# Using the Profile Likelihood in Searches for New Physics



Statistical issues relevant to significance of discovery claims BIRS, Banff, 11-16 July, 2010

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### Outline

Prototype search analysis for LHC Test statistics based on profile likelihood ratio Systematics covered via nuisance parameters Sampling distributions to get significance/sensitivity Asymptotic formulae from Wilks/Wald Examples:

> $n \sim \text{Poisson} (\mu s + b), m \sim \text{Poisson}(\tau b)$ Shape analysis

Conclusions

### Prototype search analysis

Search for signal in a region of phase space; result is histogram of some variable *x* giving numbers:

$$\mathbf{n} = (n_1, \ldots, n_N)$$

Assume the  $n_i$  are Poisson distributed with expectation values

$$E[n_i] = \mu s_i + b_i$$
strength parameter

where

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## Prototype analysis (II)

Often also have a subsidiary measurement that constrains some of the background and/or shape parameters:

$$\mathbf{m} = (m_1, \ldots, m_M)$$

Assume the  $m_i$  are Poisson distributed with expectation values

$$E[m_i] = u_i(\boldsymbol{\theta})$$
  
nuisance parameters ( $\boldsymbol{\theta}_{s}, \boldsymbol{\theta}_{b}, b_{tot}$ )

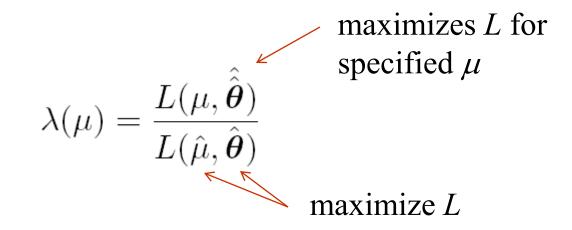
Likelihood function is

$$L(\mu, \theta) = \prod_{j=1}^{N} \frac{(\mu s_j + b_j)^{n_j}}{n_j!} e^{-(\mu s_j + b_j)} \quad \prod_{k=1}^{M} \frac{u_k^{m_k}}{m_k!} e^{-u_k}$$

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## The profile likelihood ratio

Base significance test on the profile likelihood ratio:



The likelihood ratio of point hypotheses gives optimum test (Neyman-Pearson lemma).

The profile LR hould be near-optimal in present analysis with variable  $\mu$  and nuisance parameters  $\theta$ .

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### Test statistic for discovery

Try to reject background-only ( $\mu = 0$ ) hypothesis using

$$q_0 = \begin{cases} -2\ln\lambda(0) & \hat{\mu} \ge 0\\ 0 & \hat{\mu} < 0 \end{cases}$$

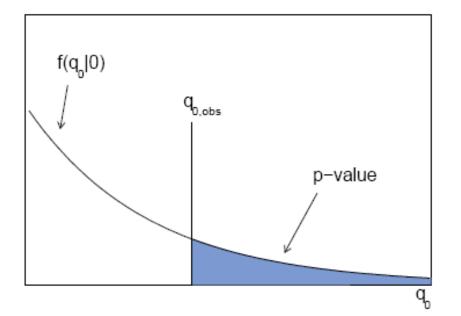
i.e. here only regard upward fluctuation of data as evidence against the background-only hypothesis.

### *p*-value for discovery

Large  $q_0$  means increasing incompatibility between the data and hypothesis, therefore *p*-value for an observed  $q_{0,obs}$  is

$$p_0 = \int_{q_{0,\text{obs}}}^{\infty} f(q_0|0) \, dq_0$$

will get formula for this later



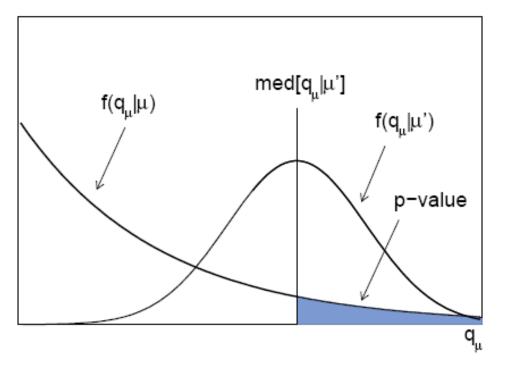
From *p*-value get equivalent significance,

$$Z = \Phi^{-1}(1-p)$$

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## Expected (or median) significance / sensitivity

When planning the experiment, we want to quantify how sensitive we are to a potential discovery, e.g., by given median significance assuming some nonzero strength parameter  $\mu'$ .



