

Stabilized finite elements for Darcy flow and application to hydrothermal flows

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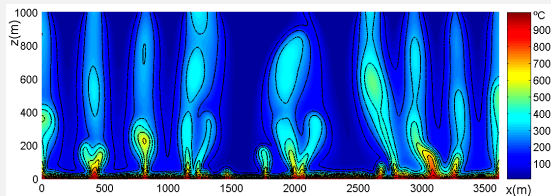
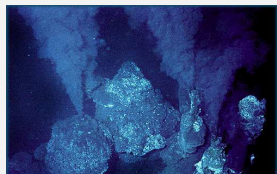
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- 2 Local projection stabilization (LPS) for Darcy-Brinkman
- 3 Stability and a priori error estimates
- 4 Model problem
- 5 Application to black smokers
 - Effect of numerical scheme to statistical quantities

1. Introduction

Black smokers



- Undersea volcanoes are interesting due to biochemical mineral formation
- Numerical simulation may help to understand their formation:
 - dynamical processes in the ocean's crust (porous media)
 - PDEs with nonlinear mathematical models for the coefficients
- Numerical schemes should be improved when model complexity increases

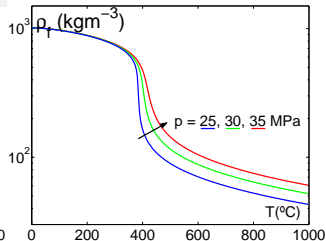
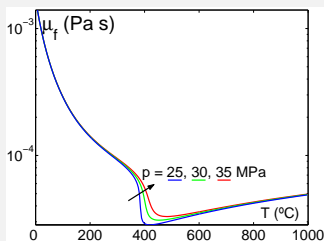
Governing equations

Darcy + temperature eq:

$$\phi \frac{\partial \rho_f}{\partial t} + \nabla \cdot (\rho_f \mathbf{v}) = 0$$
$$\frac{\mu_f}{k} \mathbf{v} + \nabla p = \rho_f \mathbf{g}$$

$$(\phi \rho_f c_{p_f} + (1 - \phi) \rho_r c_{p_r}) \frac{\partial T}{\partial t} + (\rho_f c_{p_f} \mathbf{v} \cdot \nabla) T - \nabla \cdot (D \nabla T) = 0$$

- **Model I:** $\mu = \text{const}$, $\rho_f = \rho_0 + \alpha(T - T_0) + \beta(p - p_0)$
- **Model II:** variable μ_f , c_p 's and α, β :



2. Darcy-Brinkman-problem

$$\begin{aligned}\sigma v - \nu \Delta v + \nabla p &= f && \text{in } \Omega, \\ \operatorname{div} v &= g && \text{in } \Omega, \\ v \cdot n &= 0 && \text{on } \partial\Omega, \\ \nu v \cdot t &= 0 && \text{on } \partial\Omega\end{aligned}$$

Compatibility conditions $\int_{\Omega} g \, dx = 0$, $\int_{\Omega} p \, dx = 0$.

Brinkman 1947: consistent bc for Darcy

Different regimes:

- $\sigma = 0$, $g = 0$, $\nu > 0$: standard stationary Stokes system
- $\sigma \sim (t_n - t_{n-1})^{-1}$, $g = 0$: one time step of non-stationary Stokes
- $\nu = 0$: Darcy case

Variational formulation

Variational spaces:

$$\begin{aligned}V &:= \{v \in H^1(\Omega)^d : v = 0 \text{ on } \partial\Omega\} && \text{for } \nu > 0, \\V &:= \{v \in L^2(\Omega)^d : \operatorname{div} v \in L^2(\Omega), v \cdot n = 0 \text{ on } \partial\Omega\} && \text{for } \nu = 0, \\X &:= V \times L_0^2(\Omega)\end{aligned}$$

Bilinear form for $u = (v, p), \varphi = (\phi, \xi) \in X$:

$$A(u, \varphi) := (\sigma v, \phi) + (\nu \nabla v, \nabla \phi) - (p, \operatorname{div} \phi) + (\operatorname{div} v, \xi)$$

Uniformly stable and accurate fem with respect to σ and ν ?

