

Using time-dependent Gaussian basis sets in quantum dynamics simulations.

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- ▶ The G-MCTDH algorithm
- ▶ Stability of GWP propagation
- ▶ Phase space coverage and convergence



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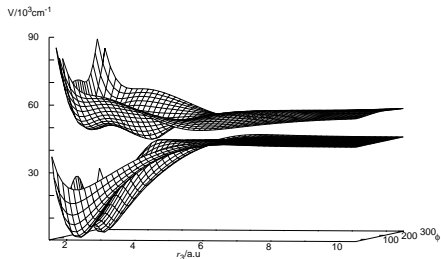
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- ▶ Phase space coverage and convergence



Collaboration: Irene Burghardt, Frankfurt
Benjamin Lasorne, Montpellier

Quantum Dynamics of Large Molecules

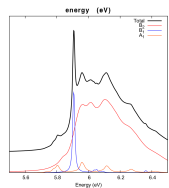
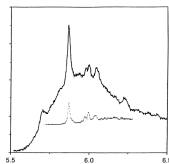
Aim is to follow nuclear dynamics over potential surfaces:



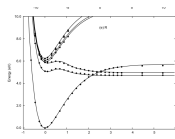
Ammonia branching ratio:

$$\Gamma_A^X(2D) \approx 2$$

$$\Gamma_A^X(6D) \approx 100$$

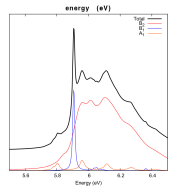
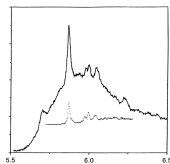
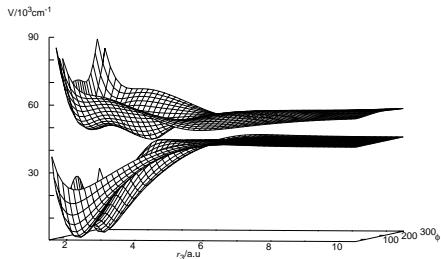


Pyrrole
absorption



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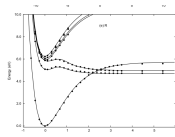
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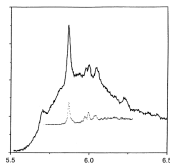
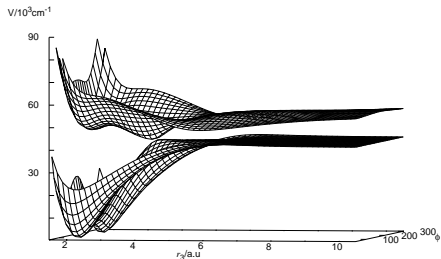
Bottlenecks:

- ▶ Dimensionality of wavefunction
- ▶ Obtaining potential

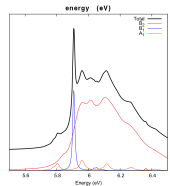


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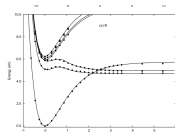
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Gaussian Wavepackets?

The MCTDH method

$$\Psi(Q_1, \dots, Q_f, t) = \sum_{j_1=1}^{n_1} \dots \sum_{j_f=1}^{n_p} A_{j_1 \dots j_f}(t) \prod_{\kappa=1}^p \varphi_{j_\kappa}^{(\kappa)}(Q_\kappa, t)$$

Variational equations of motion for A and φ .

$$i\dot{A}_J = \sum_L \langle \Phi_J | H | \Phi_L \rangle A_L - \sum_{\kappa=1}^p \sum_{l=1}^{n_\kappa} g_{j_\kappa l}^{(\kappa)} A_{J^\kappa}$$
$$i\dot{\varphi}^{(\kappa)} = (\mathbf{f}^{(\kappa)})^T \varphi^{(\kappa)} + \left(\mathbf{1} - \mathbf{P}^{(\kappa)} \right) \left(\rho^{(\kappa)} \right)^{-1} \langle \mathbf{H} \rangle^{(\kappa)} \varphi^{(\kappa)}$$

with the constraint operator matrix

$$f_{ij} = \langle \phi_i | \mathbf{f} | \phi_j \rangle = i \langle \phi_i | \dot{\phi}_j \rangle$$

computer resources $\sim n^p + nNp$

Reviews: Beck *et al* Phys. Rep. (00) 324:1

Meyer and Worth TCA (03) 109:251

The G-MCTDH Method

$$\Psi(Q_1, \dots, Q_f, t) = \sum_{j_1=1}^{n_1} \dots \sum_{j_p=1}^{n_f} A_{j_1 \dots j_p}(t) \prod_{\kappa=1}^{p-n} \varphi_{j_\kappa}^{(\kappa)} \prod_{\kappa=n+1}^p g_{j_\kappa}^{(\kappa)}$$