So for *p*-value, need  $f(q_0|0)$ , for sensitivity, will need  $f(q_0|\mu')$ ,

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Wald approximation for profile likelihood ratio To find *p*-values, we need:  $f(q_0|0)$ ,  $f(q_\mu|\mu)$ For median significance under alternative, need:  $f(q_\mu|\mu')$ 

Use approximation due to Wald (1943)

$$V^{-1} = -E\left[\frac{\partial^2 \ln L}{\partial \theta_i \partial \theta_j}\right]$$

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#### Noncentral chi-square for $-2\ln\lambda(\mu)$

If we can neglect the  $O(1/\sqrt{N})$  term,  $-2\ln\lambda(\mu)$  follows a noncentral chi-square distribution for one degree of freedom with noncentrality parameter

$$\Lambda = \frac{(\mu - \mu')^2}{\sigma^2}$$

As a special case, if  $\mu' = \mu$  then  $\Lambda = 0$  and  $-2\ln\lambda(\mu)$  follows a chi-square distribution for one degree of freedom (Wilks).

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## Distribution of $q_0$

Assuming the Wald approximation, we can write down the full distribution of  $q_0$  as

$$f(q_0|\mu') = \left(1 - \Phi\left(\frac{\mu'}{\sigma}\right)\right)\delta(q_0) + \frac{1}{2}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{q_0}}\exp\left[-\frac{1}{2}\left(\sqrt{q_0} - \frac{\mu'}{\sigma}\right)^2\right]$$

The special case  $\mu' = 0$  is a "half chi-square" distribution:

$$f(q_0|0) = \frac{1}{2}\delta(q_0) + \frac{1}{2}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{q_0}}e^{-q_0/2}$$

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Cumulative distribution of  $q_0$ , significance From the pdf, the cumulative distribution of  $q_0$  is found to be

 $F(q_0|\mu') = \Phi\left(\sqrt{q_0} - \frac{\mu'}{\sigma}\right)$ 

The special case  $\mu' = 0$  is

$$F(q_0|0) = \Phi\left(\sqrt{q_0}\right)$$

The *p*-value of the  $\mu = 0$  hypothesis is

 $p_0 = 1 - F(q_0|0)$ 

Therefore the discovery significance Z is simply

$$Z = \Phi^{-1}(1 - p_0) = \sqrt{q_0}$$

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### The Asimov data set

To estimate median value of  $-2\ln\lambda(\mu)$ , consider special data set where all statistical fluctuations suppressed and  $n_i$ ,  $m_i$  are replaced by their expectation values (the "Asimov" data set):

$$n_{i} = \mu' s_{i} + b_{i}$$

$$m_{i} = u_{i}$$

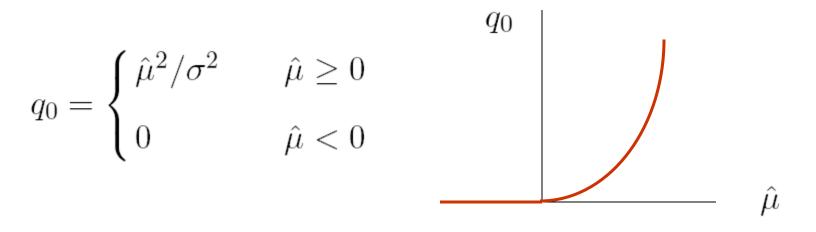
$$\rightarrow \hat{\mu} = \mu' \quad \hat{\theta} = \theta$$

$$\lambda_{A}(\mu) = \frac{L_{A}(\mu, \hat{\theta})}{L_{A}(\hat{\mu}, \hat{\theta})} = \frac{L_{A}(\mu, \hat{\theta})}{L_{A}(\mu', \theta)}$$

$$-2 \ln \lambda_{A}(\mu) = \frac{(\mu - \mu')^{2}}{\sigma^{2}} = \Lambda$$
Asimov value of -2ln $\lambda(\mu)$  gives non-centrality param.  $\Lambda$ , or equivalently,  $\sigma$ 

### Relation between test statistics and $\hat{\mu}$

Assuming Wald approximation, the relation between  $q_0$  and  $\hat{\mu}$  is



Monotonic, therefore quantiles of  $\hat{\mu}$  map one-to-one onto those of  $q_0$ , e.g.,

$$\operatorname{med}[q_0] = q_0(\operatorname{med}[\hat{\mu}]) = q_0(\mu') = \frac{{\mu'}^2}{\sigma^2} = -2\ln\lambda_{\rm A}(0)$$

$$\mathrm{med}[Z_0] = \sqrt{-2\ln\lambda_{\mathrm{A}}(0)}$$

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### Profile likelihood ratio for upper limits

For purposes of setting an upper limit on  $\mu$  use

$$q_{\mu} = \begin{cases} -2\ln\lambda(\mu) & \hat{\mu} \leq \mu \\ 0 & \hat{\mu} > \mu \end{cases} \quad \text{where} \quad \lambda(\mu) = \frac{L(\mu, \hat{\theta})}{L(\hat{\mu}, \hat{\theta})}$$

Note for purposes of setting an upper limit, one does not regard an upwards fluctuation of the data as representing incompatibility with the hypothesized  $\mu$ .

Note also here we allow the estimator for  $\mu$  be negative (but  $\hat{\mu}s_i + b_i$  must be positive).

