

On the MCTDH method

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Setting

Full quantum dynamics: Schrödinger equation

$$i\dot{\psi} = H\psi, \quad \psi = \psi(x_1, \dots, x_N, t)$$

↑

Intermediate models: via Dirac–Frenkel variational principle

Born–Oppenheimer, MCTDH, Gaussian wavepackets, ...

↓

Classical molecular dynamics: Newtonian eqs. of motion

$$M\ddot{x} = -\nabla V(x), \quad x = (x_1(t), \dots, x_N(t))$$

Time-dependent Schrödinger equation

$$i \frac{\partial \psi}{\partial t} = H\psi$$

N -particle wavefunction: $\psi = \psi(x, t)$

for $x = (x_1, \dots, x_N) \in \mathbb{R}^{3N}$, $x_n \in \mathbb{R}^3$, $t > 0$

Hamiltonian operator: $H = T + V$

with the kinetic and potential energy operators

$$T = - \sum_{j=1}^N \frac{1}{2m_j} \Delta_{x_j} \quad \text{and} \quad V = V(x) \\ \text{(smooth and bounded)}$$

high-dimensional linear PDE

Outline

The Dirac–Frenkel variational approximation principle

The MCTDH method, or dynamical low-rank approximation

Modelling error of MCTDH

Splitting integrators for MCTDH

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Dirac-Frenkel principle: abstract formulation

$\mathcal{M} \subset L^2(\mathbb{R}^{3N})$ approximation manifold

$T_u\mathcal{M}$ tangent space at $u \in \mathcal{M}$, complex linear

approximate wavefunction $u = u(\cdot, t) \in \mathcal{M}$ determined from

$$\langle i\dot{u} - Hu, v \rangle = 0 \quad \text{for all } v \in T_u\mathcal{M}$$

Galerkin method on a state-dependent approximation space

Dirac 1930, Frenkel 1934 (TDHF)

Orthogonal projection

$$\text{Re} \langle \dot{u} - \frac{1}{i}Hu, v \rangle = 0 \quad \text{for all } v \in T_u\mathcal{M}$$

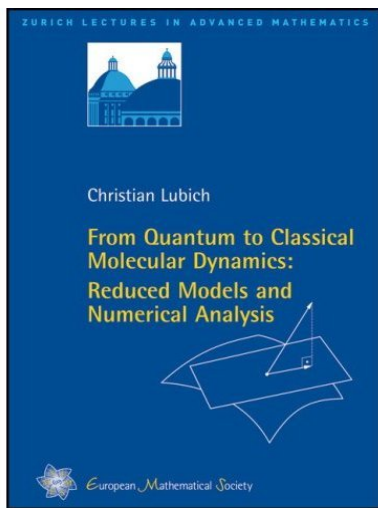
equivalent to **minimum defect condition**:

determine approximation $t \mapsto u(t) \in \mathcal{M}$ such that

$$\dot{u} = \vartheta \in T_u\mathcal{M} \quad \text{with} \quad \|\vartheta - \frac{1}{i}Hu\| = \min!$$

orthogonal projection $\dot{u} = P(u)\frac{1}{i}Hu$

Cover illustration: tangent space projection



$$\dot{u} = P(u) \frac{1}{i} H u$$

Symplectic projection

$$\operatorname{Im} \langle \dot{u} - \frac{1}{i}Hu, v \rangle = 0 \quad \text{for all } v \in T_u\mathcal{M}$$

- symplectic 2-form on $L^2(\mathbb{R}^{3N})$: $\omega(\xi, \eta) = -2 \operatorname{Im} \langle \xi, \eta \rangle$
- Hamiltonian function: $\mathcal{H}(u) = \langle u, Hu \rangle$

Hamiltonian system on the manifold \mathcal{M} :

$$\omega(\dot{u}, v) = d\mathcal{H}(u)v \quad \text{for all } v \in T_u\mathcal{M}$$

Consequences:

- ▶ energy conservation: $\mathcal{H}(u(t)) = \text{Const.}$
- ▶ symplecticity of the flow: for $\xi(t), \eta(t)$ tangent vectors propagated by the linearized flow along $u(t)$,

$$\omega(\xi(t), \eta(t)) = \text{Const.}$$

- ▶ A posteriori error bound

$$\|u(t) - \psi(t)\| \leq \int_0^t \text{dist}\left(\frac{1}{i}Hu(\tau), T_{u(\tau)}\mathcal{M}\right) d\tau$$

- ▶ Quasi-optimality

$$\|u(t) - \psi(t)\| \leq d(t) + Ce^{c\kappa t} \int_0^t d(\tau) d\tau$$

with $d(t) = \text{dist}(\psi(t), \mathcal{M})$ the best-approximation error,
 κ bound of the curvature of \mathcal{M}

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MCTDH

Multi-Configuration Time-Dependent Hartree method

Time-dependent Hartree method

Hartree product:

$$\psi(x_1, \dots, x_N, t) \approx \phi^{(1)}(x_1, t) \cdot \dots \cdot \phi^{(N)}(x_N, t)$$

Equations of motion for single-particle functions $\phi^{(n)}(x_n, t)$
by Dirac–Frenkel reduction to

$$\mathcal{M} = \{ \phi^{(1)} \otimes \dots \otimes \phi^{(N)} \mid \phi^{(n)} \in L^2(\mathbb{R}^3) \}$$

separation of variables, rank-1 approximation

MCTDH

Multi-Configuration Time-Dependent Hartree method

Approximate by linear combination of Hartree products

$$\psi(x, t) \approx \sum_{j_1=1}^{r_1} \cdots \sum_{j_N=1}^{r_N} a_{j_1 \dots j_N}(t) \phi_{j_1}^{(1)}(x_1, t) \cdot \dots \cdot \phi_{j_N}^{(N)}(x_N, t)$$

- ▶ mutually orthogonal single-particle functions $\phi_1^{(n)}, \dots, \phi_{r_n}^{(n)}$
- ▶ core tensor $(a_{j_1 \dots j_N})$ of full multilinear rank (r_1, \dots, r_N)

Tucker format, not canonical format
also hierarchical Tucker format

MCTDH equations of motion

from Dirac–Frenkel variational principle on MCTDH manifold:
coupled system of ODEs and low-dimensional nonlinear PDEs

$$i \frac{da_J}{dt} = \sum_L \langle \Phi_J | H | \Phi_L \rangle a_L$$
$$i \rho^{(n)} \frac{\partial \phi^{(n)}}{\partial t} = (I - P^{(n)}) \langle H \rangle_{-x_n} \phi^{(n)}$$

with $\Phi_J(x, t) = \prod_{n=1}^N \phi_{j_n}^{(n)}(x_n, t)$ for multi-indices $J = (j_1, \dots, j_N)$

Meyer, Manthe & Cederbaum 1990

existence and regularity: *Koch & L. 2007*

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Convergence of MCTDH to Schrödinger solution???

Naive expectation: Taking more and more terms in the linear combination of Hartree products yields an ever better accuracy

Obstructions:

- ▶ approximation properties of basis of Hartree products $\Phi_J(\cdot, t)$?
- ▶ ill-conditioned coefficient tensor (a_J)
 - ↔ ill-conditioned density matrices $\rho^{(n)} = A_{(n)}A_{(n)}^*$

Quasi-optimality of variational approximation

$u(t)$ MCTDH approximation, $\psi(t)$ exact wave function

MCTDH error bounded by best-approximation error:

$$\|u(t) - \psi(t)\| \leq d(t) + C_\kappa e^{c_\kappa t} \int_0^t d(\tau) d\tau$$

with $d(t) = \text{dist}(\psi(t), \mathcal{M})$ and κ curvature of manifold \mathcal{M}

but $\kappa \sim \text{cond}(\rho^{(n)})^{1/2} \rightarrow \infty$ as number of configurations $\rightarrow \infty$