Replace single-particle functions with Gaussian functions

$$g_j(\mathbf{Q}, t) = \exp(\mathbf{Q}^T \zeta_j \mathbf{Q} + \mathbf{Q}^T \xi_j + \eta_j)$$

The G-MCTDH Method

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$$g_j(\mathbf{Q}, t) = \exp(\mathbf{Q}^T \boldsymbol{\zeta}_j \mathbf{Q} + \mathbf{Q}^T \boldsymbol{\xi}_j + \eta_j)$$

Propagate parameters $\lambda = \{\boldsymbol{\zeta}, \boldsymbol{\xi}, \eta\}$

$$\begin{aligned} i\dot{A}_j &= \sum_{lk} S_{jk}^{-1} \langle \Phi_k | H | \Phi_l \rangle A_l - \sum_{\kappa=1}^p \sum_{l=1}^{n_\kappa} i S_{jk}^{-1} \langle g_k | \frac{\partial}{\partial t} g_l \rangle A_{J_l^\kappa} \\ &= \sum_{lk} S_{jk}^{-1} H_{kl} A_l - \sum_{\kappa=1}^p \sum_{l=1}^{n_\kappa} i S_{jk}^{-1} \tau_{kl} A_{J_l^\kappa} \\ i\dot{\boldsymbol{\Lambda}} &= \mathbf{C}^{-1} \mathbf{Y} \end{aligned}$$

Burghardt *et al* JCP (99) 99:2927

$$Y_{i\alpha} = \sum_l \rho_{il} \left(H_{il}^{(\alpha 0)} - \left[\mathbf{s}^{(\alpha 0)} \mathbf{s}^{-1} \mathbf{H} \right]_{il} \right)$$

$$C_{i\alpha, j\beta} = \rho_{ij} \left(S_{ij}^{(\alpha\beta)} - \left[\mathbf{s}^{(\alpha 0)} \mathbf{s}^{-1} \mathbf{s}^{(0\beta)} \right]_{ij} \right)$$

where α is a parameter and i a function

$$S_{il}^{(\alpha\beta)} = \left\langle \frac{\partial g_i}{\partial \lambda_{i\alpha}} \left| \frac{\partial g_l}{\partial \lambda_{l\beta}} \right. \right\rangle ; \quad H_{il}^{(\alpha\beta)} = \left\langle \frac{\partial g_i}{\partial \lambda_{i\alpha}} \left| \hat{H} \left| \frac{\partial g_l}{\partial \lambda_{l\beta}} \right. \right. \right\rangle$$

e.g. if $\lambda_{i\alpha} = \xi_{i\alpha}$

$$S_{il}^{(\alpha 0)} = \left\langle \frac{\partial g_i}{\partial \xi_{i\alpha}} \left| g_l \right. \right\rangle = \langle g_i | x_\alpha | g_l \rangle$$

If only GWP known as *variational Multi-Configurational Gaussian* (vMCG) approach

Alternative Ansatz

Return to original MCTDH equation and variational derivation:

$$\Psi(Q_1, \dots, Q_f, t) = \sum_{j_1=1}^{n_1} \dots \sum_{j_f=1}^{n_p} A_{j_1 \dots j_p}(t) \prod_{\kappa=1}^p \varphi_{j_\kappa}^{(\kappa)}(Q_\kappa, t)$$