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### Alternative test statistic for upper limits

Assume physical signal model has  $\mu > 0$ , therefore if estimator for  $\mu$  comes out negative, the closest physical model has  $\mu = 0$ .

Therefore could also measure level of discrepancy between data and hypothesized  $\mu$  with

$$\tilde{\lambda}(\mu) = \begin{cases} \frac{L(\mu, \hat{\hat{\boldsymbol{\theta}}}(\mu))}{L(\hat{\mu}, \hat{\boldsymbol{\theta}})} & \hat{\mu} \ge 0, \\ \frac{L(\mu, \hat{\hat{\boldsymbol{\theta}}}(\mu))}{L(0, \hat{\hat{\boldsymbol{\theta}}}(0))} & \hat{\mu} < 0. \end{cases} \quad \tilde{q}_{\mu} = \begin{cases} -2\ln\tilde{\lambda}(\mu) & \hat{\mu} \le \mu \\ 0 & \hat{\mu} > \mu \end{cases}$$

Performance not identical to but very close to  $q_{\mu}$  (of previous slide).  $q_{\mu}$  is simpler in important ways.

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Relation between test statistics and  $\hat{\mu}$ Assuming the Wałd approximation for  $-2\ln\lambda(\mu)$ ,  $q_{\mu}$  and  $\tilde{q}_{\mu}$ both have monotonic relation with  $\mu$ .

$$q_{\mu} = \begin{cases} \frac{(\mu - \hat{\mu})^2}{\sigma^2} & \hat{\mu} < \mu & \tilde{q}_{\mu} \\ 0 & \hat{\mu} > \mu & \tilde{q}_{\mu} \\ \hline \mu & \mu & \hat{\mu} \end{cases}$$

$$\tilde{q}_{\mu} = \begin{cases} \frac{\mu^2}{\sigma^2} - \frac{2\mu\hat{\mu}}{\sigma^2} & \hat{\mu} < 0 & \text{And therefore quantiles} \\ \frac{(\mu - \hat{\mu})^2}{\sigma^2} & 0 \le \hat{\mu} \le \mu & \text{of } q_{\mu}, \tilde{q}_{\mu} \text{ can be obtained} \\ 0 & \hat{\mu} > \mu \\ 0 & \text{of } \hat{\mu} > \mu \\ \end{cases}$$

 $q_{\mu}$ 

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# Distribution of $q_{\mu}$

#### Similar results for $q_{\mu}$

$$f(q_{\mu}|\mu') = \Phi\left(\frac{\mu'-\mu}{\sigma}\right)\delta(q_{\mu}) + \frac{1}{2}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{q_{\mu}}}\exp\left[-\frac{1}{2}\left(\sqrt{q_{\mu}} - \frac{(\mu-\mu')}{\sigma}\right)^2\right]$$

$$f(q_{\mu}|\mu) = \frac{1}{2}\delta(q_{\mu}) + \frac{1}{2}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{q_{\mu}}}e^{-q_{\mu}/2}$$

$$F(q_{\mu}|\mu') = \Phi\left(\sqrt{q_{\mu}} - \frac{(\mu - \mu')}{\sigma}\right)$$

$$p_{\mu} = 1 - F(q_{\mu}|\mu) = 1 - \Phi\left(\sqrt{q_{\mu}}\right)$$

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## Distribution of $\tilde{q}_{\mu}$

Similar results for  $\tilde{q}_{\mu}$ 

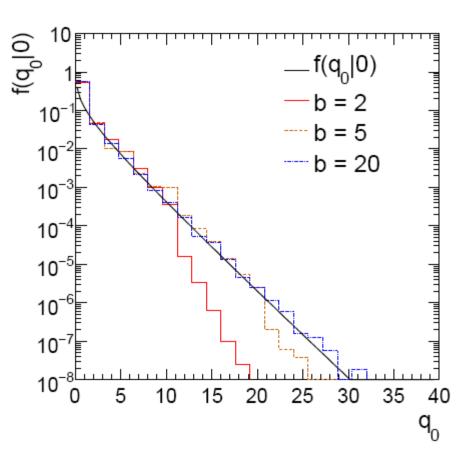
$$\begin{split} f(\tilde{q}_{\mu}|\mu') &= \Phi\left(\frac{\mu'-\mu}{\sigma}\right)\delta(\tilde{q}_{\mu}) \\ &+ \begin{cases} \frac{1}{2}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{\tilde{q}_{\mu}}}\exp\left[-\frac{1}{2}\left(\sqrt{\tilde{q}_{\mu}}-\frac{(\mu-\mu')}{\sigma}\right)^{2}\right] & 0 < \tilde{q}_{\mu} \le \mu^{2}/\sigma^{2} \\ \frac{1}{\sqrt{2\pi\sigma}}\exp\left[-\frac{1}{2}\frac{(\tilde{q}_{\mu}-(\mu^{2}-2\mu\mu')/\sigma^{2})^{2}}{(2\mu/\sigma)^{2}}\right] & \tilde{q}_{\mu} > \mu^{2}/\sigma^{2} \end{cases} \\ F(\tilde{q}_{\mu}|\mu') &= \begin{cases} \Phi\left(\sqrt{\tilde{q}_{\mu}}-\frac{(\mu-\mu')}{\sigma}\right) & 0 < \tilde{q}_{\mu} \le \mu^{2}/\sigma^{2} \\ \Phi\left(\frac{\tilde{q}_{\mu}-(\mu^{2}-2\mu\mu')/\sigma^{2}}{2\mu/\sigma}\right) & \tilde{q}_{\mu} > \mu^{2}/\sigma^{2} \end{cases} \end{split}$$

