Stable traveling spots in a planar three-component FitzHugh-Nagumo system

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Outline



- Introduction
- Stationary spots
- Bifurcation to traveling spots:
 - Asymptotic analysis
 - Direct solver
 - AUTO
- Work in progress

Model



Generalized FitzHugh-Nagumo Equation:

$$oldsymbol{U}_t = arepsilon^2 \Delta oldsymbol{U} + oldsymbol{U} - oldsymbol{U}^3 - arepsilon(lpha oldsymbol{V} + eta oldsymbol{W} + \gamma)$$

$$au V_t = \Delta V + U - V$$

 $heta W_t = D^2 \Delta W + U - W$

where $0 < \epsilon << I$; D>I; $0 < \tau, \theta$; α, β, γ are constants.

- U: fast component
 - \rightarrow bistable: U = ±I
 - \rightarrow nonlinear: U³
 - coupling to the slow components is small
- V,W: slow components
 - ⇒ linear
 - only coupled to the fast component

Gas-discharge experiments



U: current density V: voltage drop W: surface charge black: U=-I, white: U=+I

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Inspiration





Stationary spot





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Theorem [vH, Sandstede '11]: Assume that $R_1 > 0$ solves: $lpha v_0 + eta w_0 + \gamma = -rac{\sqrt{2}}{3R_1}$ where v_0 , w_0 are given by $v_0 = 1 - 2R_1K_1(R_1)I_0(R_1), \quad w_0 = 1 - 2rac{R_1}{D}K_1\left(rac{R_1}{D}
ight)I_0\left(rac{R_1}{D}
ight)$ Then there exists a stationary radially symmetric spot with radius R_{I} . This spot is stable "if and only if" $\lambda(\ell) < 0$ for all $\ell = 0, 2, 3, ...,$ where $\lambda(\ell) = 3\sqrt{2}\varepsilon^2 \alpha R_1 \left(K_1(R_1)I_1(R_1) - K_\ell(R_1)I_\ell(R_1) \right) +$ $3\sqrt{2}arepsilon^2etarac{R_1}{D^2}\left(K_1\left(rac{R_1}{D}
ight)I_1\left(rac{R_1}{D}
ight)-K_\ell\left(rac{R_1}{D}
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- Spot corresponding to the smallest zero of existence condition is unstable with respect to $\ell = 0$ (radial perturbations)
- $\alpha, \beta \leq 0$: Spot is unstable with respect to $\ell = 0$ (radial perturbations)





Drift: asymptotics

Goal: Determine for arbitrary small ϵ the points (τ, θ) at which the stationary spot bifurcates to a traveling spot

Method: Weakly nonlinear analysis



• Speed is small (second small parameter):

$$c=\delta\,,\quad 0$$

•Traveling spot retains to leading order the shape of the stationary spot:

$$(U^T,V^T,W^T)=(U^s,V^s,W^s)+\delta(u,v,w)$$

• Determine eq^{ns} for (u,v,w) and use singular perturbation techniques to derive the drift line

Result: The drift line is given by

$$\frac{\sqrt{2}}{3R_1^2} = \alpha \hat{\tau} \left(I_1(R_1) K_2(R_1) - I_0(R_1) K_1(R_1) \right) + \frac{\beta \hat{\theta}}{D^3} \left(I_1\left(\frac{R_1}{D}\right) K_2\left(\frac{R_1}{D}\right) - I_0\left(\frac{R_1}{D}\right) K_1\left(\frac{R_1}{D}\right) \right)$$

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$$(\hat{ au},\hat{ heta})=arepsilon^2(au, heta)$$