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and now using

$$|\varphi_r\rangle = \sum_{\alpha=1}^m |g_\alpha\rangle D_{\alpha r} \quad ; \quad r = 1, n$$

vary δA_J , $\delta \varphi_j$ (grid-based SPFs), δD_{ir} and $\delta g_\alpha = \sum_a \delta \lambda_{a\alpha} \frac{\partial g_\alpha}{\partial \lambda_{a\alpha}}$

obtain the same EOMs for the A_J , φ_j and $\lambda_{a\alpha}$ as before, but SPFs represented by

$$i\dot{D}_{\gamma i} = \sum_{lj\alpha} S_{\gamma\alpha}^{-1} \rho_{ij}^{-1} \langle g_\alpha | (1 - P) \langle H \rangle_{jl} | \varphi_l \rangle + \sum_m f_{mi} D_{\gamma m} - \sum_{\alpha\beta} S_{\gamma\alpha}^{-1} \tau_{\alpha\beta} D_{\beta i}$$

with $P = \sum_r |\varphi_r\rangle \langle \varphi_r|$.

Thus GWP's act as a time-dependent primitive basis.

1. This should provide a better G-MCTDH, with a normalised A-vector and so able to use the CMF integrator.
2. A-vector will be shorter than the original G-MCTDH.

This is the first layer of "Multi-layer G-MCTDH"

Römer, Ruckenbauer and Burghardt JCP (13) **138**: 064106

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Römer, Ruckebauer and Burghardt JCP (13) **138**: 064106

If $m = n$ then equivalent to original G-MCTDH

For vMCG, only 1 “SPF” expanded in basis. Can now write:

$$\begin{aligned}i\dot{A} &= EA \quad ; \quad E = \langle \varphi | H | \varphi \rangle \\i\dot{D}_\alpha &= \sum_{\alpha\beta} S_{\gamma\alpha}^{-1} (H_{\alpha\beta} - i\tau_{\alpha\beta}) D_\beta - ED_\alpha\end{aligned}$$

and the overall phase is moved into the A-coefficient.

Connection to trajectories

For frozen GWPs, taking

$$H = \sum_{\kappa} \frac{p_{\kappa}^2}{2m_{\kappa}} + V$$

Connection to trajectories

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$$H = \sum_{\kappa} \frac{p_{\kappa}^2}{2m_{\kappa}} + V + H_R$$

Connection to trajectories

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$$H = \sum_{\kappa} \frac{p_{\kappa}^2}{2m_{\kappa}} + V$$

In terms of the moments,

$$\langle H \rangle_{jl} = S_{jl} X_l^{(0)} + \sum_{\beta=1}^f S_{jl}^{(0\beta)} X_l^{(\beta)} + \sum_{\alpha,\beta=1}^f S_{jl}^{(\alpha\beta)} X_l^{(\alpha\beta)} + \dots$$

For example

$$X_l^{(\beta)} = -2i \frac{A_{\beta\beta j}}{m_r} p_{\beta j} + \frac{4A_{\beta\beta j}^2}{m_r} q_{\beta j} + \left. \frac{\partial V}{\partial x_{\beta}} \right|_{q_l} + \sum_{\alpha} \left. \frac{\partial^2 V}{\partial x_{\beta} \partial x_{\alpha}} q_{\alpha l} \right|_{q_l} + \dots$$

and as

$$i\dot{\xi}_{l\beta} = 2iA_{l\beta}\dot{q}_{l\beta} - \dot{p}_{l\beta}$$

the EOMs for the linear parameters can be written

$$\begin{aligned}\dot{q}_{l\beta} &= \frac{p_{l\beta}}{m_\beta} + \frac{1}{2A_{l\beta}} \operatorname{Im} \sum_{m\alpha} C_{l\beta m\alpha}^{-1} \tilde{Y}_{m\alpha} \\ \dot{p}_{l\beta} &= -V'_{l\beta} + \operatorname{Re} \sum_{m\alpha} C_{l\beta m\alpha}^{-1} \tilde{Y}_{m\alpha}\end{aligned}$$

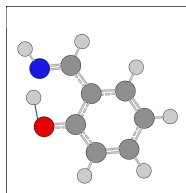
with

$$\tilde{Y}_{m\alpha} = \frac{4A_{j\alpha}^2}{m_\alpha} q_{j\beta} + \sum_{\beta \neq \alpha} v''_{j\alpha\beta} q_{l\alpha} + \sum_{\alpha, \beta=1}^f s_{jl}^{(\alpha\beta)} x_l^{(\alpha\beta)} + \dots$$