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### Monte Carlo test of asymptotic formula

 $n \sim ext{Poisson}(\mu s + b)$  $m \sim ext{Poisson}(\tau b)$ Here take  $\tau = 1$ .

Asymptotic formula is good approximation to  $5\sigma$ level ( $q_0 = 25$ ) already for  $b \sim 20$ .

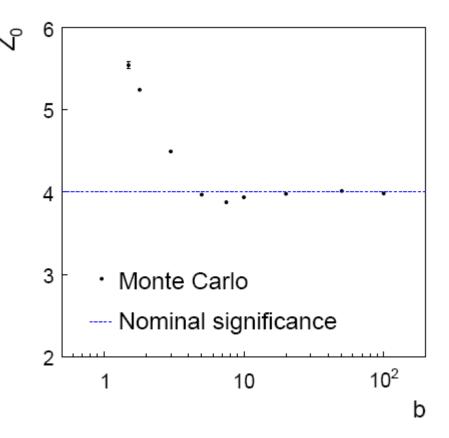


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Monte Carlo test of asymptotic formulae

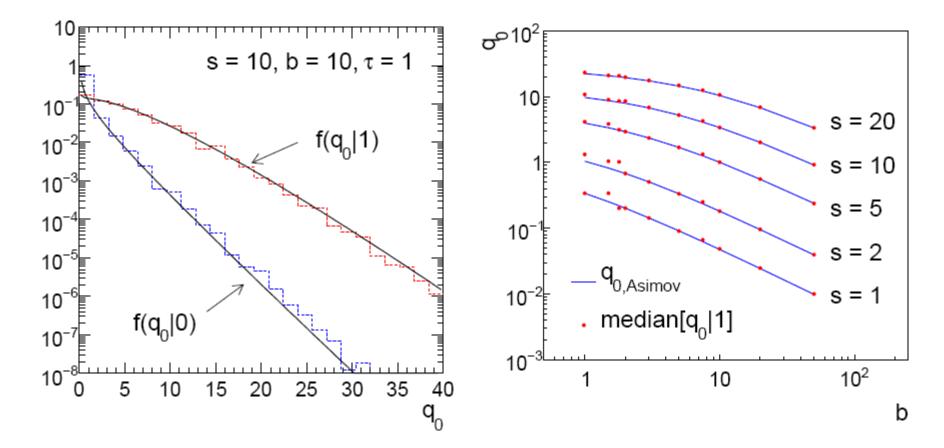
Significance from asymptotic formula, here  $Z_0 = \sqrt{q_0} = 4$ , compared to MC (true) value.

For very low b, asymptotic formula underestimates  $Z_0$ . Then slight overshoot before rapidly converging to MC value.



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Monte Carlo test of asymptotic formulae Asymptotic  $f(q_0|1)$  good already for fairly small samples. Median[ $q_0|1$ ] from Asimov data set; good agreement with MC.



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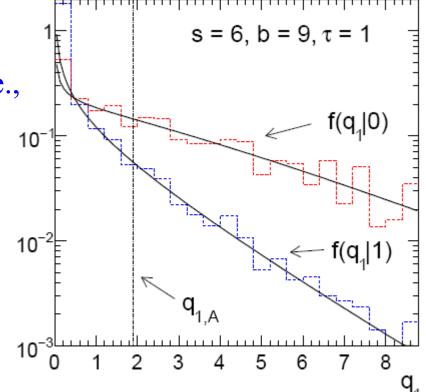
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#### Monte Carlo test of asymptotic formulae

Consider again  $n \sim \text{Poisson}(\mu s + b), m \sim \text{Poisson}(\tau b)$ Use  $q_{\mu}$  to find *p*-value of hypothesized  $\mu$  values.

E.g.  $f(q_1|1)$  for *p*-value of  $\mu = 1$ . Typically interested in 95% CL, i.e., *p*-value threshold = 0.05, i.e.,  $q_1 = 2.69$  or  $Z_1 = \sqrt{q_1} = 1.64$ . Median[ $q_1|0$ ] gives "exclusion sensitivity". Here asymptotic formulae good

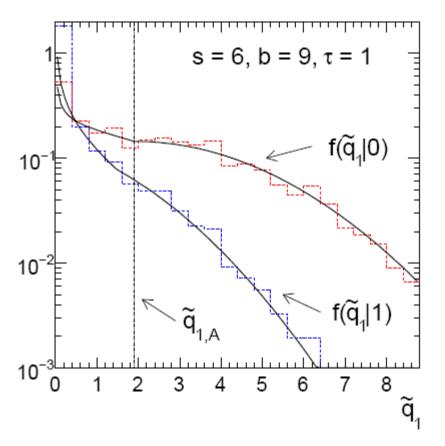
for s = 6, b = 9.