- Stable families of finite elements for Stokes and Darcy are usually not the same.
- $H(\operatorname{div})$ -conforming FE for Darcy are not suitable for Stokes (Mardal et al. 2002)
- Stokes MINI-element not suitable for Darcy

Heterogeneous approaches:

- Layton, Schieweck, Yotov 2003:
RT or Brezzi-Douglas-Marini for Darcy, Taylor-Hood or Mini for Stokes

Unified approaches:

- Karper, Mardal, Winther 2006: **non-conforming FE** of Crouzeix-Raviart type
- Burman, Hansbo 2007: P_1/P_0 +edge stab.
- Badia, Codina 2009: ASGS and OSS

Aim:

- Discretization with equal-order elements
- stable and accurate for $\nu \in [0, \infty)$, in particular for $\nu = 0$
- stable and accurate for $\sigma \in [0, \infty)$
- uniform approach

Local projection stabilization (LPS)

- Replace X by discrete space $X_h = V_h \times Q_h$:

$$u_h \in X_h : \quad A(u_h, \varphi) + S_h(u_h, \varphi) = F(\varphi) \quad \forall \varphi \in X_h$$

- Symmetric stabilization term

$$S_h(u, z) := (\delta \kappa_h(\operatorname{div} v), \kappa_h(\operatorname{div} w)) + (\alpha \kappa_h(\nabla p), \kappa_h(\nabla q))$$

Fluctuation operator $\kappa_h := I - \pi_h$ and local L^2 -projection π_h :

$$\pi_h : L^2(\Omega) \rightarrow D_h$$

D_h : Inf-sup condition, locally L^2 -stable, interpolation property.

- Two-level method: $D_h := Q_{2h, r-1}^{dc}$
Several one-level methods for Oseen (Matthies, Skrzypacz, Tobiska)
- Appropriate scaling of stabilization parameters in dependence of h , σ and ν .

3. Stability and a priori error estimates

- We show an inf-sup condition with respect to the triplnorm ($\theta > 0$):

$$\|u\|^2 := \underbrace{\nu|v|_1^2 + \sigma\|v\|_0^2 + S_h(y, y)}_{=|\cdot|_a^2} + \underbrace{\theta^2\|p\|_0^2}_{=|\cdot|_b^2}.$$

θ will be fixed later, but h independent.

- Stability with respect to $|\cdot|_a$:

$$A_h(u_h, u_h) := A(u_h, u_h) + S_h(u_h, u_h) = |u_h|_a^2 \quad \forall u_h \in X_h$$

- How do we obtain stability with respect to $|\cdot|_b$?

Abstract stability result

Let $(H, \|\cdot\|_H)$ be a Hilbert space where $\|\cdot\|_H$ is defined by means of two semi-norms $|\cdot|_a$ and $|\cdot|_b$

$$\|y\|_H^2 := |y|_a^2 + |y|_b^2$$

Lemma (Br., Schieweck 2010)

Let $B : H \times H \rightarrow \mathbb{R}$ bilinear with

$$\begin{aligned} \forall y \in H : \quad & B(y, y) \geq c_0 |y|_a^2 \\ \forall y \in H \exists x \in H : \quad & B(y, x) \geq c_2 |y|_b^2 - c_1 |y|_a^2 \quad \text{and} \quad \|x\|_H \leq \|y\|_H \end{aligned}$$

with $c_0, c_2 > 0$ and $c_1 \geq 0$. Then there is $\gamma > 0$ s.t.

$$\forall y \in H \exists z \in H \setminus \{0\} : \quad B(y, z) \geq \gamma \|y\|_H \|z\|_H$$

$$\gamma = \min \left\{ \frac{c_2}{1 + \varrho}, \frac{c_2}{1 + (c_1 + c_2)/c_0} \right\} > 0,$$

where $\varrho > 1$ is arbitrary.

Stability of Darcy-Brinkman-LPS

Proposition

For patch-wise constants

$$0 < \alpha_0 h_K^2 \leq \alpha_K \quad \text{and} \quad 0 \leq \delta_K \leq \bar{\delta} \quad \forall K \in \mathcal{T}_{2h}.$$

the stabilized Darcy-Brinkman bilinear form satisfies the inf-sup condition

$$\forall y_h \in X_h \exists z_h \in X_h \setminus \{0\} : \quad A_h(y_h, z_h) \geq \frac{1}{2} \|y_h\| \|z_h\|$$

with $\theta > 0$, and for the the corresponding discrete problem there exists always a unique solution.

$$\theta = \frac{4\gamma_c}{5c_i \alpha_0^{-2} + 9c_s (\sigma + \nu + c_\kappa^2 \bar{\delta})^{1/2}} > 0,$$

c_i, c_s, c_κ depend on κ_h but not on h .