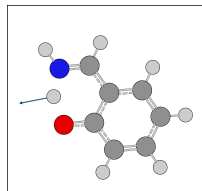
Salicylaldehyde Test Case: 2D Proton transfer

Hamiltonian in normal modes fitted to RHF/3-21G*

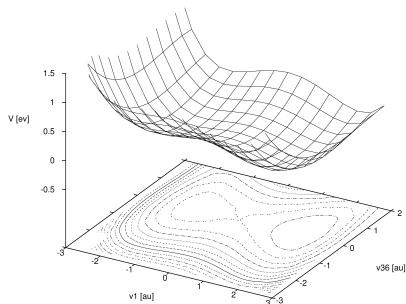
$$H = \frac{1}{2} \sum_{\kappa=1,18} \frac{\omega_{\kappa}}{2} \left(\frac{\partial^2}{\partial q_{\kappa}^2} + q_{\kappa}^2 \right) + \sum_{n=1}^4 A_n q_1^n + B_{11} q_1 q_{18} + B_{22} q_1^2 q_{18}^2 + B_{31} q_1^3 q_{18} + B_{13} q_1 q_{18}^3$$



ν_1



ν_{18}



Inversion C-matrix: Stability and Convergence

At the start, \mathbf{C} is singular. Values are due to density matrix, ρ and the projection of the derivative functions out of the GWP space

$$C_{i\alpha j\beta} = \rho_{ij} \left\langle \frac{\partial g_i}{\partial \lambda_{i\alpha}} \left| 1 - \sum_{rs} |g_r\rangle S_{rs}^{-1} \langle g_s| \right| \frac{\partial g_j}{\partial \lambda_{j\beta}} \right\rangle$$

Need only to include functions that have a significant population AND significantly contribute to the projected space (i.e. where the basis functions can move to).

Assume functions do not change much over a step, yet populations do. Estimate coefficients at end of step:

$$A_j(t + \delta t) = A_j(t) + \dot{A}_j(t)\delta t$$

Look at eigenvalues c_i of \mathbf{C} and count n_{ignore} with $c_i < \varepsilon$.

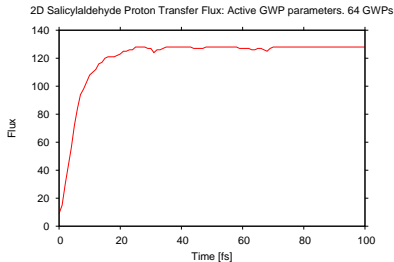
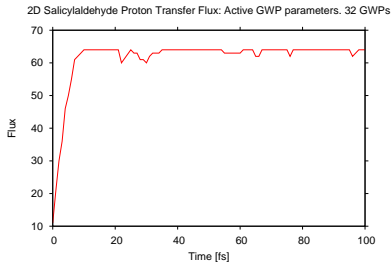
If $n_{\text{ignore}} > 0$, project $\frac{\partial g_i}{\partial \lambda_{i\alpha}}$ onto space spanned by $N - n_{\text{ignore}}$ eigenvectors,

$$C_{i\alpha} = \sum_{a=n_{\text{ignore}}+1}^N \left\langle \frac{\partial g_i}{\partial \lambda_{i\alpha}} \middle| a \right\rangle \left\langle a \middle| \frac{\partial g_i}{\partial \lambda_{j\alpha}} \right\rangle$$

and ignore n_{ignore} parameters with lowest values of $C_{i\alpha}$

With $\varepsilon = 1 \times 10^{-6}$:

32 GWPs (left) and 64 GWPs (right)



Linear Dependence

With large basis sets linear dependencies occur. Seen as step sizes drop as lowest eigenvalue of \mathbf{S} matrix drop below 1×10^{-6} .

Test 2D salicylaldehyde with 32 functions w/o dynamic selection requires 63784 steps and with 64 does not finish.

use similar procedure as dynamic selection

Look at eigenvalues s_i of \mathbf{S} and count n_{linear} with $s_i < \epsilon$.