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Other bifurcations

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The other bifurcations (ℓ =0,2,3,...), if present, will be Hopf bifurcations. These Hopf lines are implicitly given by:

$$\begin{array}{ll} 0 &=& \displaystyle \frac{1}{R_1^2} (1-\ell^2) + 3\sqrt{2} \alpha R_1 \left(I_1(R_1) K_1(R_1) - \Re \left[I_\ell \left(\sqrt{1+i\hat{\tau}|\hat{\lambda}(\ell)|} R_1 \right) K_\ell \left(\sqrt{1+i\hat{\tau}|\hat{\lambda}(\ell)|} R_1 \right) \right] \right) \\ &\quad + 3\sqrt{2} \frac{\beta}{D^2} R_1 \left(I_1 \left(\frac{R_1}{D} \right) K_1 \left(\frac{R_1}{D} \right) - \Re \left[I_\ell \left(\sqrt{1+i\hat{\theta}|\hat{\lambda}(\ell)|} \frac{R_1}{D} \right) K_\ell \left(\sqrt{1+i\hat{\theta}|\hat{\lambda}(\ell)|} \frac{R_1}{D} \right) \right] \right) \\ &\quad |\hat{\lambda}(\ell)| &=& -3\sqrt{2} \alpha R_1 \left(\Im \left[I_\ell \left(\sqrt{1+i\hat{\tau}|\hat{\lambda}(\ell)|} R_1 \right) K_\ell \left(\sqrt{1+i\hat{\tau}|\hat{\lambda}(\ell)|} R_1 \right) \right] \right) \\ &\quad - 3\sqrt{2} \frac{\beta}{D^2} R_1 \left(\Im \left[I_\ell \left(\sqrt{1+i\hat{\theta}|\hat{\lambda}(\ell)|} \frac{R_1}{D} \right) K_\ell \left(\sqrt{1+i\hat{\theta}|\hat{\lambda}(\ell)|} \frac{R_1}{D} \right) \right] \right) . \end{array}$$

$\begin{array}{rcl} \boldsymbol{U}_t &=& \varepsilon^2 \Delta \boldsymbol{U} &+& \boldsymbol{U} - \boldsymbol{U}^3 - \varepsilon (\alpha \boldsymbol{V} + \beta \boldsymbol{W} + \gamma) \\ \boldsymbol{\tau} \boldsymbol{V}_t &=& \Delta \boldsymbol{V} &+& \boldsymbol{U} - \boldsymbol{V} \\ \boldsymbol{\theta} \boldsymbol{W}_t &=& D^2 \Delta \boldsymbol{W} &+& \boldsymbol{U} - \boldsymbol{W} \end{array} \begin{array}{rcl} \boldsymbol{Specific parameters} \end{array}$

Choose the following set of parameters (for the remainder of presentation):

$$\alpha$$
 = 0.5 , β = 2, γ = 1, D = 2

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Then, there exists a stationary stable spot solution with (leading order) width

 $R_1 = 1.86$

The bifurcation diagram of this stationary spot looks like:

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Direct PDE solver

• Code written by K.-I. Ueda:

➡ 5-point discretization of the Laplacian on a 20 by 20 square with 200 equidistance mesh points

- Semi-implicit time scheme: conjugate gradients with incomplete Cholesky
- Parameter values:

$$\alpha = 0.5, \beta = 2, \gamma = 1, D = 2, \epsilon = 0.1, \hat{\tau} = 6, \hat{\theta} = 0.01$$

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Why AUTO?

• Direct simulations with the PDE solver are slow and costly since the speed of a traveling spot is slow, especially for small ϵ .

• For example, simulation shown was done far away from the drift bifurcation line, with relatively large ϵ :

Want: better numerical evidence for the drift bifurcation line and more flexibility. Tool: AUTO

AUTO

Rescale and co-moving frame:

$$(\hat{\tau},\hat{ heta}) = \varepsilon^2(au, heta), \ (x_1,x_2,t) \to (x_1 - \varepsilon^2 ct,x_2,t)$$