Proof. Using previous Lemma and continuous inf-sup condition.

Proposition (A priori estimate)

Under the assumption $v \in H^{r+1}(\Omega)^d$ and $p \in H^{r+1}(\Omega)$, the choice of stabilization constants

$$\alpha = (\sigma + h^{-2}\nu)^{-1} \quad \text{and} \quad \delta = \sigma h^2,$$

leads to the a priori bound

$$\|u - u_h\| \lesssim \varrho_h h^r |v|_{r+1} + \varrho_h^{-1} h^{r+1} |p|_{r+1},$$

where $\varrho_h := \sigma^{1/2} h + \nu^{1/2}$.

Proof.

- Standard techniques on bases of the discrete inf-sup condition for arbitrary coefficients.
- "Optimization" of the parameters.

Extreme cases:

- **Stokes:** ($\alpha = h^2\nu^{-1}$, $\delta = 0$)

$$\|u\|^2 = \nu|v|_1^2 + \theta^2\|p\|_0^2 + S_h(u, u)$$

A priori est.:

$$\|u - u_h\| \lesssim \nu^{1/2}h^r|v|_{r+1} + \nu^{-1/2}h^r|p|_r$$

optimal

- **Darcy:** ($\alpha = \sigma^{-1}$, $\delta = \sigma h^2$)

$$\|u\|^2 = \sigma\|v\|_0^2 + \theta^2\|p\|_0^2 + S_h(u, u)$$

A priori est.:

$$\|u - u_h\| \lesssim \sigma^{1/2}h^r|v|_r + \sigma^{-1/2}h^r|p|_{r+1}$$

Not yet optimal for L^2 -errors \rightarrow duality arguments needed

Duality argument (1)

(R1) For $\varphi \in \{\mathbf{0}\} \times L_0^2(\Omega)$, the solution $z \in Y$ of

$$A(w, z) = \langle \varphi, w \rangle \quad \forall w \in Y$$

is in $H^1(\Omega)^d \times H^2(\Omega)$ and it holds the estimate

$$\|z_v\|_1 + \sigma^{-1} \|z_p\|_2 \lesssim \|\varphi_p\|.$$

- Example: $\nu = 0$, $\sigma > 0$, $\Omega \subset \mathbb{R}^d$, $d \in \{2, 3\}$, is a open convex polyhedron or has C^2 -boundary \implies (R1).

Proposition

Let $\nu = 0$, $\sigma > 0$ and assume that the regularity property (R1) is satisfied. Then

$$\|p - p_h\| \lesssim \sigma h^{r+1} |v|_r + h^{r+1} |p|_{r+1}.$$

Duality argument (2)

(R2) For $\varphi \in L^2(\Omega)^d \times \{0\}$, the solution $z \in Y$ of

$$A(w, z) = \langle \varphi, w \rangle \quad \forall w \in Y$$

is in $H^2(\Omega)^d \times H^1(\Omega)$ and it holds the estimate

$$\nu \|z_v\|_2 + \|z_p\|_1 \lesssim \|\varphi_v\|.$$

- Example: $\nu > 0$, $\sigma \geq 0$, $\Omega \subset \mathbb{R}^d$, $d \in \{2, 3\}$, is a open convex polyhedron or has C^2 -boundary \implies (R2).

Proposition

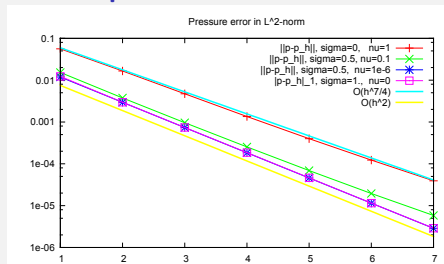
Let $\nu > 0$, $\sigma \geq 0$ and assume that the regularity property (R2) is satisfied. Then

$$\|v - v_h\| \lesssim h^{r+1} \left(\left(1 + \frac{\sigma}{\nu} h^2\right) |v|_{r+1} + \frac{h}{\nu} |p|_{r+1} \right).$$

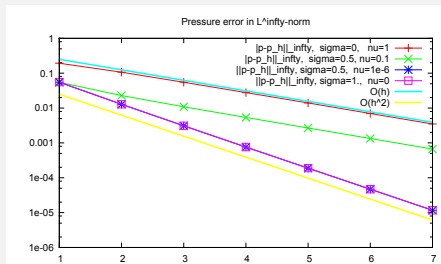
4. Numerical test

- Smooth (not divergence-free) exact solution.
- Q_1/Q_1 FEM with varying ν, σ and varying h .