If $n_{\text{linear}} > 0$, project g_j onto space spanned by $n - n_{\text{linear}}$ eigenvectors,

$$S_j = \sum_a = n_{\text{linear}} + 1^n \langle g_j | a \rangle \langle a | g_j \rangle$$

and ignore n_{ignore} functions with lowest values of S_j

with $\epsilon = 1 \times 10^{-6}$

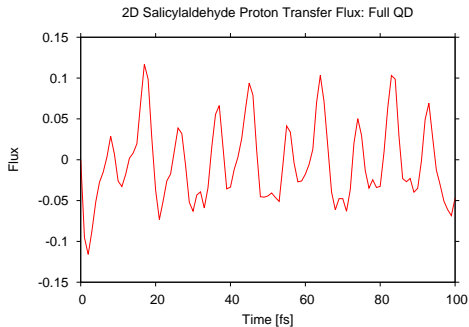
32 GWP finds 2 dependencies at 4 fs and 42 fs.

64 GWP test finds 7 dependencies at 4 fs and 42 fs.

Flux through barrier

Starting with $(q_1, q_{18}) = (0.96, -0.07)$. Corresponds to O–H bond stretched as lower minimum is at $(q_1, q_{18}) = (1.26, -0.04)$. Energy below barrier height.

4th order integrals.

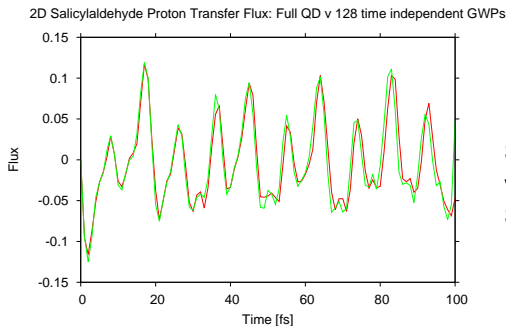


Exact solution.

Flux through barrier

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Exact (red).

Size: 2542 Time 16 s

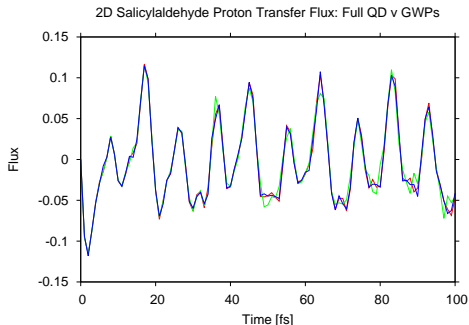
vMCG 128 t-ind (green).

Size: 512 Time: 2294 s

Flux through barrier

Starting with $(q_1, q_{18}) = (0.96, -0.07)$. Corresponds to O–H bond stretched as lower minimum is at $(q_1, q_{18}) = (1.26, -0.04)$. Energy below barrier height.

4th order integrals.



Exact (red).

Size: 2542. Time: 16 s

vMCG 16 (green).

Size: 64. Time: 68 s

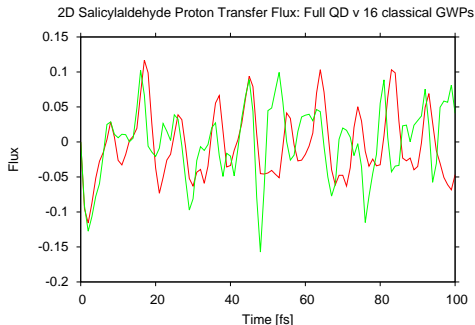
vMCG 32 (blue).

Size: 128. Time: 205 s

Flux through barrier

Starting with $(q_1, q_{18}) = (0.96, -0.07)$. Corresponds to O–H bond stretched as lower minimum is at $(q_1, q_{18}) = (1.26, -0.04)$. Energy below barrier height.

4th order integrals.