#### Monte Carlo test of asymptotic formulae

Same message for test based on  $\widetilde{q}_{\mu}$ .

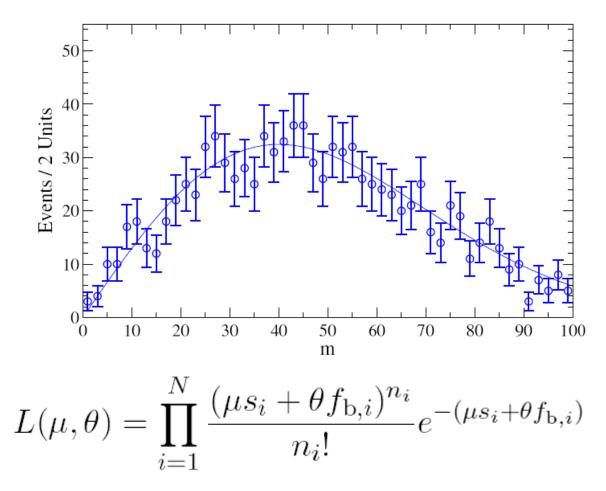
 $q_{\mu}$  and  $\tilde{q}_{\mu}$  give similar tests to the extent that asymptotic formulae are valid.



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### Example 2: Shape analysis

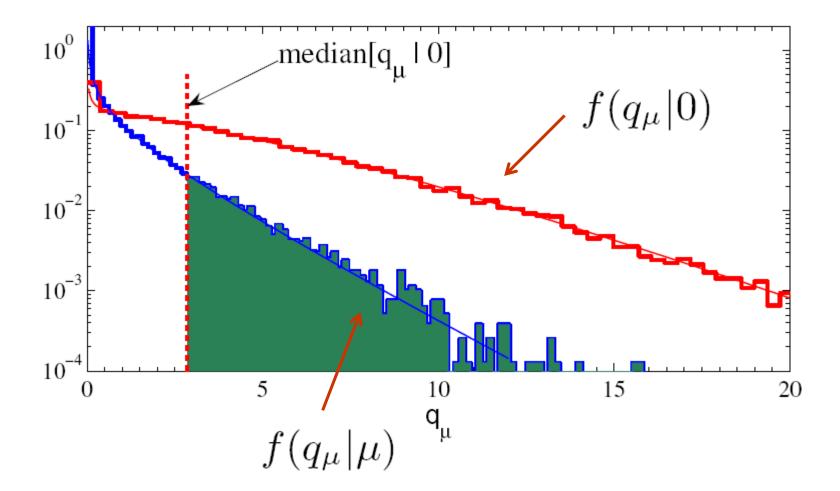
#### Look for a Gaussian bump sitting on top of:



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#### Monte Carlo test of asymptotic formulae

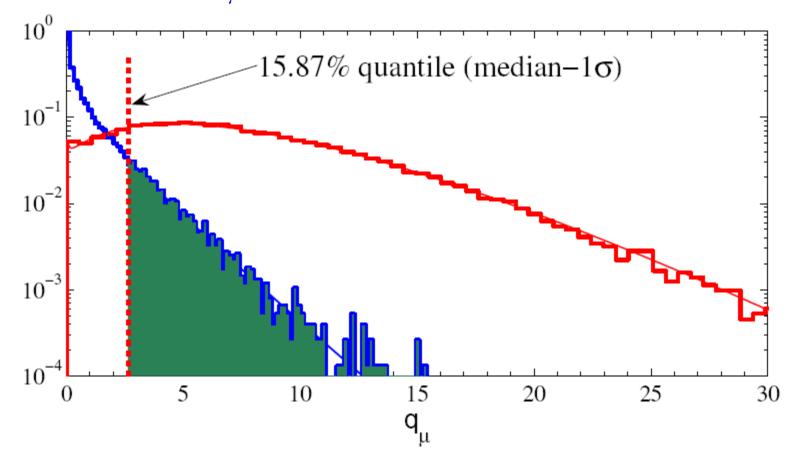
Distributions of  $q_{\mu}$  here for  $\mu$  that gave  $p_{\mu} = 0.05$ .



## Using $f(q_{\mu}|0)$ to get error bands

We are not only interested in the median[qm|0]; we want to know how much statistical variation to expect from a real data set.

But we have full  $f(q_{\mu}|0)$ ; we can get any desired quantiles.