Objectives

- ▶ error bounds in the case of ill-conditioned density matrices
- ▶ convergence as number of configurations $\rightarrow \infty$

under appropriate assumptions

Approximability assumption

Exact wavefunction can be written as

$$\psi(t) = v(t) + e(t)$$

where $v(t) \in \mathcal{M}$ has small defect:

$$\left\| i \frac{\partial v}{\partial t}(\cdot, t) - H v(\cdot, t) \right\| \leq \varepsilon$$

Small error $e(t)$ and small defect (ε) for **some** linear combination of Hartree products

Assumptions on the coefficient tensor

$B_{(n)}(t)$ n th unfolding of coefficient tensor $(b_{j_1 \dots j_N}(t))$ of $v(t) \in \mathcal{M}$,
for $n = 1, \dots, N$

- ▶ Bound of the pseudo-inverse (large: small δ):

$$\|B_{(n)}^\dagger(t)\|_2 = \frac{1}{\sigma_{r_n}(B_{(n)}(t))} \leq \delta^{-1}$$

- ▶ Small time derivatives of the coefficient tensor in components that correspond to small singular values:

$$\left\| B_{(n)}^\dagger(t) \dot{B}_{(n)}(t) \right\|_2 \leq c,$$

Error bound

$$\|u(t) - \psi(t)\| \leq \|e(t)\| + 2t\varepsilon \quad \text{for } t = O(\delta/\varepsilon)$$

Bound independent of δ (ill-conditioning)

for fixed rank, can achieve $\delta \sim \varepsilon$ by perturbing $v(t)$: $t = O(1)$

Convergence: if we *assume* $\delta \sim \varepsilon$, we can admit arbitrary number of configurations: $\varepsilon \rightarrow 0$

Proof of the error bound

use estimates for the tangent space projection $P(u)$ to show a quadratic differential inequality

$$\frac{d}{dt} \|u - v\| \leq \frac{C}{\delta} \|u - v\|^2 + \varepsilon$$

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Variational splitting: the formal game

variational approximation: $H = T + V$

$$\langle i\dot{u} - Hu, v \rangle = 0 \quad \text{for all } v \in T_u\mathcal{M}$$

time step $u_n \rightarrow u_{n+1}$ via

1. half-step with V : $u_{n+1/2}^-$ solution at $\Delta t/2$ of

$$\langle i\dot{u} - Vu, v \rangle = 0 \quad \text{for all } v \in T_u\mathcal{M}$$

2. full step with T : $u_{n+1/2}^+$ solution at Δt of

$$\langle i\dot{u} - Tu, v \rangle = 0 \quad \text{for all } v \in T_u\mathcal{M}$$

3. half-step with V : u_{n+1} solution at $\Delta t/2$ of

$$\langle i\dot{u} - Vu, v \rangle = 0 \quad \text{for all } v \in T_u\mathcal{M}$$

Variational splitting for MCTDH

- ▶ Step with T : since $Tu \in T_u\mathcal{M}$ for $u \in \mathcal{M}$, integration step with T decouples into $\dot{a}_J = 0$ and

$$i \frac{\partial \phi_j^{(k)}}{\partial t}(x_k, t) = -\frac{1}{2m_k} \Delta_{x_k} \phi_j^{(k)}(x_k, t)$$

low-dimensional free Schrödinger equations \rightarrow FFT

- ▶ Step with V : MCTDH for Hamiltonian V instead of H
 \rightarrow larger time steps independently of the space discretization
work in preparation (with I. Oseledets): integration scheme that does not suffer from ill-conditioning of the matrices $\rho^{(k)}$

Error of variational splitting for MCTDH

second order error bound

$$\|u_n - u(t_n)\| \leq C \Delta t^2 \max \|u(t)\|_{H^2}$$

same as for Strang splitting for the linear Schrödinger equation

basic tool: Lie-commutator bounds for vector fields corresponding to T and V

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