Stationary solution in moving frame:

$$egin{aligned} & -arepsilon^2 c U_{x_1} &= arepsilon^2 \Delta U \ + \ U - U^3 - arepsilon (lpha V + eta W + \gamma) \ & -c \hat{ au} V_{x_1} &= \Delta V \ + \ U - V \ & -c \hat{ heta} W_{x_1} &= D^2 \Delta W \ + \ U - W \end{aligned}$$
 Polar coordinates: $egin{aligned} x_1 &= r \cos \phi, \ x_2 &= r \sin \phi \end{bmatrix}$

$$\begin{aligned} -\varepsilon^2 c \left(\cos \phi \, U_r - \frac{\sin \phi}{r} U_\phi \right) &= \varepsilon^2 (U_{rr} + \frac{1}{r} U_r + \frac{1}{r^2} U_{\phi \phi}) &+ U - U^3 - \varepsilon (\alpha V + \beta W + \gamma) \\ -c \hat{\tau} \left(\cos \phi \, V_r - \frac{\sin \phi}{r} V_\phi \right) &= (V_{rr} + \frac{1}{r} V_r + \frac{1}{r^2} V_{\phi \phi}) &+ U - V \\ -c \hat{\theta} \left(\cos \phi \, W_r - \frac{\sin \phi}{r} W_\phi \right) &= D^2 (W_{rr} + \frac{1}{r} W_r + \frac{1}{r^2} W_{\phi \phi}) &+ U - W \end{aligned}$$

AUTO (cont.)

Write as a first order system (in the radial variable)

$$egin{aligned} u_r &=& rac{p}{arepsilon} \ p_r &=& -rac{p}{r} - rac{arepsilon}{r^2} u_{\phi\phi} - rac{u}{arepsilon} - rac{u^3}{arepsilon} + (lpha v + eta w + \gamma) - cp\cos\phi + arepsilon c rac{\sin\phi}{r} u_{\phi} \ v_r &=& q \ q_r &=& -rac{q}{r} - rac{1}{r^2} v_{\phi\phi} - u + v - \hat{ au} cq\cos\phi + \hat{ au} c rac{\sin\phi}{r} v_{\phi} \ w_r &=& z \ z_r &=& -rac{z}{r} - rac{1}{r^2} w_{\phi\phi} - rac{u}{D^2} + rac{w}{D^2} - rac{\hat{ heta} cz}{D^2} \cos\phi + rac{\hat{ heta} c}{D^2} rac{\sin\phi}{r} w_{\phi} \end{aligned}$$

Fourier in ϕ :

$$ar{U}(r,\phi) = \sum_{\ell=-\infty}^{\ell=\infty} ar{U}^\ell(r) e^{i\ell\phi}$$

recall:

$$\cos\phi = rac{1}{2}(e^{i\phi} + e^{-i\phi})\,, \ \sin\phi = rac{1}{2i}(e^{i\phi} - e^{-i\phi})\,, \ rac{\partialar{U}}{\partial\phi} = i\sum_{\ell=-\infty}^{\infty}\ell u^\ell e^{i\ell\phi} \ \ rac{\partial^2ar{U}}{\partial\phi^2} = -\sum_{\ell=-\infty}^{\infty}\ell^2 u^\ell e^{i\ell\phi}$$

AUTO (cont.)

So, we get:

$$\begin{split} u_r^\ell &= \frac{p^\ell}{\varepsilon} \\ p_r^\ell &= -\frac{p^\ell}{r} + \frac{\varepsilon \ell^2}{r^2} u^\ell - \frac{u^\ell}{\varepsilon} - \frac{nonl}{\varepsilon} + (\alpha v^\ell + \beta w^\ell + \gamma) \\ &- \frac{c}{2} (p^{\ell-1} + p^{\ell+1}) + \frac{\varepsilon c}{2r} \left((\ell-1) u^{\ell-1} - (\ell+1) u^{\ell+1} \right) \\ v_r^\ell &= q^\ell \\ q_r^\ell &= -\frac{q^\ell}{r} + \frac{\ell^2}{r^2} v^\ell - u^\ell + v^\ell - \frac{\hat{\tau}c}{2} (q^{\ell-1} + q^{\ell+1}) \\ &+ \frac{\hat{\tau}c}{2r} \left((\ell-1) v^{\ell-1} - (\ell+1) v^{\ell+1} \right) \\ w_r^\ell &= z^\ell \\ z_r^\ell &= -\frac{z^\ell}{r} + \frac{\ell^2}{r^2} w^\ell - \frac{u^\ell}{D^2} + \frac{w^\ell}{D^2} - \frac{\hat{\theta}c}{2D^2} (z^{\ell-1} + z^{\ell+1}) \\ &+ \frac{\hat{\theta}c}{2D^2 r} \left((\ell-1) w^{\ell-1} - (\ell+1) w^{\ell+1} \right) \end{split}$$