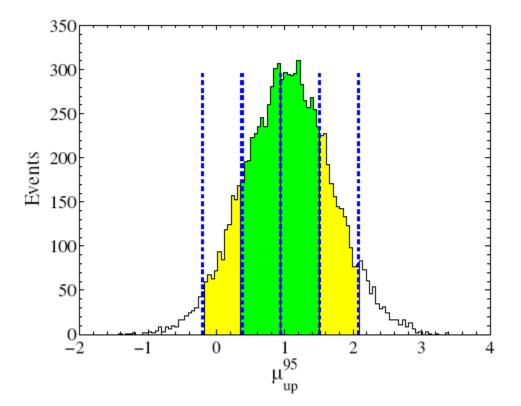


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#### Distribution of upper limit on $\mu$

 $\pm 1\sigma$  (green) and  $\pm 2\sigma$  (yellow) bands from MC; Vertical lines from asymptotic formulae

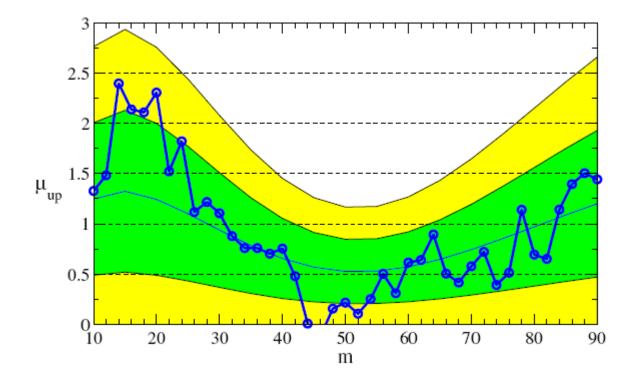


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### Limit on $\mu$ versus peak position (mass)

 $\pm 1\sigma$  (green) and  $\pm 2\sigma$  (yellow) bands from asymptotic formulae; Points are from a single arbitrary data set.



### Summary

Asymptotic distributions of profile LR applied to an LHC search. Wilks:  $f(q_{\mu}|\mu)$  for *p*-value of  $\mu$ . Wald approximation for  $f(q_{\mu}|\mu')$ .

"Asimov" data set used to estimate median  $q_{\mu}$  for sensitivity.

Gives  $\sigma$  of distribution of estimator for  $\mu$ .

Asymptotic formulae especially useful for estimating sensitivity in high-dimensional parameter space.

Can always check with MC for very low data samples and/or when precision crucial.

Implementation in RooStats (ongoing)

#### Extra slides

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Profile likelihood ratio for unified interval

We can also use directly

 $t_{\mu} = -2 \ln \lambda(\mu)$  where  $\lambda(\mu) = \frac{L(\mu, \hat{\theta})}{L(\hat{\mu}, \hat{\theta})}$ 

as a test statistic for a hypothesized  $\mu$ .

Large discrepancy between data and hypothesis can correspond either to the estimate for  $\mu$  being observed high or low relative to  $\mu$ .

This is essentially the statistic used for Feldman-Cousins intervals (here also treats nuisance parameters).

# Distribution of $t_{\mu}$

Using Wald approximation,  $f(t_{\mu}|\mu')$  is noncentral chi-square for one degree of freedom:

$$f(t_{\mu}|\mu') = \frac{1}{2\sqrt{t_{\mu}}} \frac{1}{\sqrt{2\pi}} \left[ \exp\left(-\frac{1}{2}\left(\sqrt{t_{\mu}} + \frac{\mu - \mu'}{\sigma}\right)^2\right) + \exp\left(-\frac{1}{2}\left(\sqrt{t_{\mu}} - \frac{\mu - \mu'}{\sigma}\right)^2\right) \right]$$

Special case of  $\mu = \mu'$  is chi-square for one d.o.f. (Wilks).

The *p*-value for an observed value of  $t_{\mu}$  is

$$p_{\mu} = 1 - F(t_{\mu}|\mu) = 2\left(1 - \Phi\left(\sqrt{t_{\mu}}\right)\right)$$

and the corresponding significance is

$$Z_{\mu} = \Phi^{-1}(1 - p_{\mu}) = \Phi^{-1} \left( 2\Phi \left( \sqrt{t_{\mu}} \right) - 1 \right)$$

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## Combination of channels

For a set of independent decay channels, full likelihood function is product of the individual ones:

$$L(\mu, \theta) = \prod_i L_i(\mu, \theta_i)$$

For combination need to form the full function and maximize to find estimators of  $\mu$ ,  $\theta$ .

→ ongoing ATLAS/CMS effort with RooStats framework

Trick for median significance: estimator for  $\mu$  is equal to the Asimov value  $\mu'$  for all channels separately, so for combination,

$$\lambda_{A}(\mu) = \prod_{i} \lambda_{A,i}(\mu)$$
 where  $\lambda_{A,i}(\mu) = \frac{L_{i}(\mu, \hat{\theta})}{L_{i}(\mu', \theta)}$ 

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## Discovery significance for $n \sim Poisson(s + b)$

Consider again the case where we observe n events, model as following Poisson distribution with mean s + b(assume b is known).