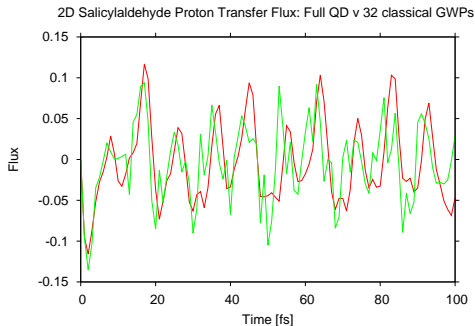


Exact (red). Time: 16 s
16 class GWPs (green).
Size: 64 Time: 8 s

Flux through barrier

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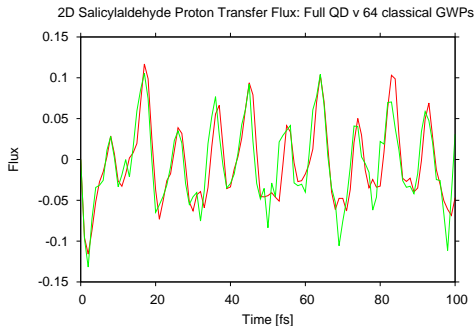


Exact (red). Time: 16 s
32 class GWPs (green).
Size: 128 Time: 24 s

Flux through barrier

Starting with $(q_1, q_{18}) = (0.96, -0.07)$. Corresponds to O–H bond stretched as lower minimum is at $(q_1, q_{18}) = (1.26, -0.04)$. Energy below barrier height.

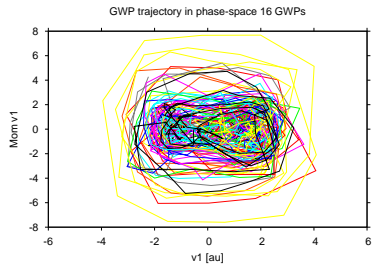
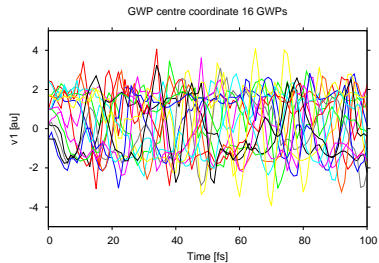
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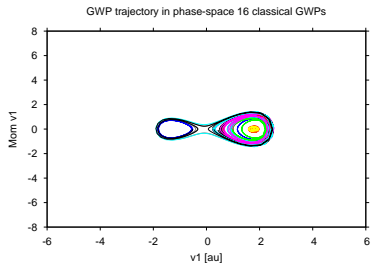
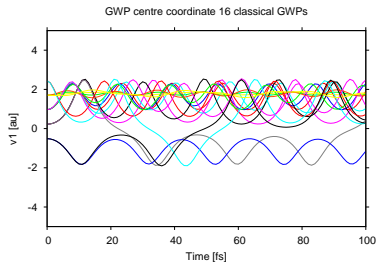
Exact (red). Time: 161 s
64 class GWPs (blue).
Size: 256 Time: 161 s

Trajectories with 16 GWPs

vMCG



Classical



Local Harmonic Approximation

For GWP calculations, common to use LHA, i.e. expand potential around centre q_l

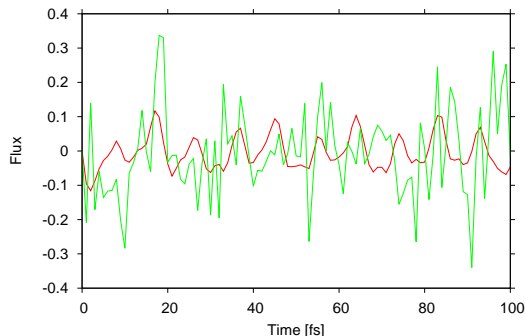
$$V_l(\mathbf{x}) = V_{0\alpha} + \sum_{\alpha} V'_{\alpha}(x_{\alpha} - q_{l\alpha}) + \frac{1}{2} \sum_{\alpha\beta} V''_{\alpha\beta}(x_{\alpha} - q_{l\alpha})(x_{\beta} - q_{l\beta})$$

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2D Salicylaldehyde Proton Transfer Flux: Full QD v 32 GWPs LHA. Full width = 0.