Error in pressure



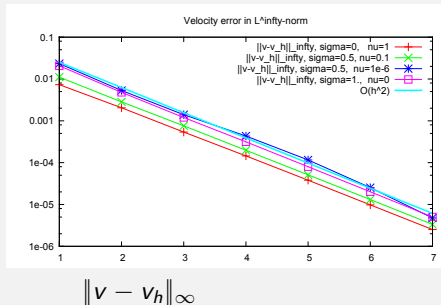
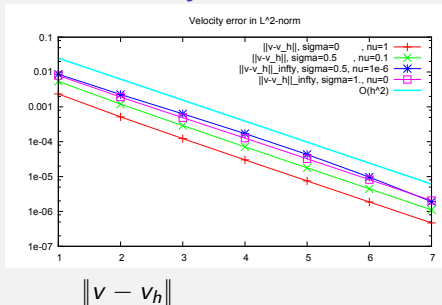
$$\|p - p_h\|$$



$$\|p - p_h\|_\infty$$

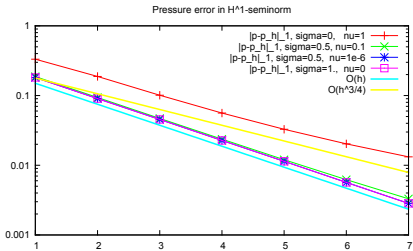
- $\nu = 0$ (Darcy): $\mathcal{O}(h^2)$ in both norms (as theory).
- $\nu = 10^{-6}$: $\mathcal{O}(h^2)$ in both norms, optimal order
- $\nu = 1$ and $\sigma = 0$ (Stokes): $\|p - p_h\| = \mathcal{O}(h^{7/4})$, $\|p - p_h\|_\infty = \mathcal{O}(h)$

Error in velocity

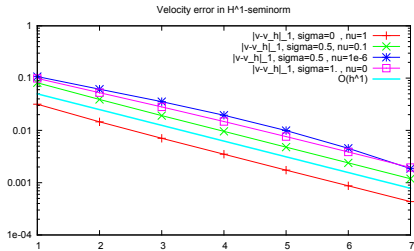


- $\mathcal{O}(h^2)$ in both norms completely independent of the parameters σ and ν .
- Optimal order of convergence as in theory ($\nu > 0$) and even for $\nu = 0$.

Errors in H^1 -seminorms



$$\|\nabla(p - p_h)\|$$



$$\|\nabla(v - v_h)\|$$

- $\|\nabla(v - v_h)\| = \mathcal{O}(h)$ independent of the parameters σ and ν .
- $\|\nabla(p - p_h)\| = \mathcal{O}(h)$ for $\sigma > 0$ and $\|\nabla(p - p_h)\| = \mathcal{O}(h^{3/4})$ for $\sigma = 0$

Numerical observations:

σ	ν	$\ p - p_h\ $	$\ p - p_h\ _\infty$	$ p - p_h _{H^1}$	$\ v - v_h\ $	$ v - v_h _{H^1}$
0	1	$h^{7/4}$	h	$h^{3/4}$	h^2	h
0.5	0.1	$h^{7/4}$	h	h	h^2	h
0.5	10^{-6}	h^2	h^2	h	h^2	h
1	0	h^2	h^2	h	h^2	h

Numerical observations:

σ	ν	$\ p - p_h\ $	$\ p - p_h\ _\infty$	$ p - p_h _{H^1}$	$\ v - v_h\ $	$ v - v_h _{H^1}$
0	1	$h^{7/4} h$	h	$h^{3/4}$	$h^2 h^2$	$h h$
0.5	0.1	$h^{7/4} h$	h	h	$h^2 h^2$	$h h$
0.5	10^{-6}	$h^2 h$	h^2	h	$h^2 h^2$	$h h$
1	0	$h^2 h^2$	h^2	h	$h^2 h$	h

theory

5. Numerical schemes for hydrothermal flows

BDF-2 in time

2 alternatives in space for T eq:

- **SUPG:**

$$\begin{aligned} & (\gamma^n T_h^n + (\beta^n \cdot \nabla) T_h^n, \varphi_h) + (D \nabla T_h^n, \nabla \varphi) + \\ & \sum_{K \in \mathcal{T}_h} \delta_K (\gamma^n T_h^n + (\beta^n \cdot \nabla) T_h^n - \nabla \cdot (D \nabla T_h^n) - f^n, (\beta^n \cdot \nabla) \varphi_h)_K = (f^n, \varphi), \end{aligned}$$

$$f^n = \frac{4}{3} \gamma^{n-1} T_h^{n-1} - \frac{1}{3} \gamma^{n-2} T_h^{n-2}$$

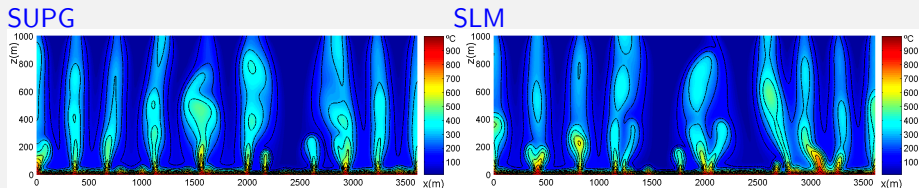
- **SLM: Semi-Lagrangian method**

$$(\gamma^n T_h^n, \varphi_h) + (D \nabla T_h^n, \nabla \varphi_h) = \left(\frac{4}{3} \gamma^{n-1} \bar{T}_h^{n-1} - \frac{1}{3} \gamma^{n-2} \bar{T}_h^{n-2}, \varphi \right)$$

$$\bar{T}_h^{n-l}(x) = T_h^{n-l}(X(x, t_n; t_{n-l})), \quad l = 1, 2.$$

SUPG vs. semi-Lagrangian method

Discrete solutions at the same time instant:



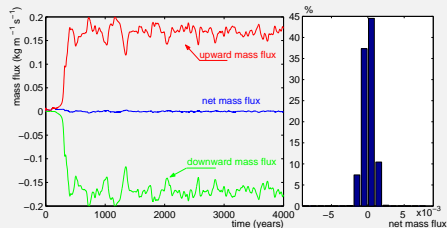
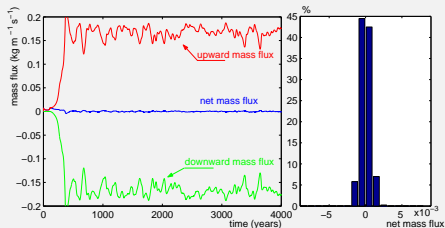
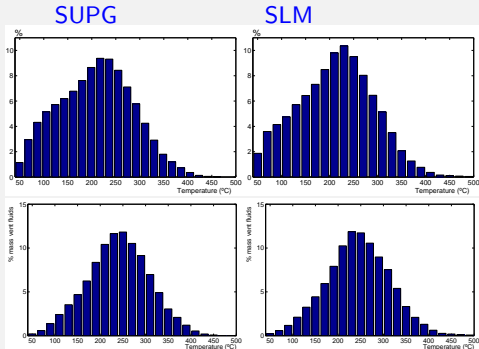
- Configuration of Coumou, Driesner, Geiger et al. (2006)
- Rel. fine mesh with 50,032 quadratic elements. Total time 4,000 a.
- System is unstable, numerical errors trigger instabilities.
- Direct comparison of the discrete solutions are not meaningful.
- One should compare statistical quantities.

Statistical quantities for constant coefficients (Model I)

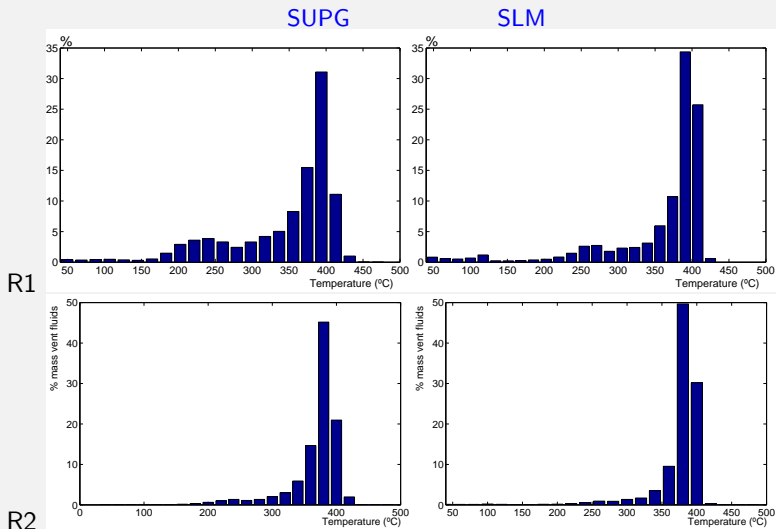
We concentrate on three different quantities:

- R_1 : T at upper bd averaged over total time and all mesh points
- R_2 : momentum $\rho_f \mathbf{v}$ at top boundary as function of T
- R_3 : up- and downward mass flux:

$$\dot{m}^{+/-} = \int_{\Gamma_{top}} (\rho_f \mathbf{v}_z)^{+/-} ds,$$



Variabel coefficients (Model II)

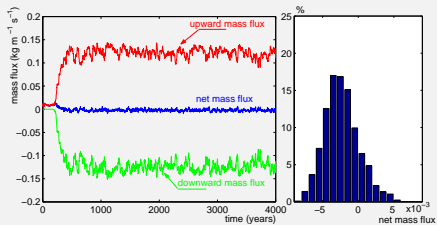
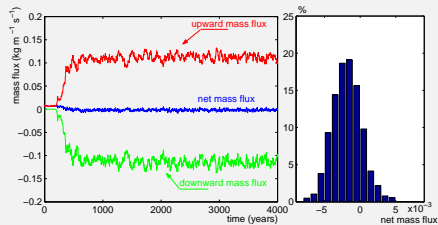


- Qualitatively similar structure
- Quantitatively: Substantial differences compared to Model I (const coef).

SUPG

SLM

R3:



Differences quantitatively:

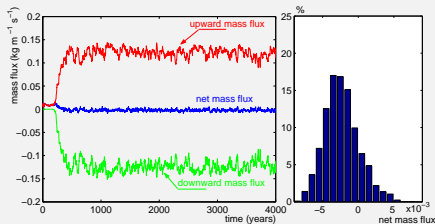
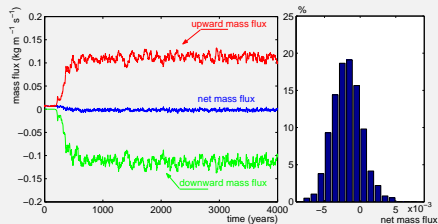
$$E_i = \sum_{k=0}^n |R_{k,SFEM} - R_{k,SLM}|$$

i	Model I	Model II
1	10.8 %	123.1 %
2	4.3 %	28.1 %

SUPG

SLM

R3:



Differences quantitatively:

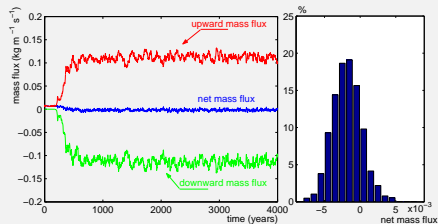
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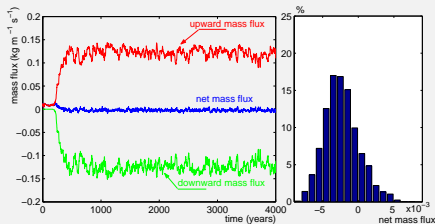
Considerable impact of numerical scheme even on statistical quantites.

SUPG

R3:



SLM



Differences quantitatively:

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Considerable impact of numerical scheme even on statistical quantites.

There is still a demand for improved discretizations for complex nonlinear models.

Conclusion

- ① Uniform approach with LPS for Darcy-Stokes, i.e. Darcy-Brinkman model,
- ② Robustness with respect to model parameters, even in extremal cases
 - for Darcy and Stokes: optimal order of convergence
- ③ FEM for hydrothermal flows:
 - ▶ **large** effect of the numerical scheme
 - ▶ **small** effect of the numerical scheme on statistical quantities for eq with constant coefficients
 - ▶ **considerable** effect of the numerical scheme on statistical quantities for eq with variable coefficients.
- ④ Improvement of numerical schemes (accuracy) has to go hand in hand with increasing model complexity.

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Thanks a lot !