- Solutions need to be even: restrict ourself to $\ell \ge 0$
- γ only appears in the $\ell = 0$ -term!
- nonl-term contains infinitely many coupled terms

AUTO: We have to truncate to a finite number of Fourier modes

AUTO: Difficulties

Implement model in AUTO for a finite number of Fourier modes and on a finite domain [0, L] with appropriate boundary conditions.

2 Major difficulties:

- \bullet AUTO does not switch onto the traveling branch for increasing $\hat{\tau}$
 - Add a small symmetry breaking term:

$$p_r^1 = p_r^1 - \delta rac{r^2}{L}$$

- \rightarrow Continue in $\hat{\tau}$ beyond bifurcation point (speed becomes non-zero)
- \blacktriangleright Continue δ down to 0 (check that speed stays nonzero)

$\begin{array}{rcl} \boldsymbol{U}_{t} &=& \varepsilon^{2}\Delta\boldsymbol{U} &+& \boldsymbol{U}-\boldsymbol{U}^{3}-\varepsilon(\alpha\boldsymbol{V}+\boldsymbol{\beta}\boldsymbol{W}+\boldsymbol{\gamma})\\ \boldsymbol{\tau}\boldsymbol{V}_{t} &=& \Delta\boldsymbol{V} &+& \boldsymbol{U}-\boldsymbol{V}\\ \boldsymbol{\theta}\boldsymbol{W}_{t} &=& D^{2}\Delta\boldsymbol{W} &+& \boldsymbol{U}-\boldsymbol{W} \end{array} \quad \begin{array}{l} \textbf{Difficulties (cont.)} \end{array}$

Second major difficulty:

• AUTO detects many branch points, so it is not possible to detect the correct drift point and continue the drift line in the $(\hat{\tau}, \hat{\theta})$ -plane.

 \Rightarrow Detect drift bifurcation as points where the linearization L₁ has a generalized eigenfunction Ψ :

$$L_1\psi=Mar{U}_r^s$$

→ M is a diagonal matrix with $I, I/\hat{\tau}, D^2/\hat{\theta}$ on its diagonal and \bar{U}_r^s is the radial derivative of stationary spot and thus lies in the null space of L_I → Add small term $\delta \bar{U}_r^s$ to the eqⁿ (makes the system onto)

$$L_1\psi + \delta \bar{U}^s_r = M \bar{U}^s_r$$

Add integral condition to ensure that the kernel is 0 (solvability condition): $\langle \psi, \bar{U}_r^s
angle = 0$

- \rightarrow Unique solution (Ψ,δ)
- \Rightarrow We are at a drift bifurcation iff $\delta = 0$, so we continue δ down to 0
- Remove the δ-term, and continue in $\hat{\tau}$ or $\hat{\theta}$

AUTO: Results

Results:

• standard parameter values:

$$\alpha = 0.5, \beta = 2, \gamma = 1, D = 2, \hat{\theta} = 0.5$$

• 15 fourier modes, domain size = 12

Results (cont.)

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$$\alpha = 0.5, \beta = 2, \gamma = 1, D = 2, \epsilon = 0.1, \hat{\tau} = 6, \hat{\theta} = 0.01$$

AUTO

PDE solver at t=8500

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Snapshots of U-component at t=7500 and t=9500

• AUTO: predicted speed = 0.38

Hopf

Different set of parameter values!

 $\hat{\tau}$ =0.31: below the Hopf line

 $\hat{\tau}$ =0.32: above the Hopf line

Work in progress I

• Super vs subcritical? [Ei, Mimura, Nagayama 2006]

• Compare AUTO with PDE solver

Work in progress I

• Super vs subcritical? [Ei, Mimura, Nagayama 2006]

• Compare AUTO with PDE solver

Work in progress II

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Interaction of traveling spots (cartoon)

Questions??