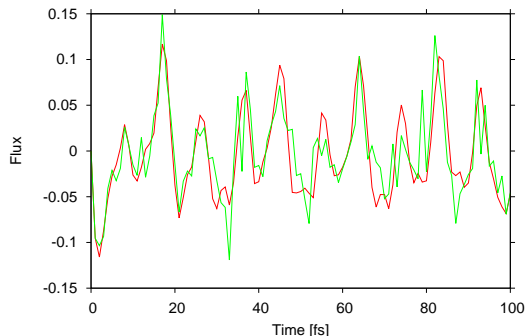


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2D Salicylaldehyde Proton Transfer Flux: Full QD v 32 GWPs LHA width = 0.4

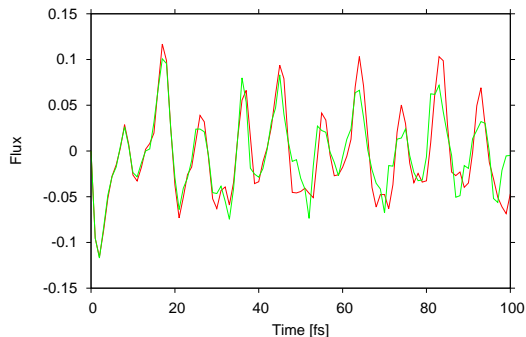


Local Harmonic Approximation

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2D Salicylaldehyde Proton Transfer Flux: Full QD v 32 GWPs LHA. Width = 0.3

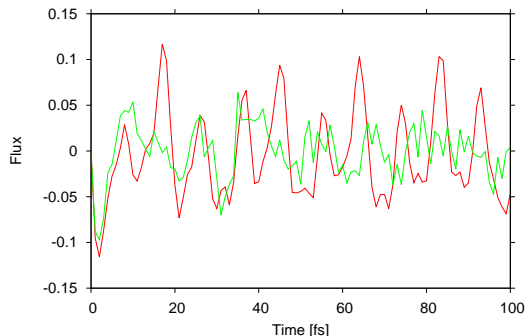


Local Harmonic Approximation

For GWP calculations, common to use LHA, i.e. expand potential around centre q_l

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2D Salicylaldehyde Proton Transfer Flux: Full QD v 32 GWPs LHA. Width = 0.2

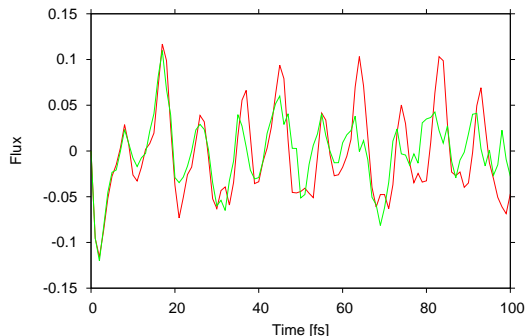


Local Harmonic Approximation

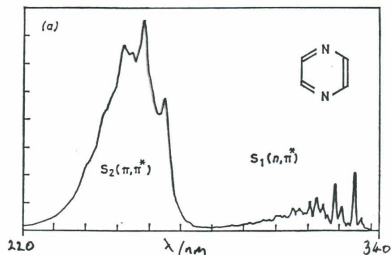
For GWP calculations, common to use LHA, i.e. expand potential around centre q_l

$$V_l(\mathbf{x}) = V_{0\alpha} + \sum_{\alpha} V'_{\alpha}(x_{\alpha} - q_{l\alpha}) + \frac{1}{2} \sum_{\alpha\beta} V''_{\alpha\beta}(x_{\alpha} - q_{l\alpha})(x_{\beta} - q_{l\beta})$$

2D Salicylaldehyde Proton Transfer Flux: Full QD v 64 GWPs LHA. Width = 0.2



Pyrazine Excitation: Model Hamiltonian



Yamazaki *et al* Farad. Discuss. (83) **75**: 395

The pyrazine molecule has 24 vibrational modes. NB. LHA is exact

$$\mathbf{H} = \sum_i \frac{\omega_i}{2} \left(-\frac{\partial^2}{\partial Q_i^2} + Q_i^2 \right) \mathbf{1} + \begin{pmatrix} -\Delta & 0 \\ 0 & \Delta \end{pmatrix} + \sum_{i \in G_1} \begin{pmatrix} \kappa_i^{(1)} & 0 \\ 0 & \kappa_i^{(2)} \end{pmatrix} Q_i +$$

$$\sum_{(i,j) \in G_2} \begin{pmatrix} \gamma_{i,j}^{(1)} & 0 \\ 0 & \gamma_{i,j}^{(2)} \end{pmatrix} Q_i Q_j + \sum_{i \in G_3} \begin{pmatrix} 0 & \lambda_i \\ \lambda_i & 0 \end{pmatrix} Q_i + \sum_{(i,j) \in G_4} \begin{pmatrix} 0 & \mu_{i,j} \\ \mu_{i,j} & 0 \end{pmatrix} Q_i Q_j.$$