- 1) For an observed *n*, what is the significance  $Z_0$  with which we would reject the s = 0 hypothesis?
- 2) What is the expected (or more precisely, median )  $Z_0$  if the true value of the signal rate is *s*?

Gaussian approximation for Poisson significance For large s + b,  $n \rightarrow x \sim \text{Gaussian}(\mu, \sigma)$ ,  $\mu = s + b$ ,  $\sigma = \sqrt{(s + b)}$ . For observed value  $x_{\text{obs}}$ , *p*-value of s = 0 is  $\text{Prob}(x > x_{\text{obs}} | s = 0)$ ,:

$$p_0 = 1 - \Phi\left(\frac{x_{\rm obs} - b}{\sqrt{b}}\right)$$

Significance for rejecting s = 0 is therefore

$$Z_0 = \Phi^{-1}(1 - p_0) = \frac{x_{\text{obs}} - b}{\sqrt{b}}$$

Expected (median) significance assuming signal rate s is

$$\mathrm{median}[Z_0|s+b] = \frac{s}{\sqrt{b}}$$

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Better approximation for Poisson significance

Likelihood function for parameter *s* is

$$L(s) = \frac{(s+b)^n}{n!}e^{-(s+b)}$$

or equivalently the log-likelihood is

$$\ln L(s) = n \ln(s+b) - (s+b) - \ln n!$$

Find the maximum by setting  $\frac{\partial n}{\partial t}$ 

$$\frac{\partial \ln L}{\partial s} = 0$$

gives the estimator for s:  $\hat{s} = n - b$ 

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Approximate Poisson significance (continued) The likelihood ratio statistic for testing s = 0 is

$$q_0 = -2\ln\frac{L(0)}{L(\hat{s})} = 2\left(n\ln\frac{n}{b} + b - n\right) \quad \text{for } n > b, \ 0 \text{ otherwise}$$

For sufficiently large s + b, (use Wilks' theorem),

$$Z_0 \approx \sqrt{q_0} = \sqrt{2\left(n\ln\frac{n}{b} + b - n\right)}$$
 for  $n > b$ , 0 otherwise

To find median[ $Z_0|s+b$ ], let  $n \rightarrow s+b$ ,

$$\mathrm{median}[Z_0|s+b] \approx \sqrt{2\left((s+b)\ln(1+s/b)-s\right)}$$

This reduces to  $s/\sqrt{b}$  for s << b.

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# Higgs search with profile likelihood

Combination of Higgs boson search channels (ATLAS) Expected Performance of the ATLAS Experiment: Detector, Trigger and Physics, arXiv:0901.0512, CERN-OPEN-2008-20.

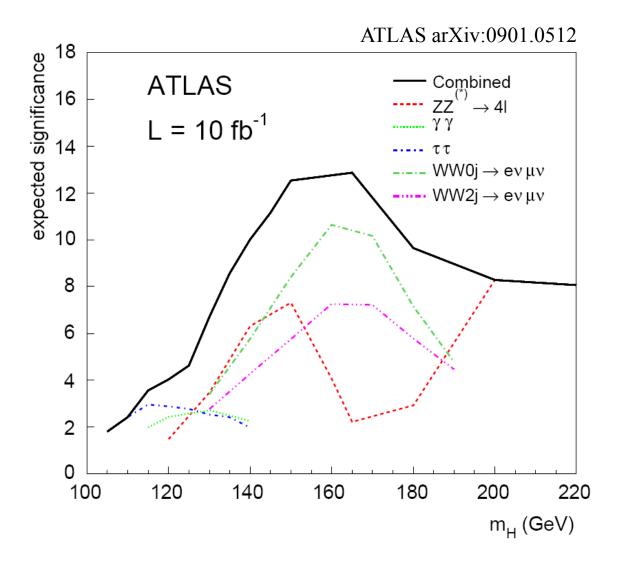
Standard Model Higgs channels considered (more to be used later):

$$\begin{split} H &\to \gamma \gamma \\ H &\to WW^{(*)} \to e\nu\mu\nu \\ H &\to ZZ^{(*)} \to 4l \ (l = e, \mu) \\ H &\to \tau^{+}\tau^{-} \to ll, lh \end{split}$$

Used profile likelihood method for systematic uncertainties: background rates, signal & background shapes.

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## Combined median significance



N.B. illustrates statistical method, but study did not include all usable Higgs channels.

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## An example: ATLAS Higgs search (ATLAS Collab., CERN-OPEN-2008-020)

Statistical Combination of Several Important Standard Model Higgs Boson Search Channels.

