture (L.-Postnikov) In QK(G/B) (finite-type), we have:

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Remark. Restricting the RHS, we retrieve the Chevalley formula in $QH^*(G/B)$ (Fulton-Woodward).

• Computer experiments (A. Buch).

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- They multiply as in the conjectured Chevalley formula.
- They are conjectured to represent Schubert classes $[\mathcal{O}_w]$ in $QK(SL_n/B)$.

Recall the KR modules, as modules for $U_q(\hat{\mathfrak{g}})$: $W^{r,s}$ and

$$W^{\otimes \mathbf{p}} = W^{p_1,1} \otimes W^{p_2,1} \otimes \ldots$$

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Fact. $W^{\otimes p}$ has a basis (crystal basis) $B = B^{\otimes p}$ such that in the limit $q \to 0$ we have

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Fact. The crystal structure on $B^{\otimes \mathbf{p}}$ is defined by a tensor product rule: $B^{\otimes \mathbf{p}} = B^{p_1,1} \otimes B^{p_2,1} \otimes \dots$

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Goal. Uniform model for all types $A_{n-1}^{(1)} - G_2^{(1)}$, based on the quantum alcove model.

The quantum alcove model for KR crystals

Given $\mathbf{p} = (p_1, p_2, \ldots)$ and an arbitrary Lie type, let $\lambda = \omega_{p_1} + \omega_{p_2} + \ldots$.

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Construction. (L. and Lubovsky, generalization of L.-Postnikov, Gaussent-Littelmann) *Crystal operators* $\tilde{f}_1, \ldots, \tilde{f}_r$ and \tilde{f}_0 on $\mathcal{A}_q(\Gamma)$.
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Main Theorem (L.-Naito-Sagaki-Schilling-Shimozono) The (combinatorial) crystal $A_q(\Gamma)$ is isomorphic to the tensor product of KR crystals $B^{\otimes p}$.

It originates in the theory of exactly solvable lattice models.

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More precisely, D_B : $B \to \mathbb{Z}_{\geq 0}$ satisfies the following conditions:

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Goal. A more efficient uniform calculation, based only on the combinatorial data associated with a crystal vertex (type *A*: Lascoux–Schützenberger charge statistic).

The energy via the quantum alcove model

Consider $J = \{j_1 < j_2 < \ldots < j_s\}$ in $\mathcal{A}_q(\Gamma)$ for $\Gamma = (\beta_1, \ldots, \beta_m)$, i.e., we have a path in the quantum Bruhat graph

$$1_W = w_0 \xrightarrow{\beta_{j_1}} w_1 \xrightarrow{\beta_{j_2}} \ldots \xrightarrow{\beta_{j_s}} w_s \, .$$

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Theorem (L.-Naito-Sagaki-Schilling-Shimozono) Given $J \in \mathcal{A}_q(\Gamma)$, which is identified with $B^{\otimes \mathbf{p}}$, we have

$$D_B(J) = -\text{height}(J).$$

This is the (unique) affine crystal isomorphism which commutes factors in the tensor product of KR crystals $B^{\otimes p}$

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Theorem (L.-Lubovsky)

We give a uniform realization, based on the quantum alcove model, of the combinatorial *R*-matrix.

We use combinatorial moves based on certain operators on W defined by QBG(W), which satisfy the Yang-Baxter equation (Brenti-Fomin-Postnikov).



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Example in type A_2 .

$$\mathbf{p} = (1, 2, 2, 1) =$$
; $\lambda = \omega_1 + \omega_2 + \omega_2 + \omega_1 = (4, 2, 0).$

A $\lambda\text{-chain}$ as a concatenation of $\omega_1\text{-},\,\omega_2\text{-},\,\omega_2\text{-},$ and $\omega_1\text{-chains:}$

 $\Gamma = ((1,2), (1,3) \mid (2,3), (1,3) \mid (2,3), (1,3) \mid (1,2), (1,3)).$

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Example. Let $J = \{1, 2, 3, 6, 7, 8\}$. ((1,2), (1,3) | (2,3), (1,3) | (2,3), (1,3) | (1,2), (1,3)).

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The corresponding element in $B^{\otimes p} = B^{1,1} \otimes B^{2,1} \otimes B^{2,1} \otimes B^{1,1}$ represented via column-strict fillings:

$$3 \otimes \frac{2}{3} \otimes \frac{1}{2} \otimes 3.$$

The energy calculation

Example. Consider the running example: $\lambda = \omega_1 + \omega_2 + \omega_2 + \omega_1$ in type A_2 .

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$$\begin{split} \Gamma &= ((\underline{1,2}), (\underline{1,3}) \mid (\underline{2,3}), (1,3) \mid (2,3), (\underline{1,3}) \mid (\underline{1,2}), (\underline{1,3})), \\ (h_i) &= (2, 4 \mid 2, 3 \mid 1, 2 \mid 1, 1). \end{split}$$
We have

 $\operatorname{height}(J) = 2.$