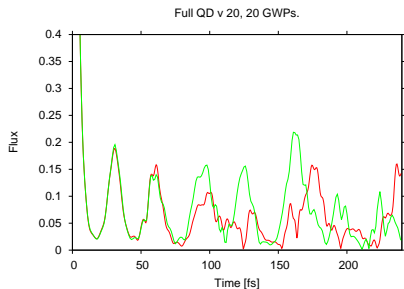
4D model: Linear Coupling

Autocorrelation function:

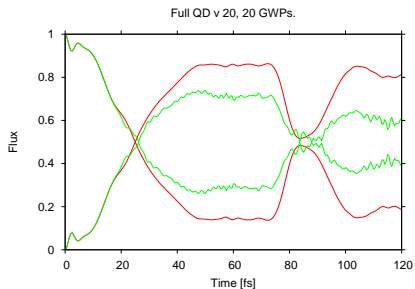
State Populations:

4D model: Linear Coupling

Autocorrelation function:

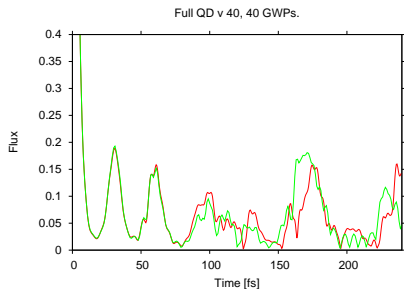


State Populations:

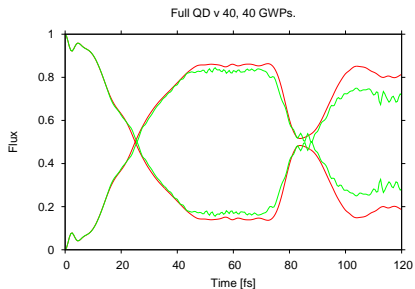


4D model: Linear Coupling

Autocorrelation function:

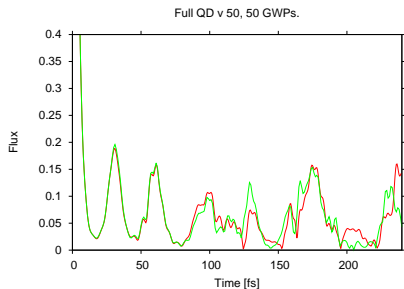


State Populations:

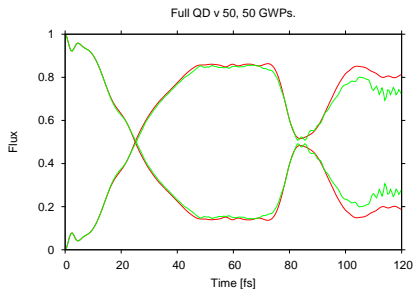


4D model: Linear Coupling

Autocorrelation function:

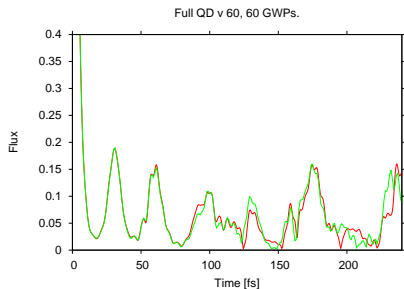


State Populations:

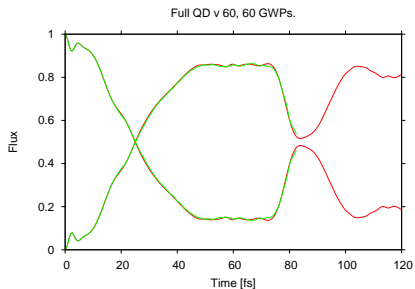


4D model: Linear Coupling

Autocorrelation function:



State Populations:



Conclusions

- ▶ Can use GWP basis in vMCG to provide converged quantum dynamics calculations
- ▶ vMCG moves to cover phase space as required
- ▶ Numerical problems due to linear dependencies and projector.