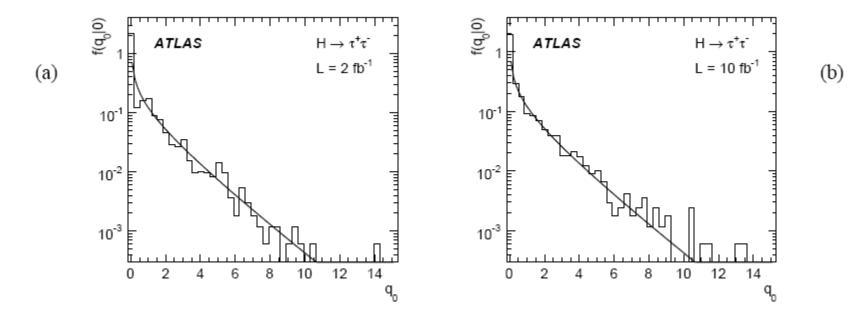


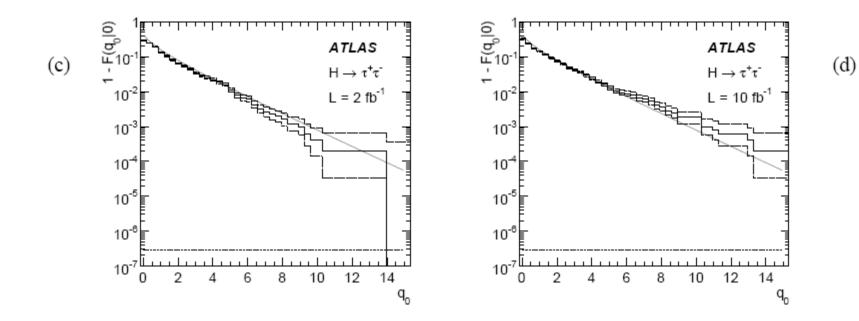
Figure 12: The distribution of the test statistic  $q_0$  for  $H \rightarrow \tau^+ \tau^-$  under the null background-only hypothesis, for  $m_H = 130 \text{ GeV}$  with an integrated luminosity of 2 (a) and 10 (b) fb<sup>-1</sup>. A  $\frac{1}{2}\chi_1^2$  distribution is superimposed. Figures (c) and (d) show  $1 - F(q_0)$  where  $F(q_0)$  is the corresponding cumulative distribution. The small excess of events at high  $q_0$  is statistically compatible with the expected curves, as can be seen by comparison with the dotted histograms that show the 68.3% central confidence intervals for  $p = 1 - F(q_0|0)$ . The lower dotted line at  $2.87 \times 10^{-7}$  shows the 5 $\sigma$  discovery threshold.

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### Cumulative distributions of $q_0$

To validate to  $5\sigma$  level, need distribution out to  $q_0 = 25$ , i.e., around 10<sup>8</sup> simulated experiments.

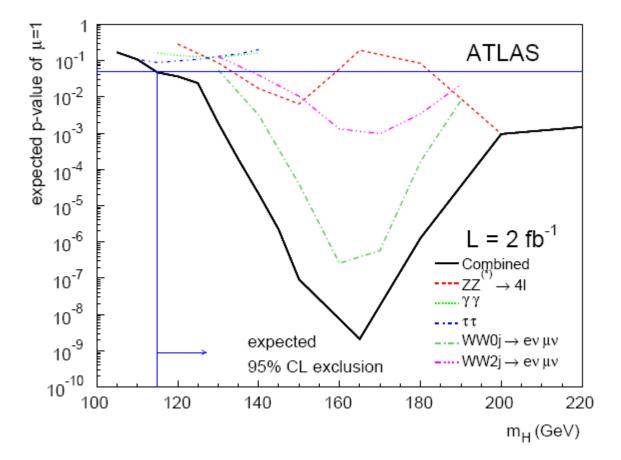
Will do this if we really see something like a discovery.



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### Example: exclusion sensitivity

Median *p*-value of  $\mu = 1$  hypothesis versus Higgs mass assuming background-only data (ATLAS, arXiv:0901.0512).



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### Confidence intervals by inverting a test

Confidence intervals for a parameter  $\theta$  can be found by defining a test of the hypothesized value  $\theta$  (do this for all  $\theta$ ):

Specify values of the data that are 'disfavoured' by  $\theta$  (critical region) such that *P*(data in critical region)  $\leq \gamma$  for a prespecified  $\gamma$ , e.g., 0.05 or 0.1.

If data observed in the critical region, reject the value  $\theta$ .

Now invert the test to define a confidence interval as:

set of  $\theta$  values that would not be rejected in a test of size  $\gamma$  (confidence level is  $1 - \gamma$ ).

The interval will cover the true value of  $\theta$  with probability  $\geq 1 - \gamma$ . Equivalent to confidence belt construction; confidence belt is acceptance region of a test.

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Relation between confidence interval and *p*-value

Equivalently we can consider a significance test for each hypothesized value of  $\theta$ , resulting in a *p*-value,  $p_{\theta}$ .

If  $p_{\theta} < \gamma$ , then we reject  $\theta$ .

The confidence interval at  $CL = 1 - \gamma$  consists of those values of  $\theta$  that are not rejected.

E.g. an upper limit on  $\theta$  is the greatest value for which  $p_{\theta} \ge \gamma$ .

In practice find by setting  $p_{\theta} = \gamma$  and solve for  $\theta$ .

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